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OPTICAL SIGNATURES PROGRAM FINAL REPORT
ROCKET PLUME OPTICAL SIGNATURES

October 1972

Prepared by:

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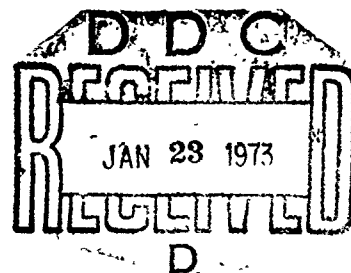
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**GENERAL
RESEARCH**



CORPORATION

WASHINGTON OPERATIONS
1501 WILSON BLVD, SUITE 700, ARLINGTON, VA 22209



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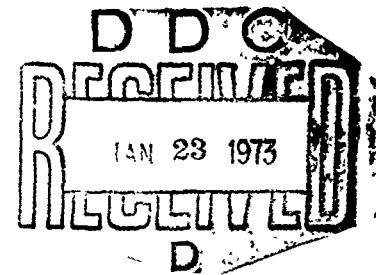
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PREFACE

This document is the final report on the "Rocket Plume Optical Signature Program." This work was performed for General Research Corporation, 1501 Wilson Boulevard, Suite 700, Arlington, Virginia 22209, under Contract 040-71-10. This contract covered the period from 20 August 1971 to 30 September 1972.

The effort described herein was accomplished by the following study personnel: S. S. Cherry, M. Thomas, and R. L. Younkin. This report was approved by H. Hurwicz, Chief Advance Technology Engineer, Aero/Thermodynamics and Nuclear Effects Research and Development, Advance Systems and Technology, McDonnell Douglas Astronautics Company-West.

ABSTRACT

As part of the Optical Signatures Program, McDonnell Douglas Astronautics Company-West has developed the initial working model of a code to describe the gross features of rocket-plume radiation for altitudes above 75 n mi. The main effort is the construction of a scheme for integration of an arbitrary function through an arbitrary axisymmetric rocket plume, with any specified look angle, plume direction, and vehicle velocity direction. Radiances are presented as integrated values in a specified spectral band. The equations used and a printout of the code and of a sample application are included.

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Section 1
INTRODUCTION AND SUMMARY

A code has been developed to describe the gross features of rocket-plume radiation for altitudes above 75 mmi (P2170-FLAME). The effort was intended to provide an initial working model, consistent with available technology. Theoretical studies were not intended to be conducted in any of the disciplines required. The intent was, however, to provide a suitable framework for the description of a plume with sufficient flexibility to incorporate the results of future research.

This objective was met by considering the two most important far-field radiative mechanisms, particle emission and atmospheric excitation, in their present state of understanding. With these major elements of plume radiation included even in crude form, future program modification to include revised theories should not require complete program revision. The main effort which has been pursued in FLAME-code development is the construction of a scheme for integration of an arbitrary function through an arbitrary axis-symmetric rocket plume, with any specified look angle, plume direction, and vehicle velocity direction. The assumption of flow symmetry about the plume axis may prove poor for late-time plumes, but to include a complete three-dimensional capability would have burdened the code's development to preclude an operational version at the end of this contract. The switch to 3-D is not difficult at a later time and will involve obtaining flow correlations as a function of the azimuthal angle, as well as r and θ , and integrating over four quadrants instead of two.

Radiances are presented in the current code as integrated values in a specified spectral band. Spectral intensities are readily available for the particles, but only a band theory is available for the atmospheric excited species, due primarily to the species cooling characteristics. Again, a switch to spectral curves is possible with ease at a later time when a more general cooling theory evolves.

A schematic of the FLAME code is presented in Figure 1-1. Radiative cooling of both the hot particles in multiphase plumes and the excited molecular radiators is included. The four boxes feeding into the FLAME code represent the main areas programmed during this contract. The use of a simplified flow model reduces both computer storage requirements and run time, and allows a more extensive checkout of the radiative mechanisms. More sophisticated flow models exist and could be substituted at a later time.

Input to the code includes simple engine parameters (Section 3, Table 3-1). Output is a plot of lines of constant radiance in the plume as viewed in the specified direction. A typical output from the code is shown in Figure 1-2 for a broadside look at a solid propellant plume in the absence of atmospheric excitation. The radiance is plotted in a plane perpendicular to the direction of view. The plume of Figure 1-2 is shown again in Figure 1-3 where the observer is at a 45-degree angle to the plume axis. Radiance from the same engine at low altitudes, including the effect of atmospheric excitation of the plume, is shown in Figure 1-4. The viewer angle, plume direction, and free stream velocity direction are all completely arbitrary. (Section 3, Figures 3-1 and 3-2.)

Only one-half the plume is treated by a single FLAME run. The remaining half can be obtained by redefining the azimuthal angles as shown in Section 3, Figure 3-2. It is not always necessary to generate both plume halves because the plume radiance is axisymmetric under certain conditions.

The radiative models used tend to underpredict the radiation. For atmospheric excitation, only single collisions are considered and the Boltzmann distribution of excited rotational states is assumed. Persistence of the plume radiance due to pumping by thermal (rather than pure kinetic) collisions with the atmosphere is not treated. For particle radiation, scattering of Earthshine is not treated.

Section 2 summarizes the theory used. Sections 3 and 4 detail the workings of the FLAME code, describing the equations used.

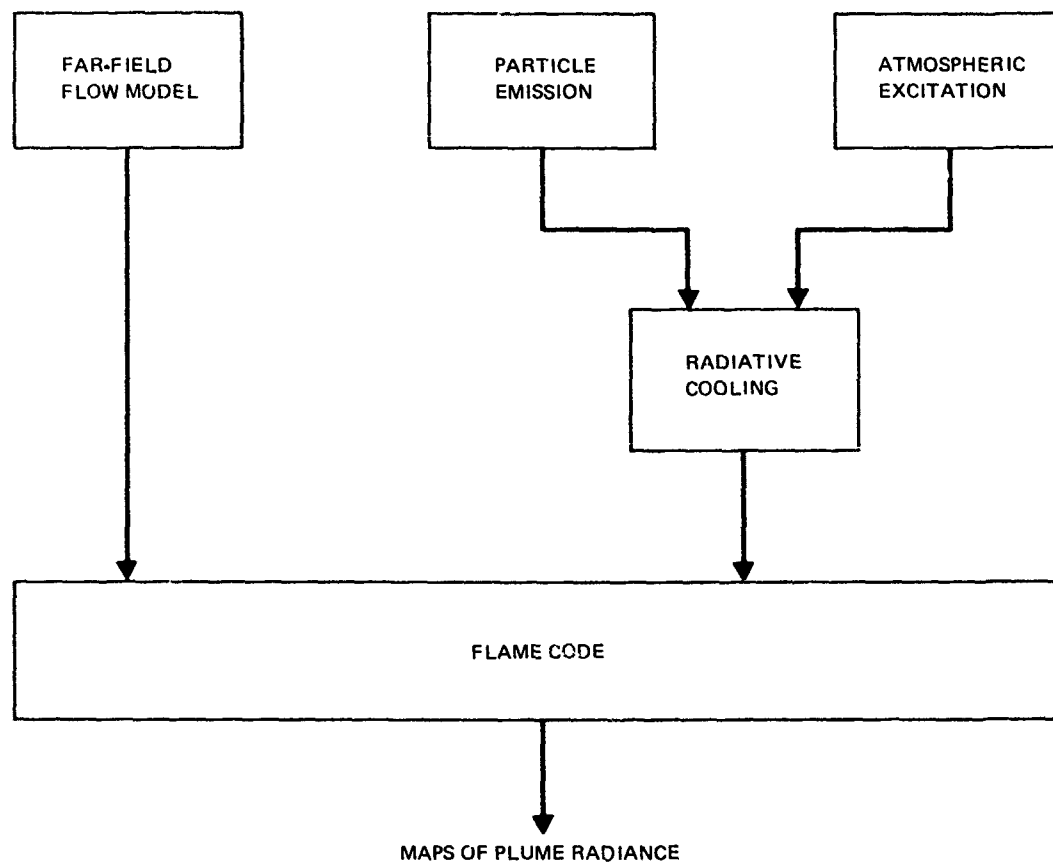


Figure 1-1. Components and Flow of FLAME Code

PLUME CONTOURS

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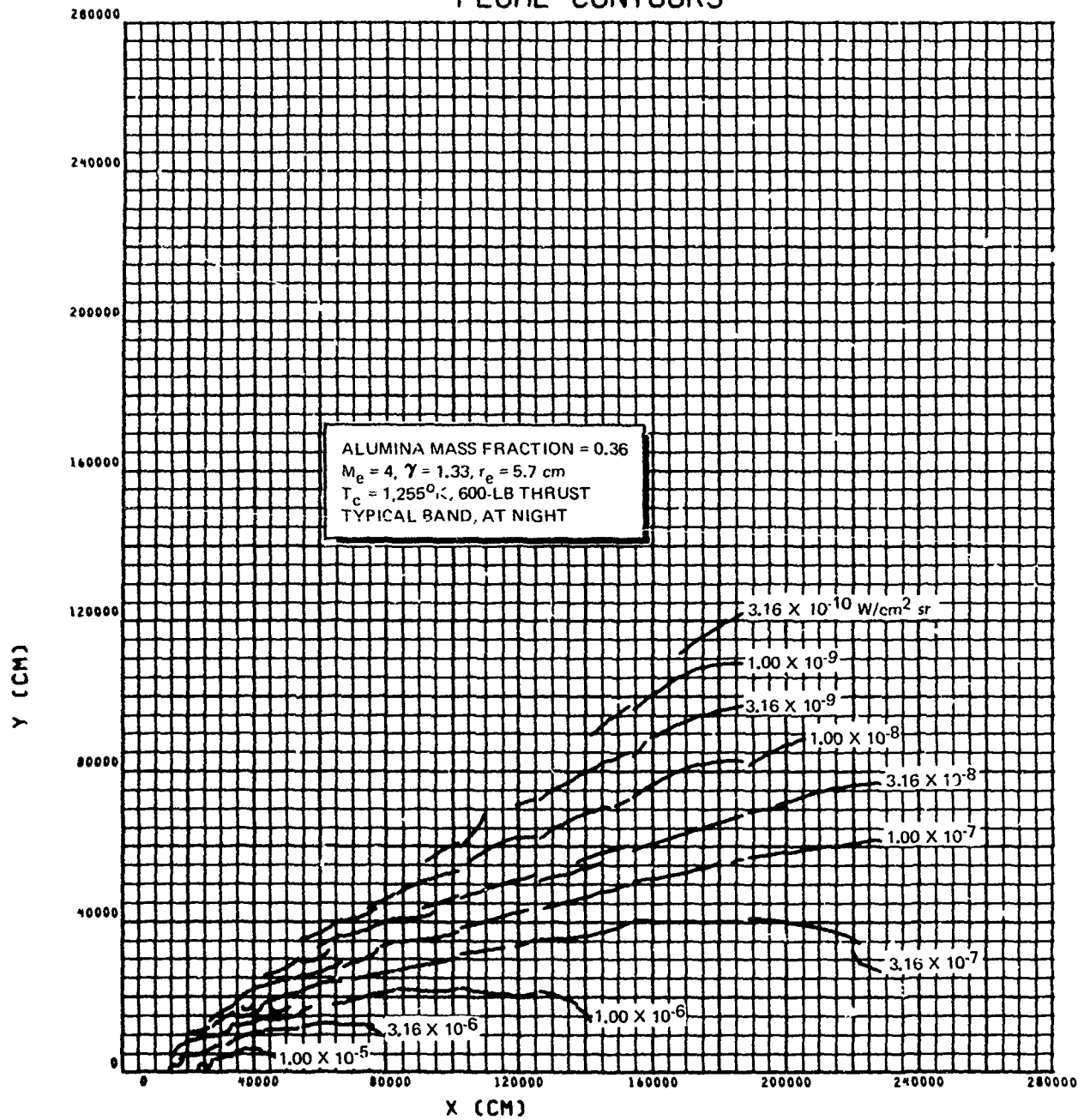


Figure 1-2. Isointensity Lines for Braodside Look at Exoatmospheric Plume of Solid Propellant Rocket

PLUME CONTOURS

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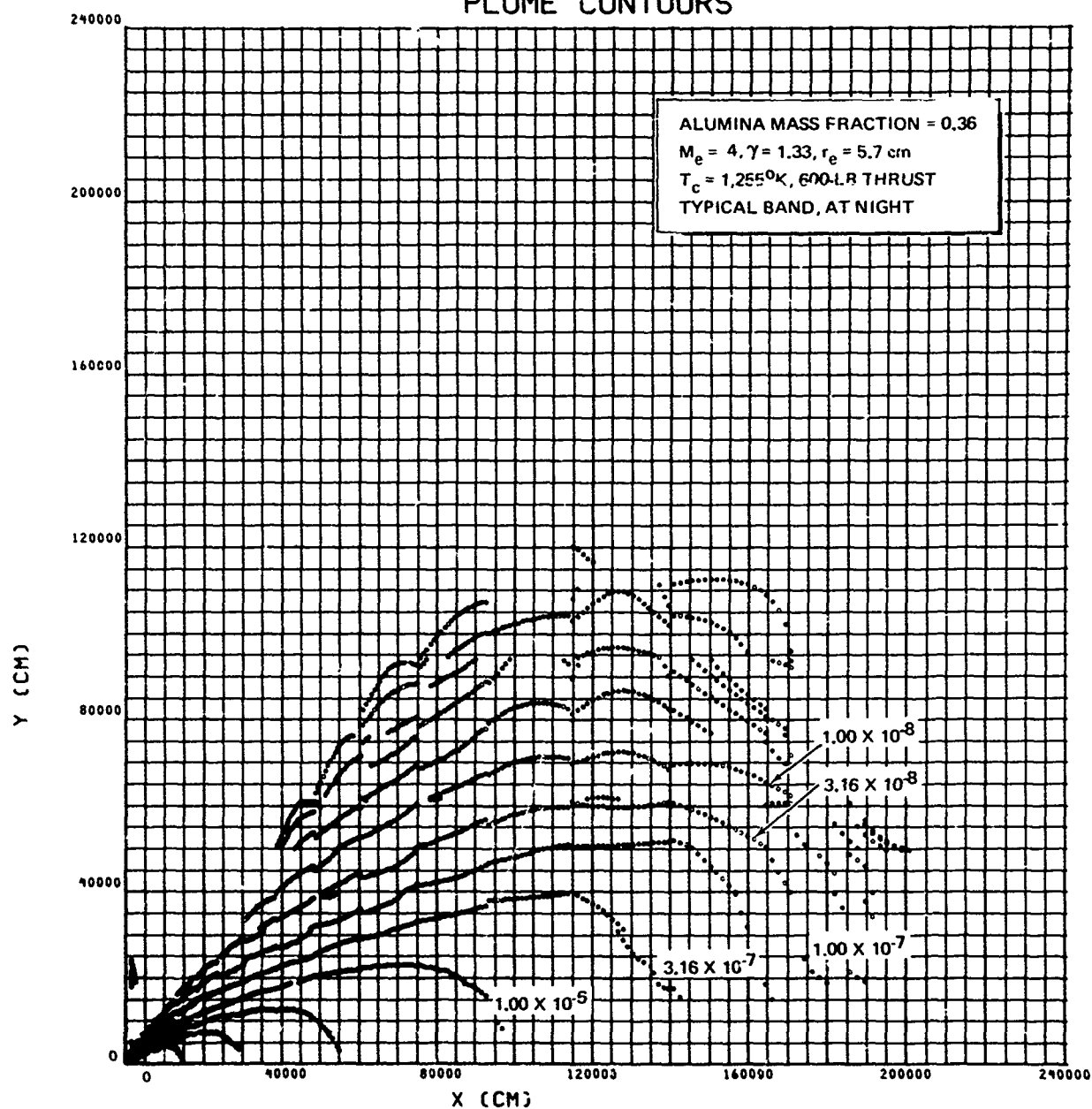


Figure 1-3. Isointensity Lines for Look at 45° from Plume Axis of Solid Propellant Rocket

PLUME CONTOURS

SERIAL A05600

2 2

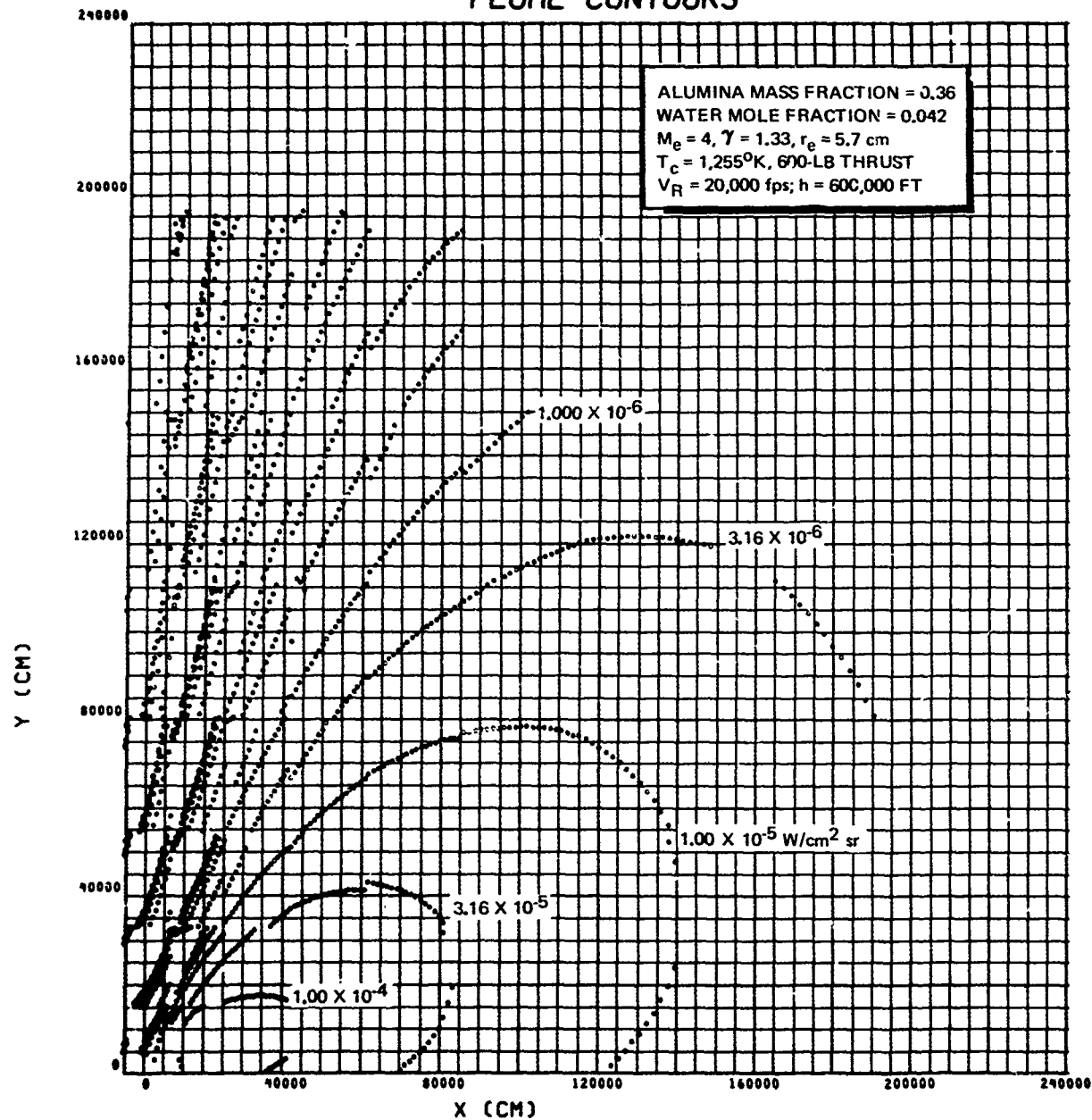


Figure 1-4. Isointensity Lines for Broadside Look at Particle and Excited Water Radiation from Solid Propellant Rocket

Section 2

THEORY

The major elements of the theory used as the basis for the FLAME code is outlined in this section. The approximate model of the far-field multiphase plume is described. For two-phase plumes, closed-form expressions for particle concentration as a function of plume coordinates and particle size are presented. The radiative mechanism for the particle cloud is simply emission, with an effective minimum temperature defined by solar illumination or scattering of Earthshine; it need not be elaborated on here. Gas radiation is considered to be caused by excitation of water rotational spectrum by collision with the atmosphere; this theory is outlined in subsection 2.3. The model can be generalized to include other radiating rotors.

The approaches used were based on existing theoretical work. No attempt has been made to derive new theoretical expressions or include radiative mechanisms not yet understood in the current FLAME. The program is flexible enough, however, to allow the inclusion of other mechanisms at some future time.

2.1 PLUME GAS FLOW

Closed-form approximation to the plume gas is relied upon for flow applicable to the far-field plume. Atmospheric mixing and the resulting plume slow down have not been addressed. Atmospheric excitation of the unaltered plume is considered.

Closed-form solutions have been developed for the density profiles as a function of engine operating conditions. It appears that all properties of the far-field gaseous plume can be thus treated. The computer program will respond to input of engine thrust (or chamber pressure), chamber temperature, expansion ratio (ϵ), γ (or exit mach number, M_e), and nozzle lip angle by generating a complete history of the rocket plume.

To predict transition to free-molecular-flow, a relationship was derived to express the Knudsen number (K_n) variation in a bipropellant rocket engine plume as a function of the polar coordinates (r, θ) and basic engine parameters F (thrust), C_F (thrust coefficient), T_c (combustion temperature), M_e (exit Mach number) and γ (ratio of specific heats). The gas viscosity, μ , was assumed to vary as $T^{1/2}$. The engine chamber pressure, P_c , could be introduced through the basic thrust equation:

$$F = P_c A_t C_F$$

where:

$$A_t = \text{Throat Area}$$

These derivations used Roberts' approximation (Reference 1) for the plume gas density which indicated that the density was inversely proportional to r^2 and proportional to $\cos \theta$ raised to the power $\gamma(\gamma - 1) M_e^{2-2\gamma}$.

The derived relationship is shown in Figure 2-1 as a function of M_e , γ , and ϵ (engine expansion area ratio). This information may be employed to determine the shape of the surface (r, θ) required to achieve a given K_n . It is noted that r increases with increasing γ for specified values of T_c , K_n , F , and θ . Indeed, this relationship indicates that $r \rightarrow 0$ as $\theta \rightarrow 90$ degrees, i. e., that transition occurs at the nozzle lip where, theoretically, there is an infinite pressure gradient as the flow undergoes a Prandtl-Meyer expansion.

Figure 2-2 presents the variation of plume density on a spherical cap ($r = \text{constant}$) as a function of θ and M_e with $\gamma = 1.4$. As shown, the density has been normalized by its value on the engine centerline, i. e., $\theta = 0$. When $M_e = 2.312$, the density distribution is linear with $\cos \theta$ and becomes increasingly dependent on $\cos \theta$ as M_e increases, e. g., for $M_e = 5.0$, $\rho \sim (\cos \theta)^{12}$.

2.2 PARTICLES

A combined nozzle and plume flowfield calculation was performed for an aluminized solid propellant motor to obtain correlations of particle characteristics in the near field. Closed-form solutions, analogous to those obtained for gas-phase-only plumes were not possible due to the nonequilibrium nature of the expansion involving momentum (velocity) and energy (temperature) differences between the phases.

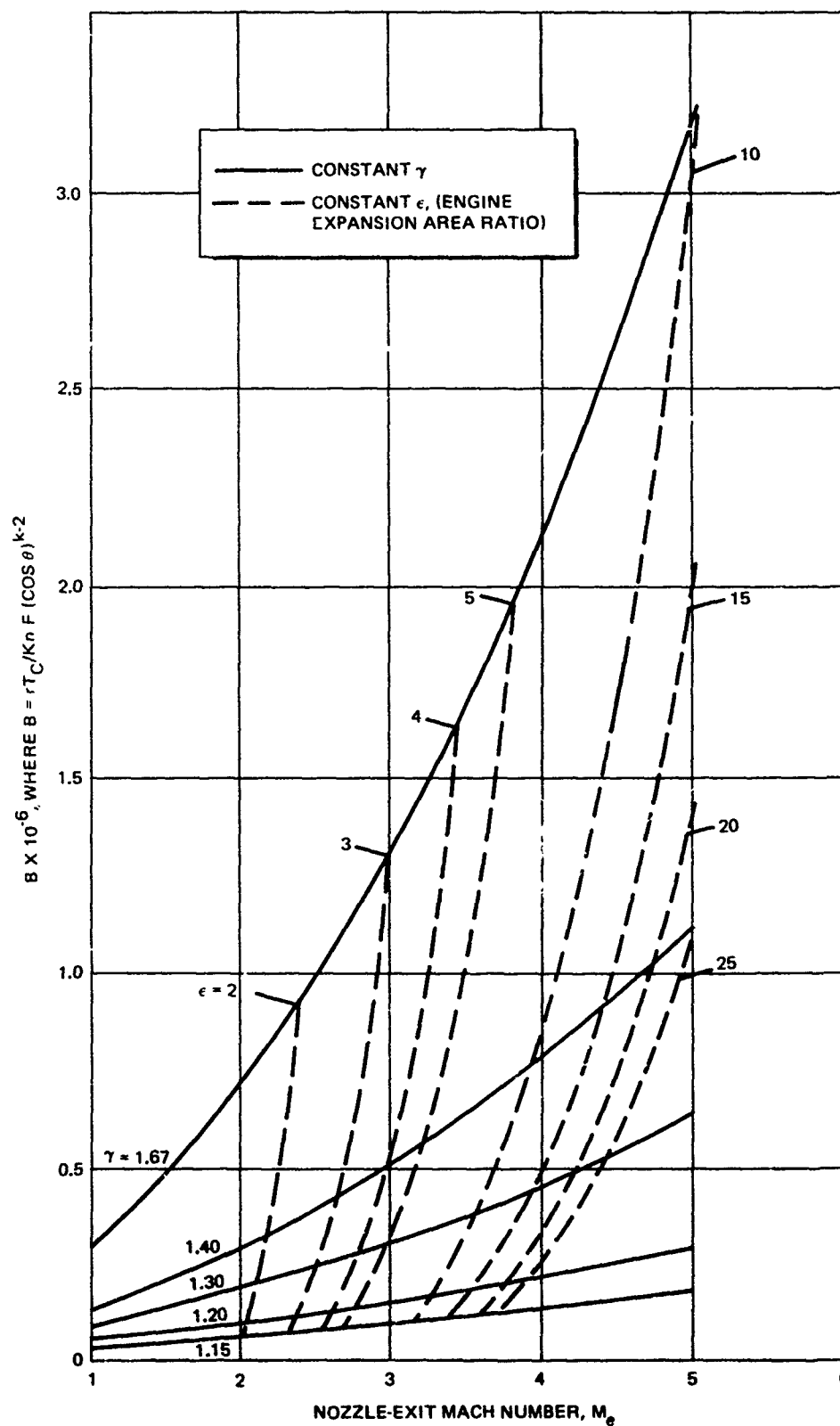


Figure 2-1. Knudsen Number Variation with Engine Conditions

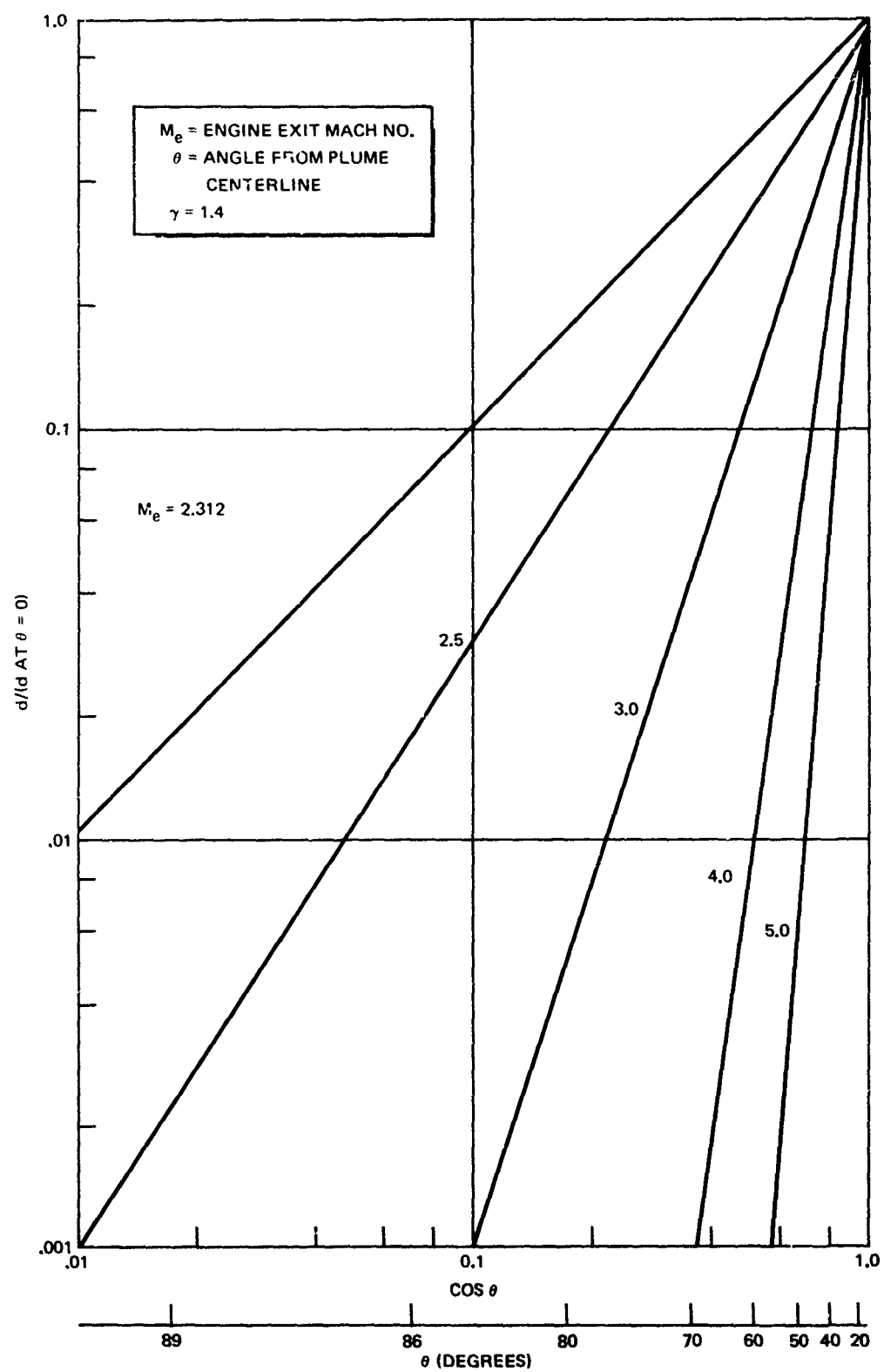


Figure 2-2. Plume Density Distribution on Spherical Cap ($r = \text{constant}$)

Correlations were obtained which related particle velocity, temperature, density, and limiting streamline direction to the seven particle-size groups considered. An expression was derived which related the particle Knudsen number to plume coordinates and motor operating conditions. The primary use of this expression was to determine the extent of the particle continuum region.

Initial estimates of particle emissivity were made, based primarily on existing calculations for alumina. The details of these analyses are presented in later subsections.

2.2.1 Two-Phase Flowfield Particle Correlations

The computer programs described in References 2 and 3 were employed to calculate the nozzle and plume flowfields for an aluminized solid-propellant motor with the following characteristics:

Thrust = 15,000 lb_f

Chamber pressure = 550 psia

Expansion area ratio = 30.8

Aluminum weight content = 16 percent

A total of seven aluminum-oxide particle groups were utilized with particle radii, ρ_i , ranging from 0.6 to 4.7 μm . The corresponding number count versus ρ_i was obtained from Reference 4.

The computed output contained particle velocity, temperature, flow directions, and nondimensional number density for each group at flowfield mesh points within the plume. This information was processed to extract certain particle asymptotic characteristics, namely; velocity, temperature, and flow direction along the particle-limiting streamlines. (The limiting streamline is defined as the boundary outside of which a given particle size will not be present.) It was found that these streamlines became straight lines after the particles had traversed a relatively short distance due to the rapid decay of interphase momentum and energy exchange.

The values of particle temperature, T_i , flow angle, θ_i , and velocity, $V(\rho_i, \theta)$, on the limiting streamlines are shown in Figure 2-3 as a function of ρ_i . It is noted that particles larger than 4μ are still undergoing the liquid-solid phase transition at the aluminum-oxide melting temperature of 4,170 R and that these particles are contained within a 23.5-degree half-angle cone. Only forced convection heat transfer between the particles and the gas is considered in the computer programs and, therefore, a further cooling by radiation must be considered separately.

The plume computer output was employed to obtain the particle velocities for the seven discrete-size groups considered. The correlation was obtained in the form:

$$V(\rho_i, \theta) = [a_i + b_i \exp(-c_i R/r_t)] (\cos \theta)^{n_i}, \quad 0 \leq \theta \leq \theta_{\rho_i} \quad (2-1)$$

where: R, θ = Polar coordinates in plume

r_t = Nozzle throat radius

θ_{ρ_i} = Particle limiting streamline inclination (Reference 1)

Table 4-1 in Section 4 shows the correlation coefficients. The a_i coefficients decrease with increasing particle diameter which indicates that the larger particles achieve a lower limiting velocity as $R/r_t \gg 1$. In addition, the n_i coefficients were all negative which indicates that the particle velocities increase off-axis on a given spherical cap. This effect is due to the higher particle concentration on-axis which has caused more interphase momentum transfer from the gas to the particles, i. e., on a per particle basis the gas is less about to "drag" the particle. The lower off-axis particle density allows the gas to "drag" the particles to higher velocities.

The particle density correlation was redone specifically for the off-axis dependency:

$$\frac{d\rho_i}{d\rho_o} = a_i \left(\frac{r_t}{R}\right)^{n_i} (\cos \theta)^{b_i}, \quad 0 \leq \theta \leq \theta_{\rho_i} \quad (2-2)$$

The correlation coefficients are given in subsection 4.5.1.4.

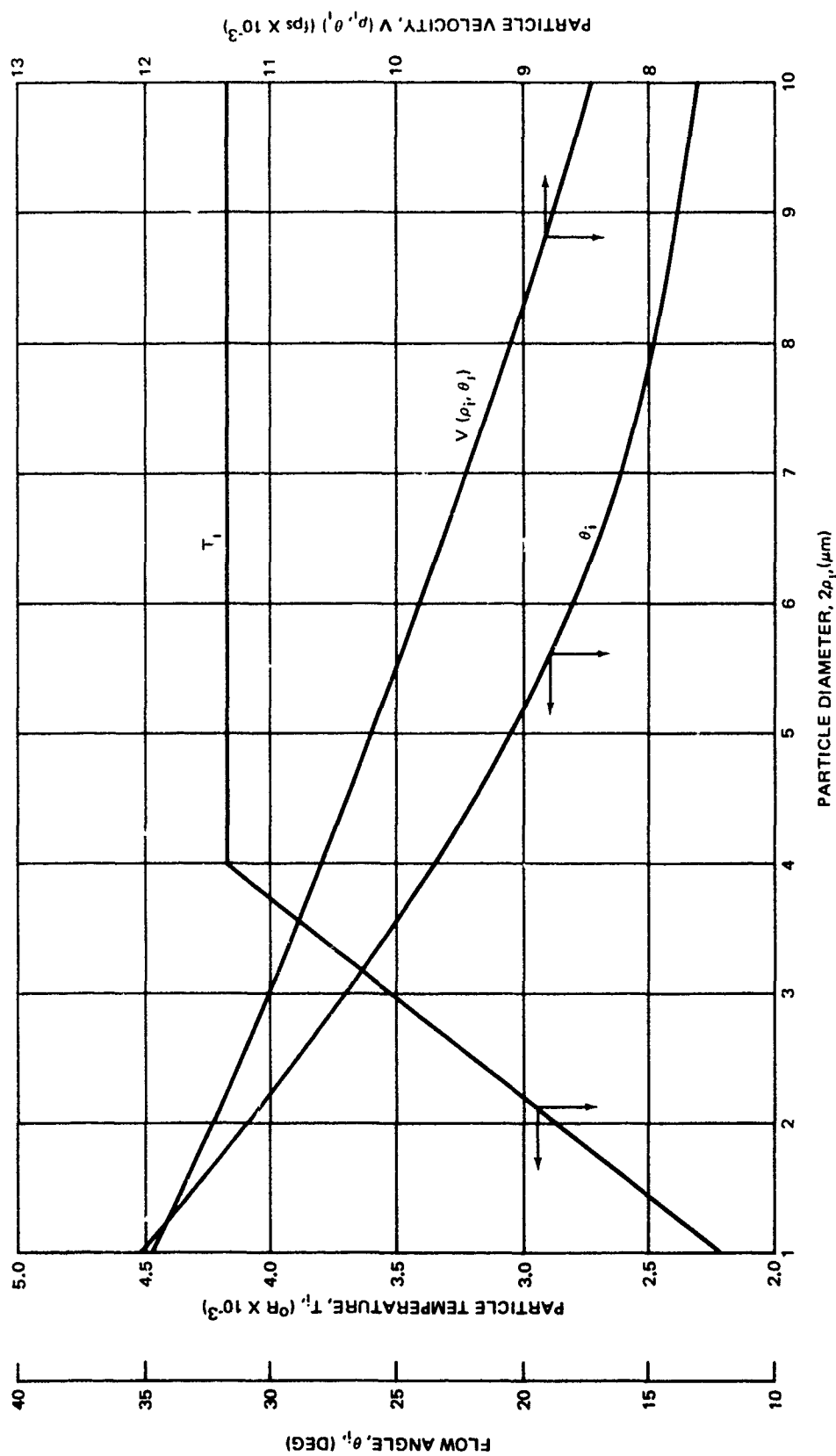


Figure 2-3. Particle Temperature, Flow Angle, and Velocity on Limiting Streamline

Aside for the largest particle size considered, the correlation showed that $b_i > 0$, i. e., the particle density decreased off-axis.

The two-phase computer results were analyzed to obtain a correlation of the gas phase mass flux on the engine centerline with the result being:

$$\left(\bar{v} \frac{d}{d_c} \right)_{\theta=0} = a \left(\frac{R}{r_t} \right)^{-2.7443} \quad (2-3)$$

where d_c = chamber density

The above result indicates that the mass flux is decaying more rapidly than inverse square with distance for the conditions employed. Additional computer runs should be performed for lower value of the solids loading in the exhaust products to determine if the gas phase mass flux will approach a $\left(\frac{R}{r_t} \right)^{-2}$ dependency.

Following Reference 5, the off-axis gas phase mass flux was assumed to be:

$$\left(\bar{v} \frac{d}{d_c} \right)_{\theta} = \exp [-\lambda^2 (1 - \cos c \theta)^2] \quad (2-4)$$

$$\left(\bar{v} \frac{d}{d_c} \right)_{\theta=0} = 0$$

where $c = \frac{90^\circ}{\theta_e + (\nu_{\max} - \nu_e)}$

and θ_e = Nozzle lip angle at exit

ν_{\max} = Maximum Prandtl-Meyer angle

ν_e = Prandtl-Meyer angle at exit Mach number

An expression was derived which relates the particle Knudsen number, Kn_{p_i} , to engine operating conditions, location, and particle size for a two-phase plume. This expression was obtained by defining Kn_{p_i} as the ratio of the relative Mach number and Reynolds number, i. e., :

$$Kn_{pi} = \frac{Mr_i}{Re_i} = \frac{|V - V_{pi}| / \sqrt{g_c \gamma R T}}{d |V - V_{pi}| 2 \rho_i / \mu} \quad (2-5)$$

where:

- V = Gas velocity
- V_{pi} = Velocity of the i th particle
- g_c = Conversion factor ($= 32.2 \text{ [lb}_m/\text{lb}_f]$). (ft/sec^2)
- γ = Gas ratio of specific heats
- R = Gas constant
- T = Gas static temperature
- d = Gas density
- ρ_i = Radius of the i th particle
- μ = Gas viscosity

Roberts' solution (Reference 1) was employed to relate d to location in the plume and to certain engine operating parameters; the viscosity was assumed to vary as \sqrt{T} ; and the isentropic relation was used to eliminate T . With these substitutions Equation 1 becomes:

$$Kn_{pi} = \left[\frac{12 \times 46.6 \times 10^{10} \sqrt{R_o/g_c} T_c}{P_c \gamma^{3/2} (\gamma-1) M_e^2 \left(1 + \frac{\gamma-1}{2} M_e^2 \right)^{-\frac{1}{\gamma-1}}} \right] \frac{(r/r_e)^2}{\rho_i \cos^{k-2} \theta}$$

where:

- r_e = Nozzle exit radius
- R_o = Universal gas constant ($= 1545 \text{ ft. lb}_f/\text{mole } ^\circ\text{R}$)
- T_c = Combustion temperature
- P_c = Chamber pressure
- M_e = Nozzle exit mach number
- r, θ = Polar coordinates of point in plume
- k = Hypersonic similarity parameter ($= \gamma(\gamma-1)m^2$)

The above equation indicates that, for given engine operating conditions and plume coordinates, the smaller particles will achieve the higher Knudsen number, i. e., they will "transition" sooner. Numerical calculations for a 1μ radius particle and typical engine operating conditions indicates that Kn_{pi} reaches unity very close to the nozzle exit plane. By comparison a gas-phase plume achieves a unity Knudsen number only at large distances from the exit plane.

2.2.2 First Estimate of Emissivity of Alumina Particles

The Mie theory for scattering of electromagnetic waves by dielectric spheres forms the basis for theoretical determination of particle emissivity. The theory actually computes the scattering and absorption efficiency factors, $Q_s(\alpha)$ and $Q_a(\alpha)$, defined by

$$Q = \sigma / \sigma_g$$

where σ_g is the geometrical cross section and σ the actual cross section. α is the size parameter $2\pi a / \lambda$, where a is the particle radius. The values obtained for Q are determined by the values of the real and imaginary parts of the index of refraction n , related by

$$n = n_1 - in_2$$

Since $n = n(\lambda, T)$, a more accurate functional dependence of Q would be $Q = Q(\lambda, a, T)$.

Plass (References 6 and 7) has computed efficiency factors for alumina spheres of 0.1 to 10μ for a range of wavelengths 0.5 to 10μ . Values of n_1 and n_2 used in this wavelength range were based upon measurements made with large sapphire crystals. They indicate $n_2/n_1 \ll 1$.

For determination of the emissivity, ϵ , Plass used an equation originally derived for the emissivity of a semi-infinite homogenous slab, with

$$\epsilon = 2.3 \sqrt{\sigma_a / \sigma_s}$$

Thus the following values of ϵ are obtained from the following table:

λ	r	ϵ
2μ	1.1 μ	0.005
	5.1	0.006
	9.9	0.018
5μ	1.1	0.054
	5.1	0.089
	9.9	0.115

(Figures for $\lambda = 10\mu$ are not immediately available).

It is to be noted that in the range $10 < \lambda < 20\mu$, values of n_2 increase by more than an order of magnitude (Brannon and Goldstein, Reference 8) so ϵ would be expected to be larger there, aside from size effects.

The results of Plass have been questioned by Morizumi and Carpenter (Reference 9) on the grounds that measurements of plume emissivities of rocket exhaust indicate larger values for particle emissivities. This conclusion is rationalized on the grounds that absorbing properties of rocket alumina particles are higher than those of sapphire, due to the polycrystalline properties of the former. They conclude that emissivity of alumina particles in rocket exhausts lies between 0.1 and 0.3.

It is concluded that conflicting and insufficient data available at this time precludes a definitive answer for the emissivity in the infrared. For provisional values, the following are suggested, using linear interpolation and extrapolation to obtain values at other sizes or wavelengths.

λ	r	ϵ	λ	r	ϵ	λ	r	ϵ
5 μ	1 μ	0.05	10 μ	1 μ	0.1	20 μ	1 μ	0.2
	5	0.09		5 μ	0.15		5	0.25
	10	0.12		10 μ	0.20		10	0.30

An approximate analytic correlation of the above data for Al_2O_3 emissivity was found. The correlation is shown in Figure 2-4, and is adequate, considering the coarseness of the available data. The emissivity correlation is of the form

$$\epsilon_{\lambda}(r) = (\lambda + 2.5)(.00689 + 11.32\rho - 2220\rho^2) \quad (2-6)$$

for λ in microns and ρ (particle radius) in cm.

2.3 ATMOSPHERIC EXCITATION

Atmospheric excitation of rocket-plume species by collisional high relative velocities has been identified as a strong source of plume radiation. A detailed analysis of the problem was presented in Reference 10 and summarized in Reference 11. The unique feature of this analysis is the inclusion of the time-dependent radiative decay in the treatment of the excited species. Thus, there is a finite time required before the plume can generate radiation by this mechanism.

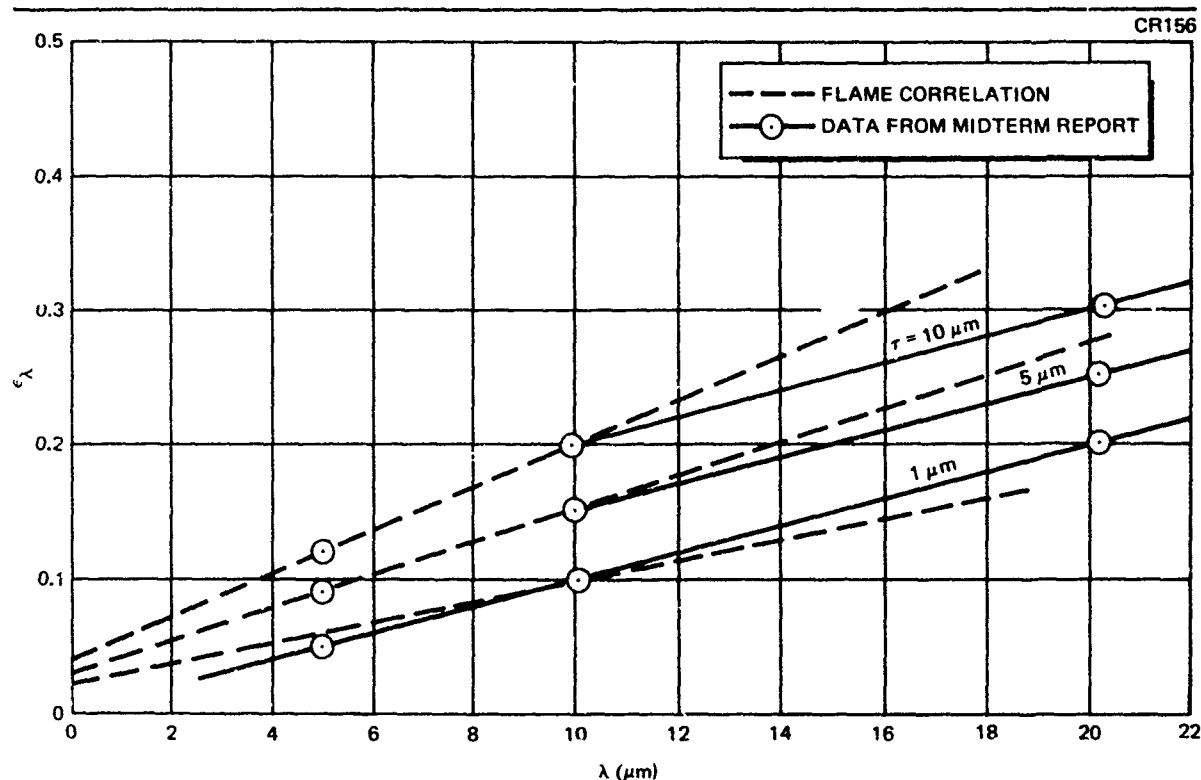


Figure 2-4. Alumina Particle Emissivity

Detailed calculations of trajectories of the colliding species, averaged over initial orientations, were done using Program P2160. Figure 2-5 shows the resulting rotational temperatures induced as a function of impact parameter and V_∞ . Also shown on Figure 2-5 is the linear variation found for T_m/V_∞ as a function of V_∞ (T_m is the maximum rotational temperature for any impact parameter).

The effective collision cross section (S) to use with T_m is found by requiring

$$T_m S = 2\pi \int_0^\infty T(b) b db$$

Calculated values of S are shown in Figure 2-6. For the current version of FLAME, S will be assumed to be constant ($=0.7 \times 10^{-5} \text{ cm}^2$).

The atmospheric interaction model used in FLAME follows that of Reference 10 namely:

- A. For the wavelength band of interest, the total energy (E_T) and radiance (j) per molecule are calculated as a function of T .
- B. Radiative decay times are calculated as a function of T_m .

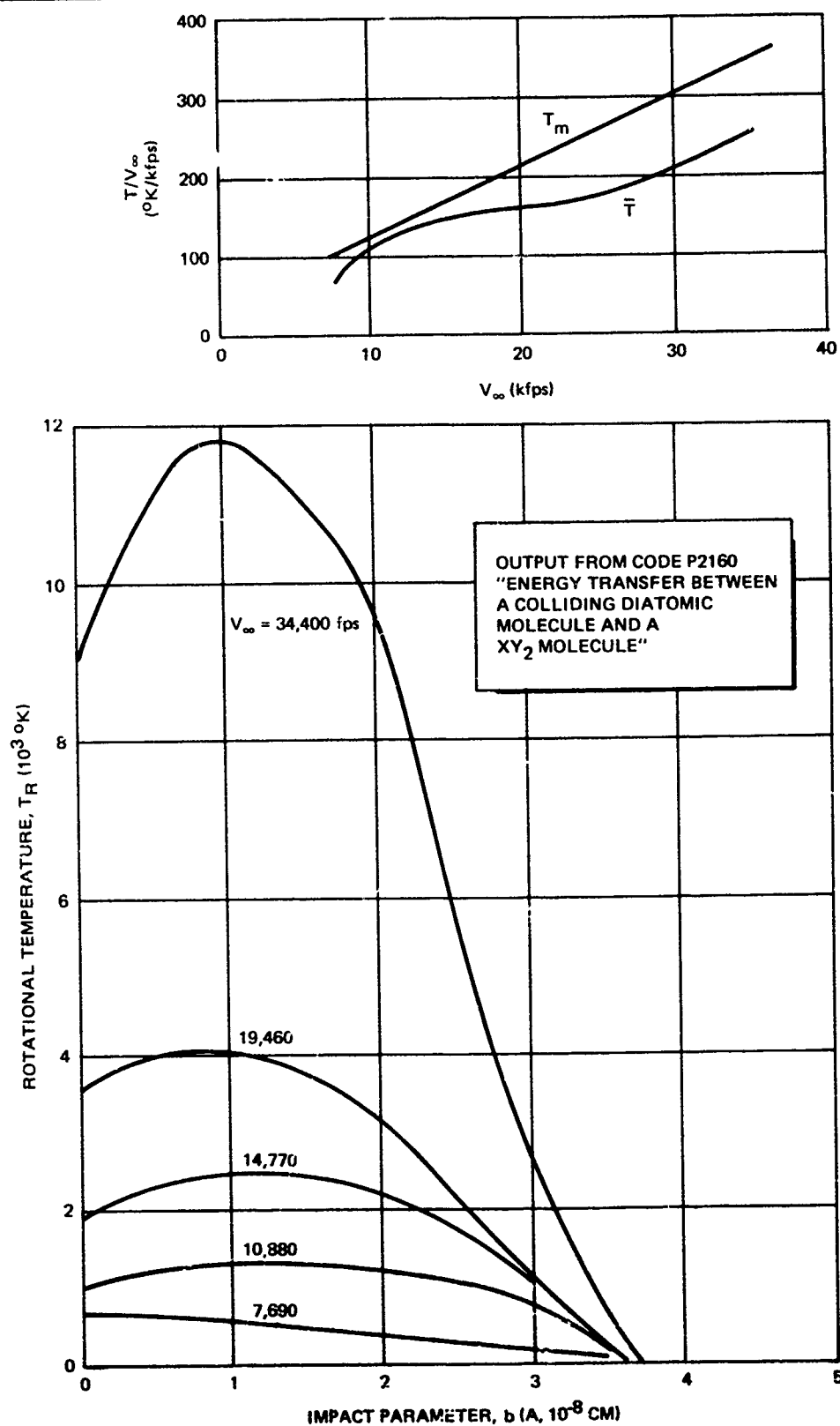


Figure 2-5. Water Rotational Energy from Single Collision

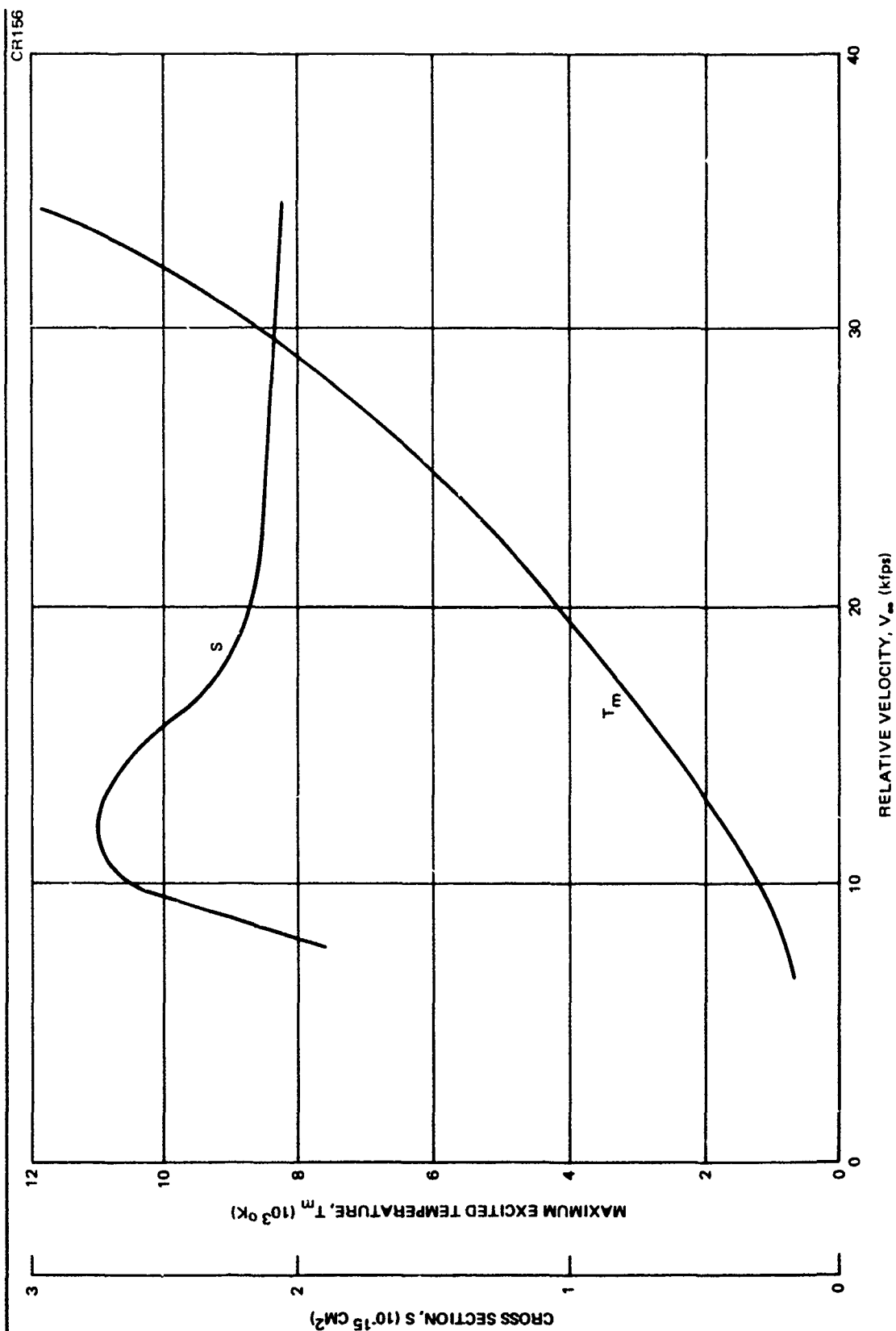


Figure 2-6. Rotational Excitation Parameters for H₂O

- C. The time of flight to a position in the plume is compared to the decay time so that an average radiance per excited molecule can be found.
- D. The fraction of molecules that are excited are calculated assuming initially cold plume species.
- E. No mechanism currently is programmed to account for plume/atmosphere mixing, but plans are to add this later. Currently if a molecule undergoes multiple collisions in the atmosphere, the resultant radiance is found by simple addition.
- F. The radiance decay of an excited species is assumed to be linear in time.
- G. Data currently in program is applicable to a single species, H_2O .

Further details of the analysis are available in Reference 10.

The analysis of Reference 10 was changed in one way. Since the radiative model used assumes single collisions, and the plume may be viewed from long distance, an attenuation factor (A_t) must be introduced to account for depletion of the supply of uncollided molecules as the gas expands far from the nozzle. Assuming each collision removes that molecule from further consideration, the reduction of the number of available molecules from N_0 (initial value) to N (at point r in plume) is given by

$$\frac{N}{N_0} = \exp(-A_t) \quad (2-7)$$

where

$$A_t = S \int_0^r n_\infty ds,$$

where S is the collision cross section and n_∞ is the atmospheric density.

A correlation was obtained for the spectral linear absorption coefficient of the H_2O rotational spectrum. The results are that

$$k_\lambda = \log^{-1}(y) \text{ in cm}^{-1} \quad (2-8)$$

where

$$y = A \frac{(\omega - \omega_0)^2}{\omega} + y_0$$

$$\omega_0 = 10.6 \sqrt{T}, \text{ T in } ^\circ K$$

$$y_o = \frac{279.96}{\sqrt{T} + 50.031} - 3.208$$

$$A = \left[1.030 - \frac{105.582}{\sqrt{T} - 2.896} \right] \times 10^{-4}$$

$$\omega = 10,000/\lambda, \lambda \text{ in } \mu\text{m}$$

The data correlated are shown in Figure 2-7, where the results of the correlation are shown for temperatures of 600°K and 1,800°K.

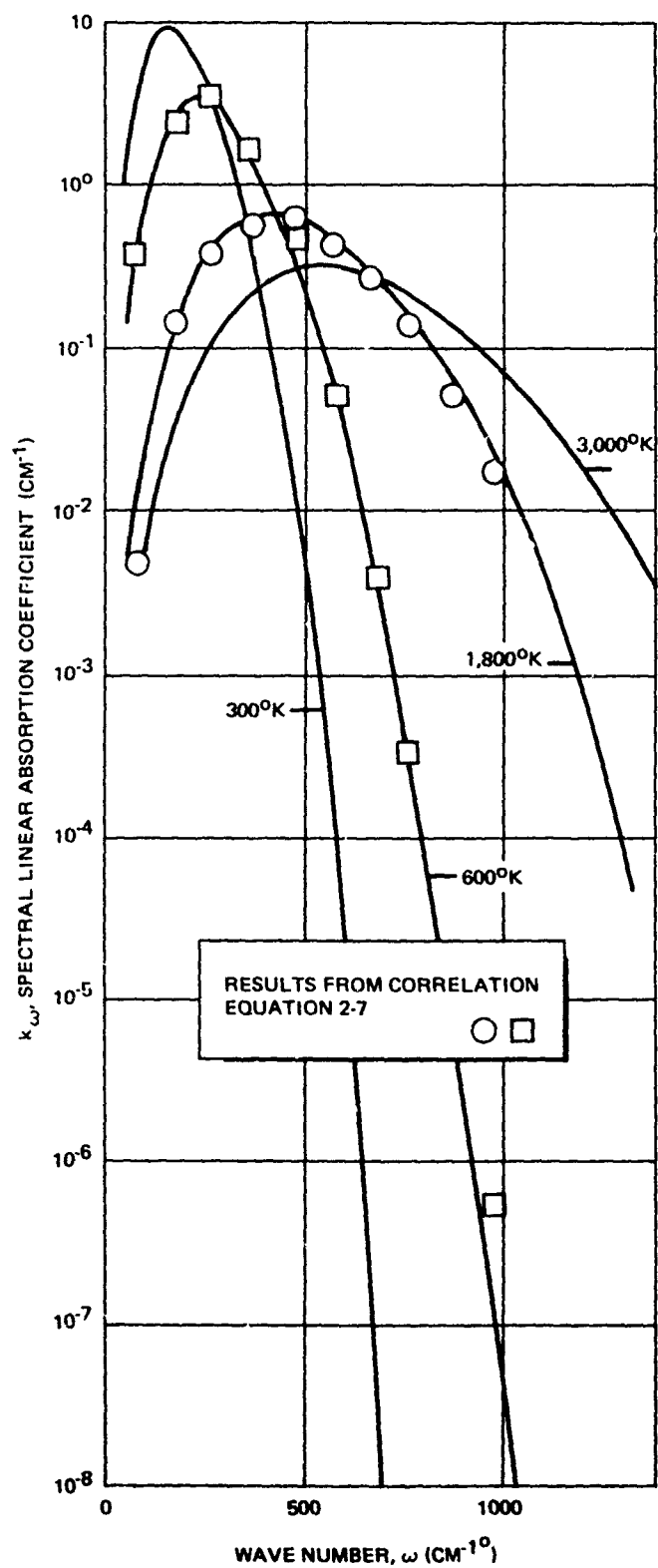


Figure 2-7. Absorption Coefficient from Water Rotation at One Atmosphere Pressure

Section 3 FLAME CODE

The following subsections deal with the logic underlying the FLAME code as it is now configured. Subsection 3.1 describes the logic that went into the choice of coordinate systems, and displays the geometric model used. Subsection 3.2 presents the flow charts used as the basis for FLAME.

3.1 GEOMETRY CONSIDERATIONS

The rocket plume is an axisymmetric body about the z -axis, as illustrated in Figure 3-1. Integration through the plume is carried out through the conical volume in the first and second quadrants defined by the polar coordinates θ_m and r_m . The viewing angle is specified completely by the angle α with respect to the x -axis in the xz plane.

We want to integrate emittance along families of lines of sight, such as s ,

$$\int E(r, \theta) ds$$

We can express ds in terms of either cartesian or spherical coordinates,

$$ds = dx \sqrt{1 + \left(\frac{dy}{dx}\right)^2 + \left(\frac{dz}{dx}\right)^2} \quad (3-1)$$

or

$$ds = dr \sqrt{1 + r^2 \left(\frac{d\phi}{dr}\right)^2 + r^2 \sin^2 \phi \left(\frac{d\theta}{dr}\right)^2} \quad (3-2)$$

Since emittance is a function of r, θ along, equation (3-2) is more inviting to use, if it is practical. The following analysis shows, however, that cartesian coordinates will probably result in less computation time, and should be used.

The conical segment in Figure 3-1 has equation

$$x^2 + y^2 = z^2 \tan^2 \theta_1 \quad (3-3)$$

The family of lines along direction of observation, α , have equations

$$y = y_1 \quad (3-4)$$

$$x = x_1 + s \cos \alpha \quad (3-5)$$

$$z = z_1 + s \sin \alpha \quad (3-6)$$

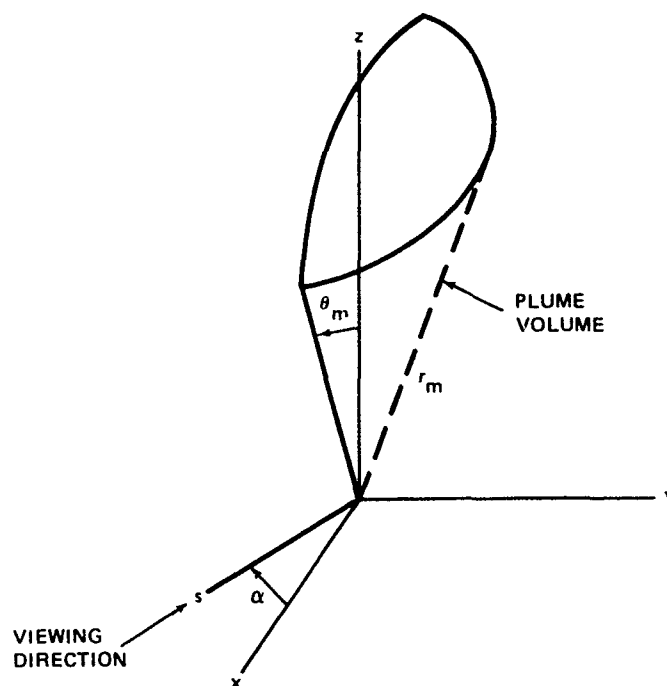


Figure 3-1. Plume Orientation and Viewing Direction

These equations can be switched to spherical coordinates using the standard relations

$$x = r \sin \theta \cos \phi \quad (3-7)$$

$$y = r \sin \theta \sin \phi \quad (3-8)$$

$$z = r \cos \theta \quad (3-9)$$

The following discussion derives the spherical and cartesian forms for ds along straight lines of sight:

3.1.1 Spherical Coordinates

$$\text{From Equation 3-8 } \frac{y_1^2}{(r \sin \theta)^2} = \sin^2 \phi; \quad (3-10)$$

From Equations 3-5, 3-6, 3-7, and 3-9

$$\frac{[\cot \alpha (r \cos \theta - z_1) + x_1]^2}{(r \sin \theta)^2} = \cos^2 \phi \quad (3-11)$$

Adding Equations 3-10 and 3-11

$$y_1^2 + \left[\cot \alpha (r \cos \theta - z_1) + x_1 \right]^2 = (r \sin \theta)^2 = r^2 - (r \cos \theta)^2 \quad (3-12)$$

let

$$\cot \alpha \ r \cos \theta = F$$

$$-\cot \alpha \ z_1 + x_1 = A$$

$$-y_1^2 + r^2 = B$$

Equation 3-12 becomes

$$F^2 + 2AF + A^2 - B + \frac{F^2}{\cot^2 \alpha} = 0$$

or

$$F^2(1 + \tan^2 \alpha) + 2AF + A^2 - B = 0; \quad F^2 / \cos^2 \alpha + 2AF + A^2 - B = 0$$

Therefore

$$F = \frac{-2A \pm \sqrt{4A^2 - 4A^2 \left(\frac{1}{\cos^2 \alpha} \right) + \frac{4B}{\cos^2 \alpha}}}{1/\cos^2 \alpha} =$$

$$-2A \cos^2 \alpha \pm 2 \cos \alpha \sqrt{B - A^2 \sin^2 \alpha}$$

Therefore

$$\sin \alpha \ r \cos \theta = -2A \cos \alpha \pm 2 \sqrt{B - A^2 \sin^2 \alpha}$$

or

$$\cos \theta = -\frac{2A \cot \alpha}{r} \pm \frac{2}{r} \sqrt{\frac{B}{\sin^2 \alpha} - A^2} \equiv u(r)$$

$$\theta = \cos^{-1} \left(-\frac{2A \cot \alpha}{r} \pm \frac{2}{r} \sqrt{\frac{B}{\sin^2 \alpha} - A^2} \right)$$

From consideration of the case $\alpha = 90$ degrees and noting that $\cos\theta > 0$, only + root is valid.

$$\frac{d\theta}{dr} = \frac{-1}{\sqrt{1-u^2}} \frac{du}{dr}$$

or

$$\frac{d\theta}{dr} = \frac{-1}{\sqrt{1-u^2}} \left[-\frac{u}{r} + \frac{1}{r} \left(\frac{B}{\sin^2 \alpha} - A^2 \right)^{-1/2} \frac{2r}{\sin^2 \alpha} \right] \quad (3-13)$$

also

$$\sin \phi = \frac{y_1}{r \sin \theta} = \frac{y_1}{r \sqrt{1-u^2}} \equiv v(r)$$

$$\frac{d\phi}{dr} = \frac{1}{\sqrt{1-v^2}} \frac{dv}{dr}$$

where

$$\frac{dv}{dr} = \frac{-y_1}{r^2 \sqrt{1-u^2}} + \frac{y_1 u}{r(1-u^2)^{3/2}} \frac{du}{dr}$$

From Equation 3-13

$$\frac{d\theta}{dr} = \frac{uw}{r(1-u^2)^{1/2}}$$

where

$$w(r) = \left[1 - \frac{2r}{u \sin^2 \alpha \sqrt{B/\sin^2 \alpha - A^2}} \right]$$

then

$$\frac{dv}{dr} = \frac{-y_1}{r^2 \sqrt{1-u^2}} - \frac{y_1 u^2 w}{r^2 (1-u^2)^{3/2}}$$

$$\frac{du}{dr} = - \frac{uw}{r}$$

or finally

$$ds = dr \left[1 + r^2 \frac{1}{1-v^2} \frac{y_1^2}{r^4(1-u^2)} \left(1 + \frac{u^2 w^2}{1-u^2} \right)^2 + r^2 \frac{y_1^2}{r^2(1-u^2)} \frac{u^2}{r^2} w^2 \right]^{1/2} \quad (3-14)$$

From Equations 3-3, 3-4, 3-5, and 3-6 the limits of integration, r , and r_2 can be found.

We have

$$x^2 + y^2 + z^2 = r^2 \quad (3-15)$$

$$x^2 + y^2 = z^2 \tan^2 \theta_m = r^2 - z^2$$

$$r^2 = z^2 (1 + \tan^2 \theta_m)$$

Therefore

$$z = r \cos \theta_m$$

Also

$$y = y_1$$

and

$$\frac{x-x_1}{\cos \alpha} = \frac{z-z_1}{\sin \alpha}$$

or

$$x = \cot \alpha (z-z_1) + x_1 \quad (3-16)$$

Substituting x , y , and z into equation of the cone, Equation 3-3 we have a quadratic equation for r , namely

$$(\cot^2 \alpha \cos^2 \theta_1 + \cos^2 \theta_1 - 1)r^2 + (2x_1 \cot \alpha \cos \theta_1 - 2\cot^2 \alpha z_1 \cos \theta_1)r + (\cot^2 \alpha z_1^2 - 2x_1 \cot \alpha z_1 + x_1^2 + y_1^2) = 0 \quad (3-17)$$

The solution gives r_1 and r_2 , as function of z_1 , x_1 , and y_1

3.1.2 Cartesian Coordinates

From Equation 3-16

$$z = \tan \alpha (x - x_1) + z_1 \quad (3-18)$$

or

$$\frac{dz}{dx} = \tan \alpha$$

Therefore

$$ds = dx \sqrt{1 + \tan^2 \alpha} = \frac{dx}{\cos \alpha} \quad (3-19)$$

or

$$ds = \frac{dz}{\sin \alpha} \quad (3-20)$$

We obtain r and θ from Equation 3-18

$$r^2 = z^2 + y_1^2 + [\cot \alpha (z - z_1) + x_1]^2 \quad (3-21)$$

$$\cos \theta = z/r \quad (3-22)$$

It is apparent that it is easier to use cartesian coordinates, and Equations 3-19 or 3-20 together with Equations 3-18, 3-21, and 3-22 to integrate through the plume, than to use the single, complex Equation 3-14.

To avoid singularities at $\alpha = 0$ degree or 90 degrees, x will be used as the independent variable for $0 \leq \alpha \leq 45$ degrees and z for the independent variable for $45 \text{ degrees} < \alpha \leq 90 \text{ degrees}$. Since the plume is transparent, negative values of α are not needed, the integrals being equal to values for 90 degrees $+\alpha$.

For cartesian coordinates, the intercept equations are more complex than for spherical coordinates; substituting

$$y = y_1 \quad (3-23)$$

$$x = \cot \alpha (z - z_1) + x_1 \quad (3-24)$$

into equation of spherical cap

$$x^2 + y^2 + z^2 = r_m^2 \quad (3-25)$$

and conical sides

$$x^2 + y^2 = z^2 \tan^2 \theta_m \quad (3-26)$$

yields the following equations for the resulting four intercept points of a straight line of sight:

For $0 < \alpha < 45$ degrees (choose $x_1 = 0$)

$$(1 + \tan^2 \alpha) x^2 + (2z_1 \tan \alpha) x + y_1^2 + z_1^2 - r_m^2 = 0 \quad (3-27)$$

$$(\cot^2 \theta_m - \tan^2 \alpha) x^2 - (2z_1 \tan \alpha) x + y_1^2 \cot^2 \theta_m - z_1^2 = 0 \quad (3-28)$$

and for $45 \text{ degrees} < \alpha \leq 90 \text{ degrees}$ (choose $z_1 = 0$)

$$(1 - \cot^2 \alpha) z^2 + (2x_1 \cot \alpha) z + y_1^2 + x_1^2 - r_m^2 = 0 \quad (3-29)$$

$$(\tan^2 \theta_m - \cot^2 \alpha) z^2 - (2x_1 \cot \alpha) z - (y_1^2 + x_1^2) = 0 \quad (3-30)$$

For each range of α , solution of the two quadratic equations will yield four values for x or z . The correct limits of integration are found by selecting those pairs of x or z for which the corresponding values of r and θ (by Equations 3-18, 3-21, and 3-22) are on the real plume boundary, namely

$$r_m \text{ and } \theta \leq \theta_m$$

or

$$\theta_m \text{ and } r \leq r_m$$

3.1.3 Integration Schemes and Step Sizes

Most of the integrals of the form

$$I = \int_{x_1}^{x_2} f(x) dx \quad (3-31)$$

are evaluated using Gaussian eight-point quadrature. Defining

$$y = (2x - x_1 - x_2)/(x_2 - x_1) \quad (3-32)$$

we have

$$I = \frac{x_2 - x_1}{2} \int_{-1}^1 f(x) dy \quad (3-33)$$

Then

$$I \approx \frac{x_2 - x_1}{2} \sum_{i=1}^8 w_i f(x_i) \quad (3-34)$$

where

$$y_1 = y_8 = 0.18343 \ 46424 \ 95650$$

$$y_2 = y_7 = 0.52553 \ 24099 \ 16329$$

$$y_3 = y_6 = 0.79666 \ 64774 \ 13627$$

$$y_4 = y_5 = 0.96028 \ 98564 \ 97536$$

and

$$w_1 = w_8 = 0.36268\ 37833\ 78362$$

$$w_2 = w_7 = 0.31370\ 66458\ 77887$$

$$w_3 = w_6 = 0.22238\ 10344\ 53374$$

$$w_4 = w_5 = 0.10122\ 85362\ 90376$$

The values of x_i are related to y_i through Equation 3-32.

For integrating through the plume, Simpson's rule is used

$$\int_x^{x+2\Delta x} f(y)dy = (f(x) + 4f(x+\Delta x) + f(x+2\Delta x)) \frac{2\Delta x}{3} \quad (3-35)$$

The step size, Δx , is varied after each successive integration from x to $x + 2\Delta x$. A geometric variation of Δx assures that the smallest step size occurs as the integration along the line of sight passes through the plume center ($x = 0$ or $z = 0$ planes).

Let the total maximum distance along the line of sight be x_T . Starting at the middle of the plume and integrating out, a distance $x_T/2$ is covered in about $n/2$ steps (n is a program input). If each Δx step is r times bigger than the previous, the rule for geometric progression gives

$$\frac{x_T}{2} = \Delta x_{\min} \frac{r^{n/2} - 1}{r - 1} \quad (3-36)$$

where Δx_{\min} is the initial, smallest step size. Choosing a 10-percent increase in Δx each step, $r = 1.1$

or

$$\Delta x_{\min} = \frac{x_T}{20(1.1^{n/2} - 1)} \quad (3-37)$$

Starting at position $x = -|x_0|$, a large step size is used which decreases until $x = 0$ and then increases again. The initial value of Δx is found from the formula for the sum in terms of the last term (Δx_{in})

$$x_0 = \frac{\Delta x_{in} r - \Delta x_{min}}{r - 1} \quad (3-38)$$

or

$$\Delta x_{in} = (0.1 x_c + \Delta x_{min})/1.1 \quad (3-39)$$

A similar stepping scheme is used to obtain a higher density of lines of sight near the rocket exit plane or plume centerline than used in the rarified plume.

3.1.4 Relative Velocity

The rocket-body velocity must be specified in the coordinate system of Figure 3-1 so that relative velocity of plume gases and atmosphere can be calculated. Figure 3-2 shows the position of the vehicle velocity vector V_R , and the definition of straight up, \hat{u} .

Integration is performed only over a plume where $y > 0$. To obtain other half of plume, ψ_R and ψ_u should be increased 180° .

Relative velocity is just the vector sum of $V_p \hat{r}$ and \vec{V}_R where \hat{r} is the unit vector along the polar coordinate r to a point in the plume being considered. Altitude at this same point is missile altitude h_0 plus projection of \vec{r} onto \hat{u} .

$$h = h_0 + \vec{r} \cdot \hat{u} \quad (3-40)$$

3.2 FLOW CHARTS

This section presents the logic flow of the FLAME code. Input data are listed in Tables 3-1 through 3-4. The Figure 3-3 is a simplified diagram of the overall code. The charts (Figures 3-4 through 3-9) expand upon certain regions of the computer.

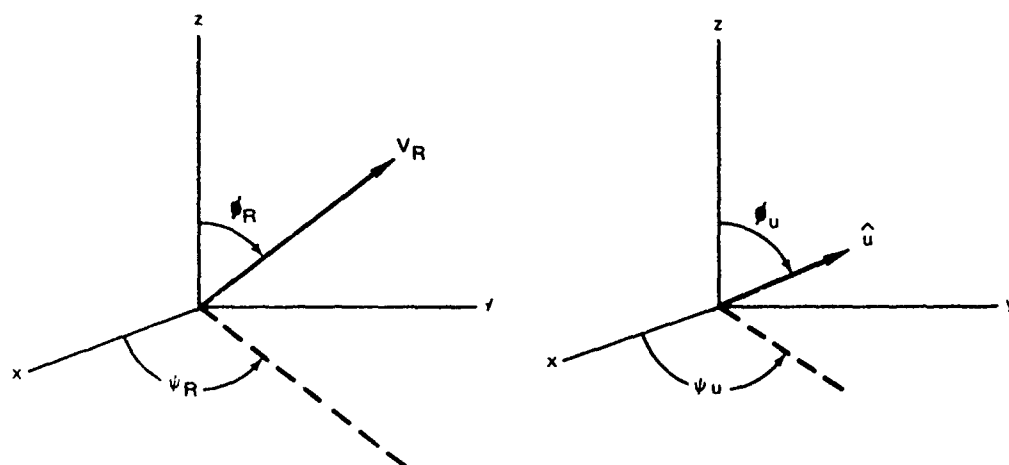


Figure 3-2. Location of Vehicle Velocity Vector and Upward Direction in Coordinate System of Figure 3-1.

Table 3-1

LISTING OF INPUT DATA

Read in		FORTTRAN	Real or Integer
t ,	time after ignition (sec)	TIG	R
r_e ,	nozzle exit radius (cm)	RE	R
γ ,	ratio of specific heats (>1)	GAMMA	R
M_e ,	exit mach number	ME	R
θ_e ,	nozzle lip angle (deg)	LIP	R
$d_e(t)$,	exit plane density as a table versus time after ignition (g/cm^3 vs sec)	PDE TID	R R
I_{ND} ,	night/day flag (0 or 1)	IND	I
I_{SL} ,	solid/liquid flag 0 or -2 (gas); 1 or -1 (solid)	ISL	I
λ_{\min} ,	band limits (μm)	WLMIN	R
λ_{\max} ,		WLMAX	R
n ,	approximate integration steps through layer	INT	I
m ,	approximate number of lines of sight used in one direction	LINES	I
T_c	chamber temperature ($^{\circ}\text{K}$)	TC	R
x_m ,	radiating gaseous species mole fraction*	XM	R
T_E ,	Earth effective temperature ($^{\circ}\text{K}$)	TE	R
α ,	look angle ($0 \leq \alpha \leq 90$ degrees)	AL	R
h_0 ,	altitude of engine (ft) (from 0.3 Mft to 3 Mft)	HO	R
V_R ,	rocket velocity (fps)	VR	R
ϕ_u	polar angle of vertical direction (deg)	PHIU	R
ψ_u	azimuthal angle of vertical direction (deg)	XIU	R
ϕ_R ,	polar angle of rocket-body velocity (deg)	PHIR	R
ψ_R ,	azimuthal angle of rocket-body velocity (deg)	XIR	R
END,	flag signifying end of run (0 or 1)	END	I

*FLAME assumes an average molecular weight (WMOL) of 17 for the plume species (approximately correct for pure gas plumes, or the gaseous constituents of solid-propellant plumes). When gaseous radiation from a solid propellant plume is calculated ($|I_{SL}| = 1$ and $x_m \neq 0$), the molecular weight should be corrected to include solids. This is done by modifying x_m , since only the ratio $XM/WMOL$ is used by FLAME. For a typical solid-propellant plume, $WMOL = 26$, so XM should be reduced by the ratio (17/26).

Table 3-2

PROGRAM REGION 1. INITIAL CALCULATIONS-PARTICLES

The following data will be calculated for each particle radius, p_i , corresponding to the i th particle size:

EPT,	table of total emissivity versus temperature
TCOOL,	table of particle temperature versus time after solidification
BPASS,	black body intensity in band of interest
ETL	table of average emissivity in region from λ_{\min} to λ_{\max} versus temperature
TMIN,	equilibrium temperature of particle
DELT,	time required to solidify droplet

Table 3-3

PROGRAM REGION 2. INITIAL CALCULATIONS-GAS

Calculate limiting plume-gas velocity, VP.
For the band width specified, calculate

ET,	table of total available rotational energy as function of temperature
JM,	table of radiative rate (radiance per molecule) as function of temperature
JT,	table of radiative rate as a function of time

Table 3-4

PROGRAM REGION 3. SELECT x_1 OR z_1 for $\alpha \geq 90 - \theta_m$

Test IA to determine magnitude of α . For IA = 0, define z_1 .
For IA = 1, define x_1 . Both x_1 and z_1 are labeled X1 in the program.

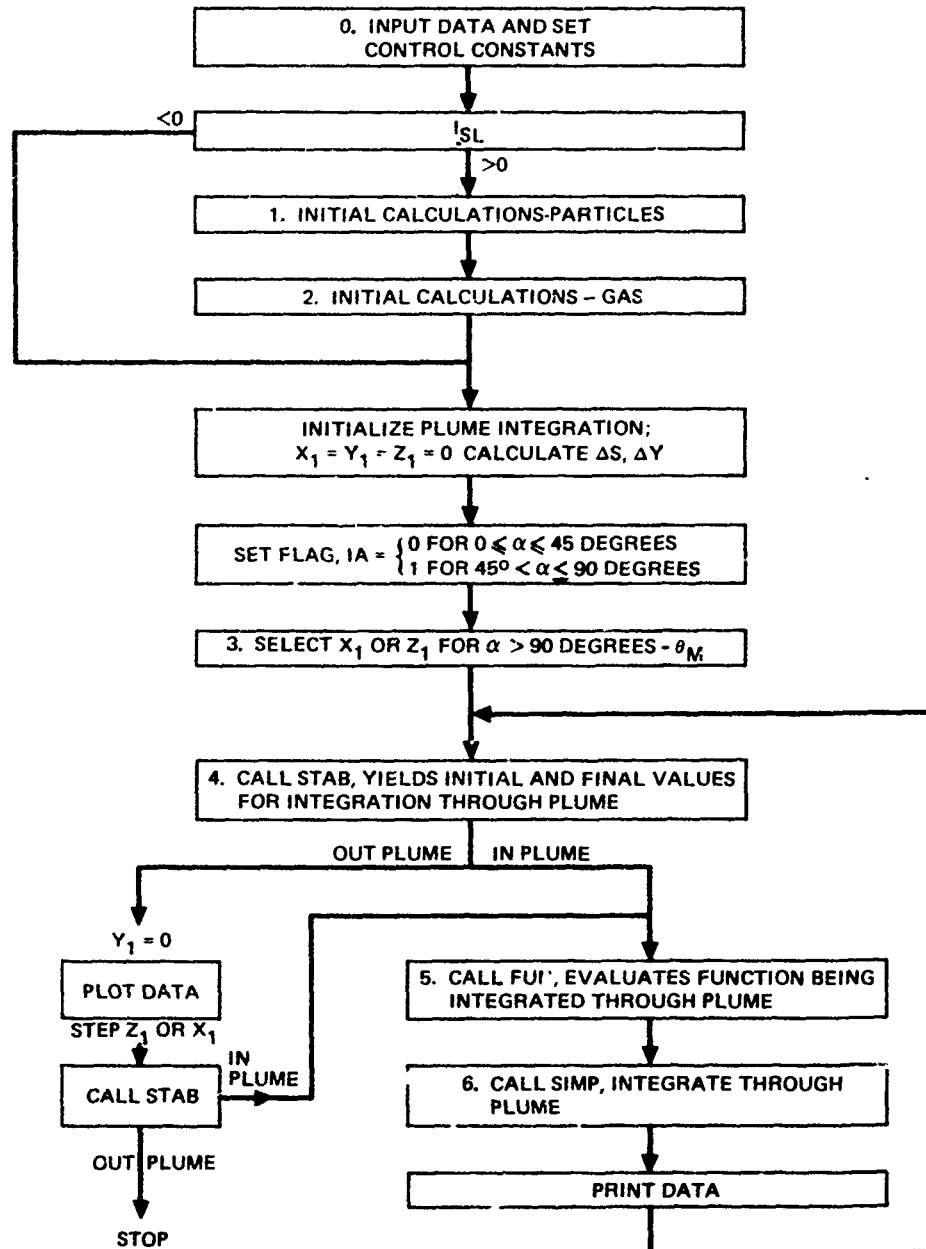


Figure 3-3. Master Flow for Program P2170, FLAME

CALL STAB AND CHECK (YIELDS INITIAL AND FINAL VALUES FOR
INTEGRATION THROUGH PLUME)

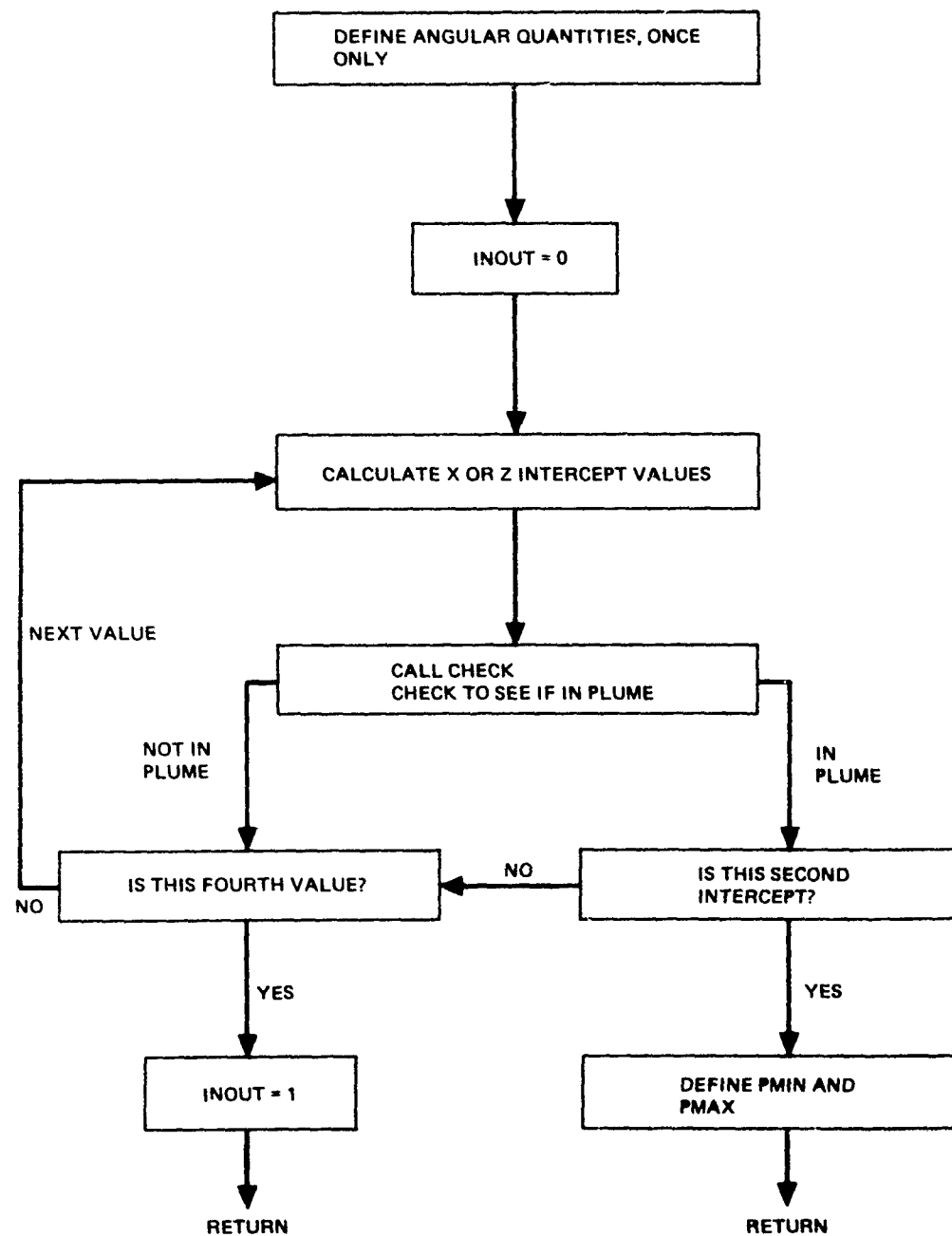


Figure 3-4. Program Region 4 Flow

CALL FUN (EVALUATES FUNCTION BEING INTEGRATED THROUGH PLUME)

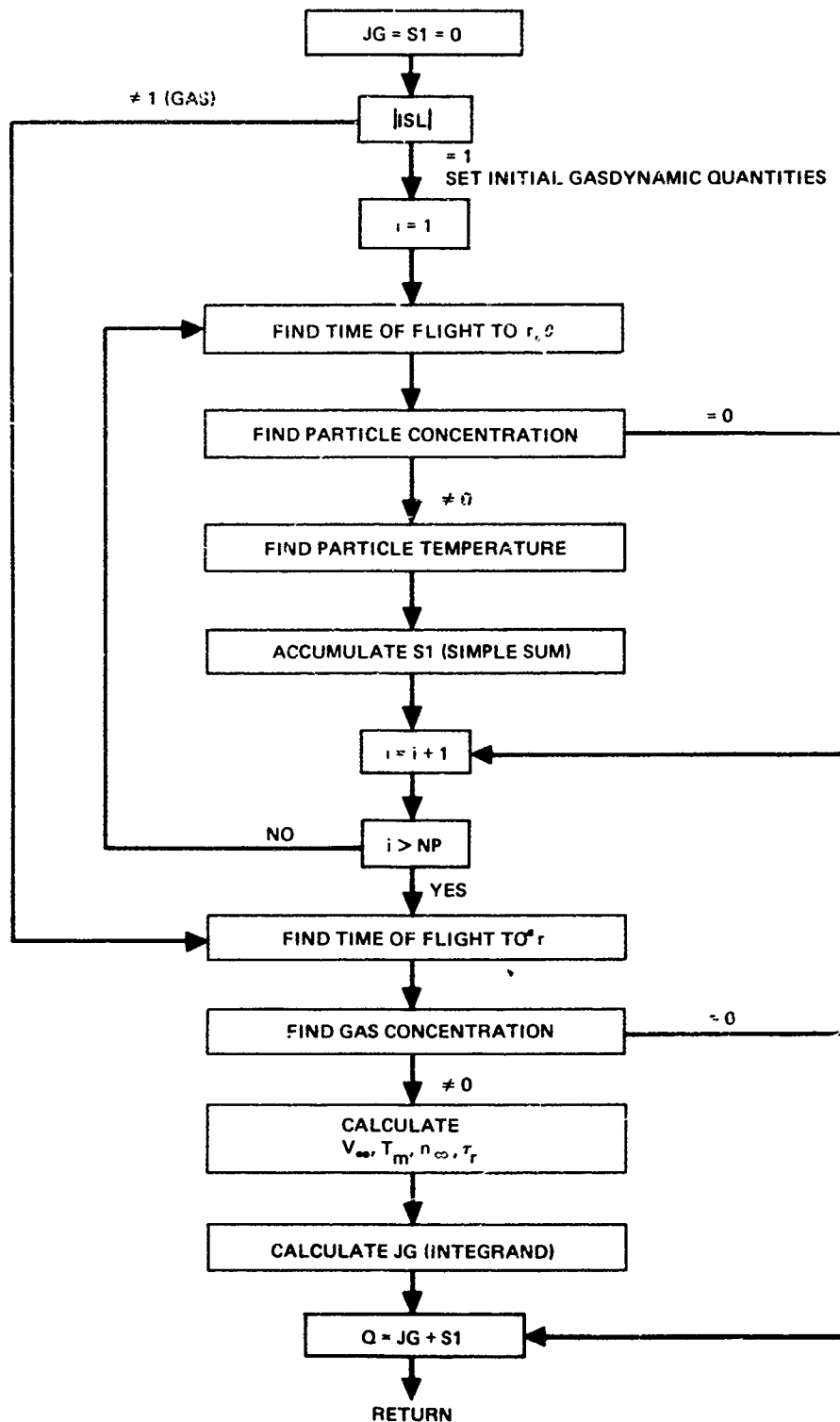


Figure 3-5. Program Region 5 Flow

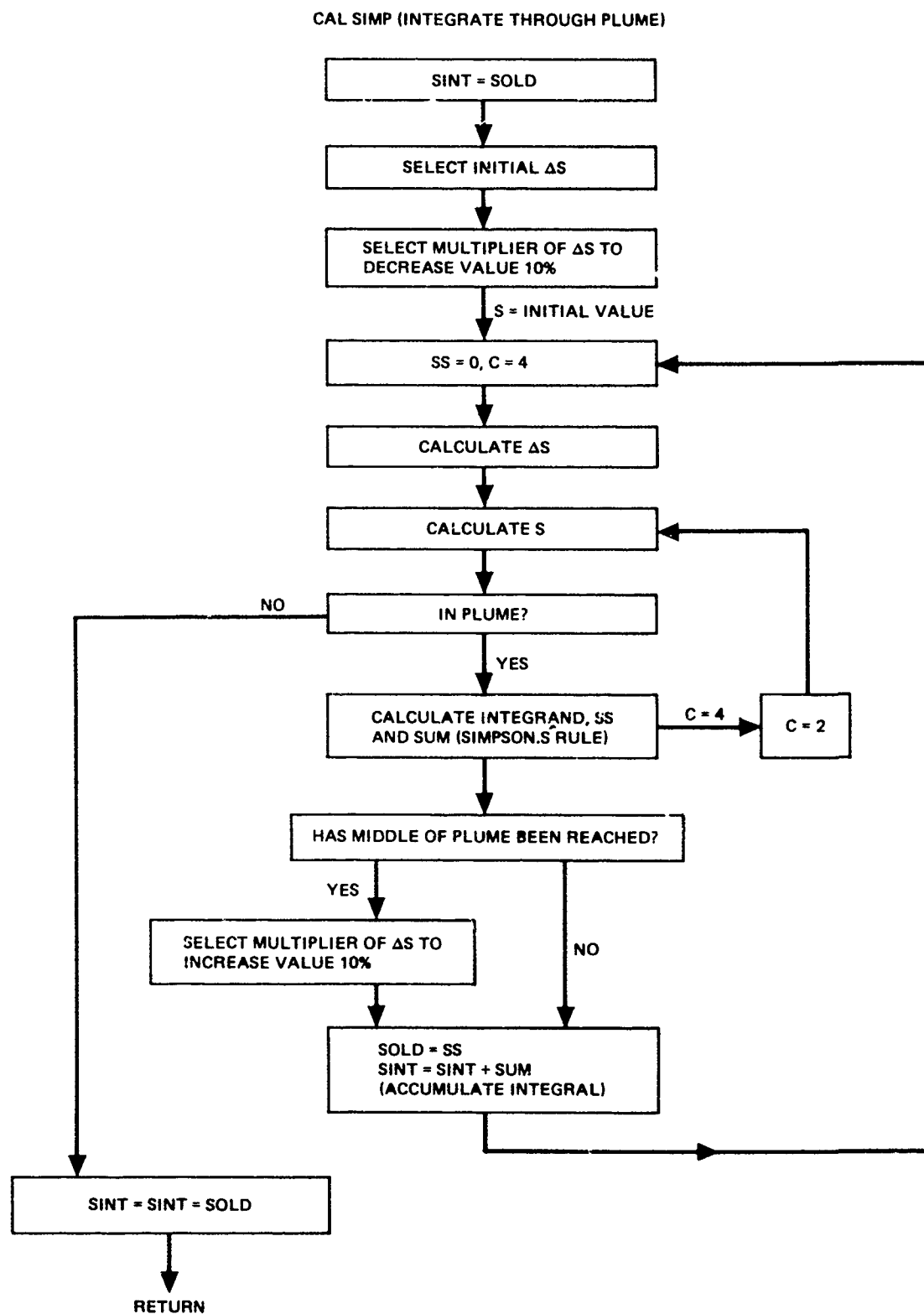


Figure 3-6. Program Region 6 Flow

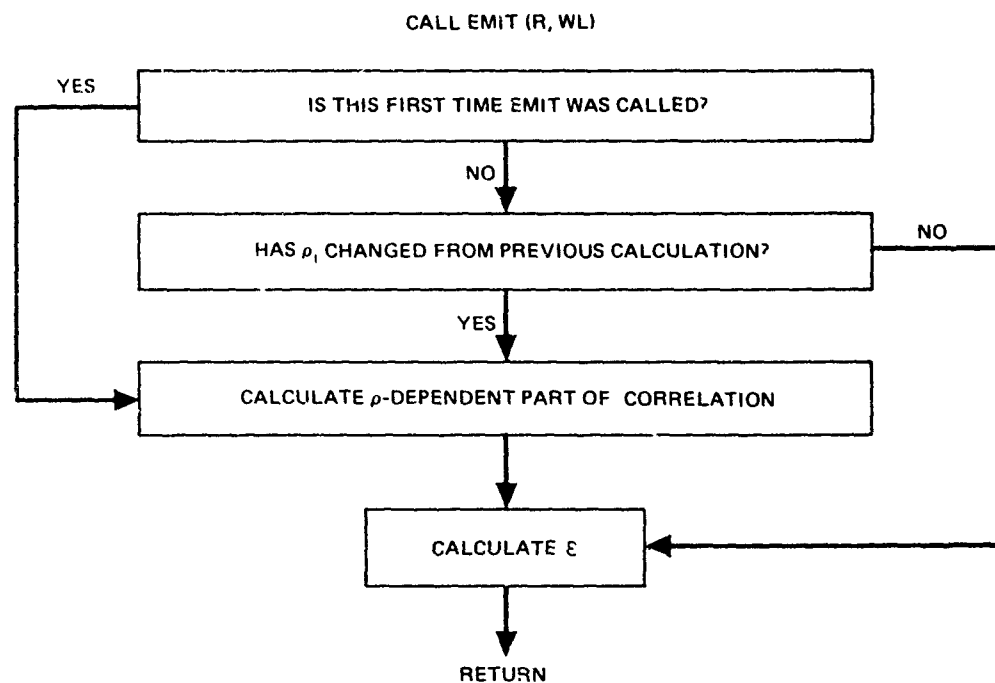


Figure 3-7. Particle Emissivity Routine Flow Chart

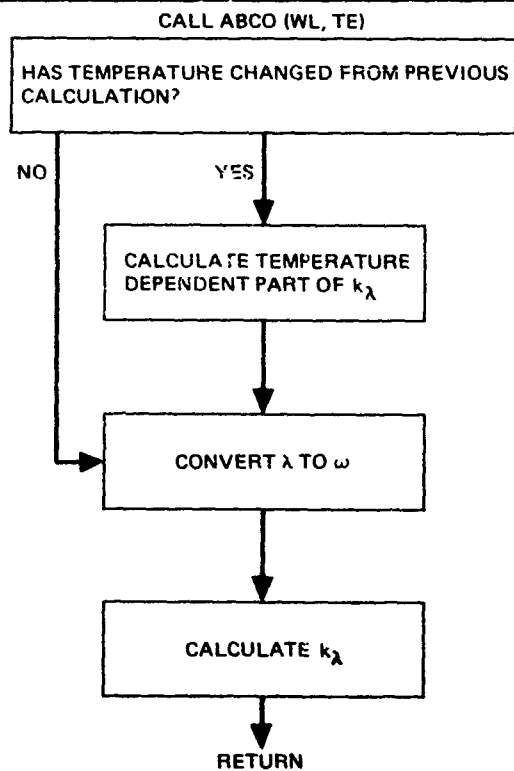


Figure 3-8. Absorption Coefficient Routine Flow Chart

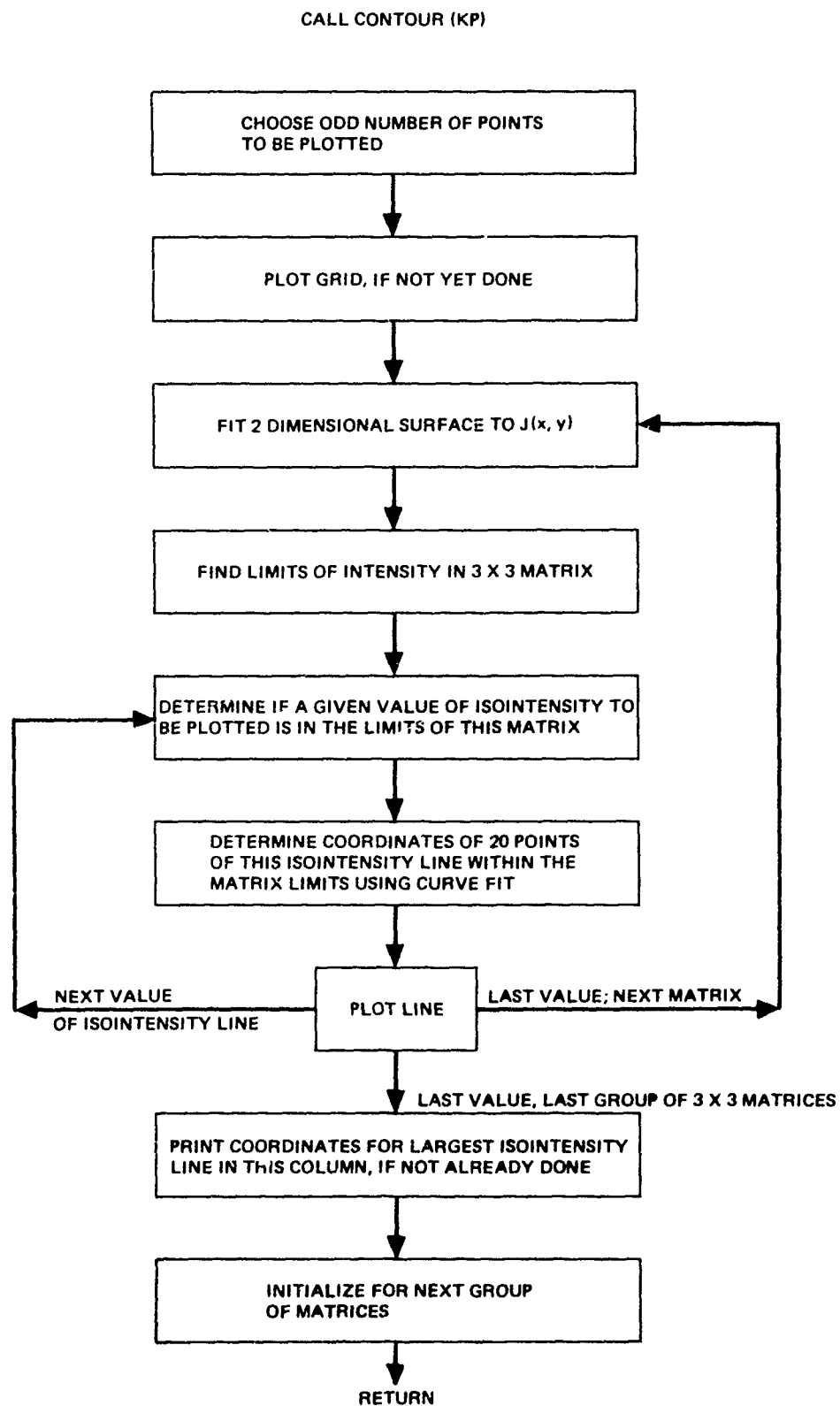


Figure 3-9. Contour Plotter Flow Chart

Section 4

EQUATIONS TO BE SOLVED

The following subsections relate to the program regions shown on the master flow chart, Figure 3-3.

4.1 PROGRAM SECTION 1. INITIAL CALCULATIONS--PARTICLES

4.1.1 EPS, Particle Emissivity Versus Temperature

$$\epsilon_i(T) = \int_0^{\infty} \epsilon_{i\lambda} R_{\lambda}^o(T) d\lambda / \sigma T^4 \quad (4-1)$$

where

$$R_{\lambda}^o(T) = \frac{2\pi c^2 h}{\lambda^5} \left[e^{\frac{hc}{\lambda kT}} - 1 \right]^{-1} = \frac{c_1}{\lambda^5} \left[e^{\frac{c_2}{\lambda T}} - 1 \right]^{-1} \quad (4-2)$$

c = velocity of light

σ = Stefan-Boltzmann constant

$$= 5.672 \times 10^{-12} \text{ j/cm}^2 \text{ } ^\circ\text{K}^4 \text{ sec}$$

$\epsilon_{i\lambda}$ = correlation eq (2-6)

$$c_1 = 3.742 \times 10^4 \text{ w } (\mu\text{m})^4 / \text{cm}^2$$

$$c_2 = 14390 \mu\text{m } ^\circ\text{K}$$

See Subsection 3.1.3 for description of Gaussian Integration.

4.1.2 TCOOL, Particle Temperature Versus Time After Solidification

$$\Delta t_i(T) = \frac{c_p \rho_i}{3\sigma} \int_{T_f}^T \frac{dT}{\epsilon_i T^4} \quad (4-3)$$

where

c_p = particle specific heat

T_f = fusion temperature

4.1.3 BPASS Black Body Intensity in Band of Interest

$$B_{\Delta\lambda}^o(T) = \frac{1}{\pi} \int_{\lambda_{\min}}^{\lambda_{\max}} R_{\lambda}^o(T) d\lambda \quad (4-4)$$

4.1.4 EPSI, Average Particle Emissivity

$$\epsilon_{i,\Delta\lambda}(T) = \frac{1}{\pi} \int_{\lambda_{\min}}^{\lambda_{\max}} \epsilon_{i\lambda} R_{\lambda}^o(T) d\lambda / B_{\Delta\lambda}^o(T) \quad (4-5)$$

4.1.5 TMIN, Particle Equilibrium Temperature at day,

$$\epsilon_i T_{\min,i}^4 = \frac{\alpha_i \Gamma}{2\sigma} + \frac{\alpha_i'}{2} T_E^4 \equiv K \quad (4-6)$$

where

α_i = solar absorptivity [$\sim \epsilon_i(6000^\circ)$]

α_i' absorptivity to Earthshine [$\sim \epsilon_i(T_E)$]

Γ = solar constant = 0.1353 w/cm^2

T_E = effective Earth temperature

at night, Equation 4-6 is used with

$$\Gamma = 0$$

Newton-Rapheson solution is of form

$$T_j = \frac{1}{3} \left[2T_{j-1} + \frac{K}{T_{j-1}^3 \epsilon_i(T_{j-1})} \right] \quad (4-7)$$

4.1.6 DELT, Time to Solidify

$$\Delta t_{i,s} = \frac{\rho_i m \lambda}{3 \epsilon_i \sigma T_f^4} \quad (4-8)$$

where

m = particle density

λ = particle heat of fusion

4.2 PROGRAM SECTION 2. INITIAL CALCULATIONS--GAS

4.2.1 Plume Gas Limiting Velocity

$$V_p = a_o \left(\frac{2}{\gamma - 1} \right)^{1/2} \quad (4-9)$$

where

$$a_o = \sqrt{\gamma R T_c} \quad (4-10)$$

4.2.2 ET, Rotational Energy Versus Temperature

$$E_T = \frac{hcB}{2\sigma^{3/2}} \left[x_1^{1/2} e^{-x_1} - x_2^{1/2} e^{-x_2} + \frac{\sqrt{\pi}}{2} (\text{erf } \sqrt{x_2} - \text{erf } \sqrt{x_1}) \right] \quad (4-11)$$

where

$$B = \text{molecular rotational constant} \quad (4-12)$$

$$\sigma = hcB/kT \quad (4-12)$$

$$x_{1,2} = J_{1,2}^2 \sigma \quad (4-13)$$

$$J_{1,2} = \frac{1}{2B\lambda_{\min, \max}} \quad (4-14)$$

$$\frac{hc}{2} = 0.9928 \text{ w sec cm}$$

$$\frac{hc}{k} = 1.4387 \text{ cm } ^\circ\text{K}$$

4.2.3 JM, Radiative Rate in Band

$$j(T) = \frac{1}{\pi} \int_{\lambda_{\min}}^{\lambda_{\max}} k_\lambda(T) R_\lambda^O(T) d\lambda \quad (4-15)$$

where $k_\lambda(T)$ = the spectral linear absorption coefficient per molecule from Equation 2-8, multiplied by kT/p to change from k_λ per atm.

4.2.4 JT, Radiative Rate Versus Time

$$- \int_{T_{\max}}^T \frac{dE_T}{4\pi j} = t \quad (4-16)$$

where T is found for value of E_T from results of Subsection 4.2.2 above, and then j is formed from results of Subsection 4.2.3.

4.2.5 Closed Form for Radiance

A least-square fit is made to the resulting data of Subsection 4.2.4 of the form

$$j(t) = A (\tau_R - t) \quad (4-17)$$

Substituting Equation 4-17 into 4-16 yields

$$\tau_r^2 = \frac{E_T}{2\pi A} \quad (4-18)$$

which relates τ_r to T_m , the maximum temperature excited by atmospheric collision.

4.3 PROGRAM SECTION 3. SELECTION OF x_1 OR z_1

These are the initial values of x_1 or z_1 to be used to define lines of sight just grazing plume.

For $\alpha \leq 45$ degrees, we have the following:

$$\begin{aligned} \theta_m < 90 \text{ and } 90 - \theta_m > \alpha \\ z_1 &= 0 \end{aligned} \quad (4-19a)$$

$$\begin{aligned} \theta_m > 90 \text{ and } \theta_m - 90 < 90 - \alpha \\ z_1 &= -r_m / \cos \alpha \end{aligned} \quad (4-19b)$$

Otherwise

$$z_1 = r_m \cos(\alpha + \theta_m) / \cos \alpha \quad (4-19c)$$

For $\alpha > 45$ degrees, we have the following:

$$\begin{aligned} \theta_m < 90 \text{ degrees and } \alpha > \theta_m \\ x_1 &= -r_m \cos(\alpha - \theta_m) / \sin \alpha \end{aligned} \quad (4-20a)$$

Otherwise

$$x_1 = -r_m / \sin \alpha \quad (4-20b)$$

In addition, if there is no gaseous radiation ($XM = 0$), the plume angular extent is taken equal to the flow angle of the limiting streamline for the smallest particle (33.52 degrees).

4.4 PROGRAM SECTION 4. CALL STAB--INTERCEPT VALUES

For $0 \leq \alpha < 45$ degrees

$$(1 + \tan^2 \alpha) x^2 + (2 z_1 \tan \alpha) x + y_1^2 + z_1^2 - r_m^2 = 0 \quad (3-27, \text{Section 3})$$

$$(\cot^2 \theta_m - \tan^2 \alpha) x^2 - (2 z_1 \tan \alpha) x + y_1^2 \cot^2 \theta_m - z_1^2 = 0 \quad (3-28, \text{Section 3})$$

and for $45 \text{ degrees} \leq \alpha \leq 90 \text{ degrees}$

$$(1 - \cot^2 \alpha) z^2 + (2 x_1 \cot \alpha) z + y_1^2 + x_1^2 - r_m^2 = 0 \quad (3-29, \text{Section 3})$$

$$(\tan^2 \theta_m - \cot^2 \alpha) z^2 - (2 x_1 \cot \alpha) z - (y_1^2 + x_1^2) = 0 \quad (3-30, \text{Section 3})$$

4.5 PROGRAM SECTION 5. CALL FUN--RADIANCE CALCULATIONS

$$z = \tan \alpha (x - x_1) + z_1 \quad (3-18, \text{Section 3})$$

$$r = \sqrt{z^2 + y_1^2 + [\cot \alpha (z - z_1) + x_1]^2} \quad (4-21)$$

$$\theta = \cos^{-1} (z/r) \quad (4-22)$$

4.5.1 Solid Particles

4.5.1.1 Calculate Throat Radius

$$r_t = r_e / \sqrt{\epsilon}$$

where

$$\epsilon = \frac{1}{M_e} \left(\frac{\gamma + 1}{2} \right)^{-\frac{\gamma + 1}{2(\gamma - 1)}} \left(1 + \frac{\gamma - 1}{2} M_e^2 \right)^{\frac{\gamma + 1}{2(\gamma - 1)}} \quad (4-23)$$

4.5.1.2 Particle Limiting Velocity

$$V_{(\rho_i, \theta)} = [a_i + b_i \exp(-c_i r/r)] \cos^{n_i} \theta \quad (4-24)$$

for $0 \leq \theta \leq \theta_i$ (limiting particle streamline inclination) where the coefficients are given in Table 4-1.

4.5.1.3 Particle Time of Flight

$$t_f = r/V_{(\rho_i, \theta)} \quad (4-25)$$

4.5.1.4 Concentration

$$\text{Density ratio} = \frac{d_i}{d_o} = f(\rho_i, r, \theta) \equiv g_i \left(\frac{r}{r_t} \right)^{e_i} (\cos \theta)^{f_i} \quad (4-26)$$

for $0 \leq \theta \leq \theta_i$ and the coefficients are given in Table 4-2. Concentration is related to density by

$$N_o = d_o \left(\frac{4}{3} \right) \pi \rho_i^3 d' \text{ (XML)}, \quad d' = \text{alumina density} \\ = 4.0045 \text{ g/cm}^3$$

XML = alumina mass fraction

4.5.1.5 Particle Temperature

If $t_f < \Delta t_{i,s}$ (Equation 4-8),

$$T = T_f$$

If $t_f > \Delta t_{i,s}$,

$T =$ maximum of

a. TCOOL ($t_f - \Delta t_{i,s}$), Equation 4-3

or b. TMIN, Equations 4-6 or 4-7

4.5.1.6 Evaluate Integrand

$$J_i = N_o (t' - t_f) f(\rho_i, r, \theta) \pi \rho_i^2 \epsilon_{i,\Delta\lambda}(T) B_{\Delta\lambda}^o(T) \quad (4-27)$$

(See Equations 4-26, 4-4, and 4-5)

Table 4-1

PARTICLE LIMITING VELOCITY COEFFICIENTS

$\rho_i(\mu\text{m})$	θ_i	$a_i(\text{cm/sec})$	$b_i(\text{cm/sec})$	c_i	n_i
0.690	33.52°	30.776×10^4	-4.6692×10^4	0.01874	-0.9048
1.330	28.22	29.701	-3.8094	0.01930	-1.3494
2.000	23.53	28.932	-3.2891	0.01976	-1.5419
2.660	19.76	28.248	-3.0367	0.02012	-1.4886
3.332	16.81	27.520	-2.9355	0.02036	-1.3175
4.000	14.75	26.764	-2.9337	0.02064	-1.0308
4.666	13.56	26.007	-2.9745	0.02092	-0.6425

Table 4-2

CONCENTRATION COEFFICIENTS

$\rho_i(\mu\text{m})$	g_i	e_i	f_i
0.690	0.2201	2.33	2.84
1.330	0.4503	2.24	5.64
2.000	0.5395	2.17	6.75
2.660	1.0308	2.12	7.30
3.332	1.3339	2.08	6.95
4.000	0.8870	2.00	6.30
4.666	0.5396	2.04	-2.05

4.5.2 Excited Gases

4.5.2.1 Time of Flight

$$t_f = r/V_p, \text{ Equation 4-9} \quad (4-28)$$

4.5.2.2 Gas Concentration

$$n_p = n_{p,e} \frac{k}{2} \left(\frac{r_e}{r} \right)^2 [\cos(a\theta)]^{k-2} \quad (4-29)$$

where

$$k = \gamma (\gamma - 1) M_e^2 \quad (4-30)$$

$$a = \frac{90 \text{ degrees}}{\theta_e + (\nu_{\max} - \nu_e)} \quad (4-31)$$

$$\nu_{\max} = \left[\sqrt{\frac{\gamma + 1}{\gamma - 1}} - 1 \right] 90 \text{ degrees} \quad (4-32)$$

$$\nu_e = \sqrt{\frac{\gamma + 1}{\gamma - 1}} \tan^{-1} \sqrt{\frac{\gamma - 1}{\gamma + 1} (M_e^2 - 1)} - \tan^{-1} \sqrt{M_e^2 - 1} \quad (4-33)$$

and $n_{p,e} = d_e N_o / \bar{M}$

N_o = Avogadro's number

\bar{M} = mean molecular weight

and remaining parameters are input data

4.5.2.3 Relative Velocity

$$x = \cot a (z - z_1) + x_1$$

$$V_{\infty} = \left[\left(\frac{xV_p}{r} + V_1 \right)^2 + \left(\frac{yV_p}{r} + V_2 \right)^2 + \left(\frac{zV_p}{r} + V_3 \right)^2 \right]^{1/2} \quad (4-34)$$

where

$$V_1 = V_R \sin \phi_R \cos \psi_R$$

$$V_2 = V_R \sin \phi_R \sin \psi_R$$

$$V_3 = V_R \cos \phi_R$$

4.5.2.4 Excitation Temperature

T_m = maximum of

A. T of ambient atmosphere (See Equation 4-40)

B. $T_M = 9.61194 \times 10^{-9} V_{\infty}^2 + 1.07636 \times 10^{-3} V_{\infty} \quad (4-35)$

(V_{∞} in cm/sec)

4.5.2.5 Attenuation Factor

$$A_t = S \int_0^r n_\infty(h_s) ds \quad (2-7)$$

where the altitudes, h_s , are found from

$$h_s = h_o + s \left(\frac{h - h_o}{r} \right)$$

4.5.2.6 Radiative Relaxation Time

$$\tau_r = \tau_r(T_m), \text{ Equation 4-18} \quad (4-36)$$

4.5.2.7 Radiance Calculation

$$J_g = V_\infty n_\infty S \tau_r [1 - \exp(t_f/\tau_r)] n_p j \exp(-A_t) x_n \quad (4-37)$$

where for $t_f \geq \tau_r$

$$j = \frac{A\tau_r}{2}, \quad (\text{See Equation 4-17}) \quad (4-38)$$

and for $t_f < \tau_r$

$$j = A \left(\tau_r - \frac{t_f}{2} \right) \quad (4-39)$$

n_∞ is found from atmosphere table, $n_\infty(h)$

S = built-in value of cross section

$$h = h_o + x \sin \phi_u \cos \psi_u + y \sin \phi_u \sin \psi_u + z \cos \phi_u \quad (4-40)$$

(See Equation 3-40)

4.5.3 Total Radiance

If both particles and gas are included the total integrand is

$$J = J_p + J_g \quad (4-41)$$

4.6 PROGRAM SECTION 6. CALL SIMP--INTEGRATION THROUGH PLUME Simpson's rule

$$\int_s^{s+\Delta s} f(x) dx = \frac{\Delta s}{3} \left[f(s) + 4 f\left(s + \frac{\Delta s}{2}\right) + f(s + \Delta s) \right] \quad (4-42)$$

Step size is controlled by a geometric progression, so that each step is 10 - larger or smaller than the previous, and minimum step size occurs at $z = 0$ or $x = 0$ planes.

4.7 CURVE FITTING ROUTINE, CF

An auxiliary routine was written to obtain a curve fit to the intensity data to facilitate plotting. The curve fit for intensity is of the form

$$I = \sum_{j=1}^6 a_j \phi_j(x, y)$$

where

$$\phi_j = x^2, x, y^2, y, xy, 1$$

The least square fit condition is (for nine points)

$$\frac{d}{da_k} \sum_{i=1}^9 \left[I_i - \sum_{j=1}^6 a_j \phi_j(x_i, y_i) \right]^2 = 0$$

The resulting matrix of equations is

$$(C)(a) = (I)$$

where

$$\begin{pmatrix} x_i^4 & x_i^3 & x_i^2 y_i^2 & x_i^2 y_i & x_i^3 y_i & x_i^2 \\ & x_i^2 & x_i y_i^2 & x_i y_i & x_i^2 y_i & x_i \\ & & y_i^4 & y_i^3 & x_i y_i^3 & y_i^2 \\ & & & y_i^2 & x_i y_i^2 & y_i \\ & & & & x_i^2 y_i^2 & x_i y_i \\ & & & & & 1 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ a_4 \\ a_5 \\ a_6 \end{pmatrix} = \begin{pmatrix} I_i x_i^2 \\ I_i x_i \\ I_i y_i^2 \\ I_i y_i \\ I_i x_i y_i \\ I_i \end{pmatrix}$$

(Symmetric)

and each term in the C and I matrices is summed over all points, e.g.

$$C_{11} = \sum_{i=1}^9 z_i^4$$

This technique of extrapolation is similar to that used by Dailey (Reference 12). Nine points (3 x 3 square) are selected and the curve fit made. Coordinates of the isointensity lines are then deduced from the curve fit by solving the quadratic equation for y, given I and x. These lines are plotted, and then the next group of nine points are selected. The procedure is repeated until the plume is covered.

Section 5

SAMPLE PROBLEM AND OPERATING INSTRUCTIONS

A list of all input quantities is given in Table 3-1* in Subsection 3.2, with their appropriate units. Additional description of the input parameters especially the Flags, is available in the problem output shown in Appendix B. The input is, in general, standard NAMELIST format (see listing of sample data in Appendix A). Following the NAMELIST input is a single card consisting of a case description (any 80 alphanumeric characters) which is printed at the start of the case. To signify the end of the run (no more cases), a NAMELIST-format input is required setting the flag, END, equal to 1. No case-description card should follow here.

A sample problem follows for an engine typical of a solid-propellant interceptor sustainer rocket. Since the radiating species mole fraction (XM) has been set equal to zero, only radiation from the alumina particles will show up. The specification of look angle equals zero implies a broadside look. The vertical angles equal to 90 degrees means that the plume axis is parallel to the Earth's surface. The values of 90 degrees for the velocity vector imply that the free-stream velocity is perpendicular to the plume axis. The final columns of printed output present x, y coordinates in cm followed by the intensity of w/cm^2 sr; many zero values of intensity result because the particles cannot expand through the total plume. The output from the sample case follows: (Note: The "error summary" in the output refers to reading in only five values for density versus time instead of maximum allowable value of 50 — no calculation errors result).

All intermediate calculations leading up to calculation of radiance (emissivities, etc.) need not be repeated on successive cases, if not changed. These calculations will be skipped if ISL is entered as -1 (solid propellant) or -2 (gas only)

*Note: If ISL = +1 or -1, the density table (PDE) should be for chamber density and not exit plane density.

Section 6

RECOMMENDATIONS FOR FUTURE WORK

This version of FLAME forms a useful framework for the calculation of plume radiance including a variety of mechanisms. The main accomplishments of this contract were to formulate the integration procedure and geometric-orientation choice to relieve the engineer of tedious extrapolations to attempt to relate flight data to calculations. Because of the large number of options included in FLAME, a comprehensive checkout was not possible during this contract (over 50 runs were made). A streamlining of the output and calculation procedures would be possible through a more extensive exercising of the FLAME options. A variety of values for look angle, altitude, times, orientations engine characteristics, relative velocity, etc. should be used as FLAME input to see if further refinements are needed.

Currently the code calculates only half of the plume in a single run which is sufficient for axisymmetry, such as for particle radiation only. The code should be refined to allow total plume extent to be calculated.

A better contour-plotting routine might be found by limiting the closeness of the isointensity lines, and constraining the curve fit to obtain smoother data. Automatic labeling of the contours would be a nice added feature.

The theoretical models used in FLAME are becoming rapidly antiquated. A molecular radiation model based on non-Boltzmann distributions should be investigated and incorporated.

Additional extensive flow-field revision would be useful based on sophisticated multiphase transient plume codes,

Section 7
REFERENCES

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Appendix A
LISTING OF FLAME CODE

C PROGRAM FLAME, VERSION 1

```

PROGRAM P2170(INPUT,OUTPUT,FILMPL,TAPE5=INPUT,TAPE6=OUTPUT,
1          TAPE48=FILMPL,TAPE16)
  DIMENSION PX(3),PY(100),PL(3,100),FC(20),IP2(20)
  DIMENSION PDE(50),TID(50), WLE(8),EPT(20,10),TET(20),TCOOL(
120,10),TCI(20),BPASS(20),FTL(20,10),TMIN(10),DELT(10),ET(20),IR(20
2),JM(20),JT(20),TJ(20),XP(8),W(8),RP(10),TCOOL(20),TCIX(20,10),CASE
3(10),TIDP(51),PDEP(51),T1(10),T2(10),WP(20),CU(2)
  EQUIVALENCE (TCIX,TIDP),(TCIX(1,4),PDEP),(T1,JT),(T2,JT(11))
  REAL IA,ME,JM,JT,LIP
  INTEGER END
  COMMON/MCG/ JO,IPL,PX,PL,PY,FC,XL,XR,YB,YT,IP,IP2
  COMMON/MSTCF/IA,AL,RD,CONF,RM,X1,Y1,RA
  COMMON/MF/ISL,NP,G,RE,GAMMA,ME,AN,TIG,TID,PDE,NR,XML,DELT,TMIN,TCI
1,TCOOL,PI,RP,TEI,ETL,BPASS,TF,VP,P2, XI2,TR,ET,A,V1,V2,V3,HO,WMOL
2,XM,XP,W,GL
  DATA XP,W,RA,PI,RP,WP/-.9602898565,-.7966664774,-.5255324099,-.183
14346425,.1834346425,.5255324099,.7966664774,.9602898565,.101228536
23,.2223810344,.3137066458,.3626837834,.3626837834,.3137066458,.222
33810344,.1012285363,90.,3.1415927,57.2957795131,.001,19*1./
  DATA NP,RP/7.46,604E-5,40.E-5,33.315E-5,26.6E-5,20.001E-5,13.3015E
1-5,6.9007E-5/,R/9.49/
  DATA TR/100.,400.,800.,1200.,1600.,2000.,2400.,2700.,3000.,3300.,
13400., 3900.,4200.,4500.,5000.,5500.,6000.,7000.,9000.,11000./
  DATA TFI/99.,150.,200.,250.,300.,350.,400.,450.,500.,600.,700.,
1800.,1600.,1200.,1400.,1700.,2000.,2500.,3200.,6001./,
2 TF,DF,CP/2323.,4240.,3.2/
C ***** REGION 0, INPUT DATA *****
  NAMELIST /IN/TIG,RE,GAMMA,ME,LIP,TE,AL,HO,VR,PHIU,XIU,TC,PHIR,XIR,
1WLMIN,WLMAX,XM,IND,ISL,INT,LINES,NR,TID,PDE,END,IPL
  GL = ALOG(10.)
  IPT = 0
1 READ(5,IN)
  IF (FND.EQ.1) STOP
  JP = 1
  KP = 1
  IPL = 0
  DO 32 I = 1,20
32 IP2(I) = 0
  IP = 0
  READ(5,200) CASE
  WRITE(6,100) CASE,TIG,RE,GAMMA,ME,LIP,TE,AL,HO,VR,PHIU,XIU,TC,PHIR
1,XIR,WLMIN,WLMAX,XM,IND,ISL,INT,LINES,NR,IPL
200 FORMAT(10A8)
100 FORMAT(1H1,3X,67HPROGRAM P2170 FLAME ROCKET EXHAUST PLUME
1OPTICAL SIGNATURES ///20X,10A8//11H INPUT DATA//29H TIME AFTER IGN
2ITION-TIG(SEC)F10.4,11X,25HNOZZLE EXIT RADIUS-RE(CM)F10.4,13X, 5HG
3AMMAF10.4/ X,19HEXIT MACH NUMBER-MFF10.4,20X,25HNOZZLE LIP ANGLE-
4LIP(DEC)F10.4,13X,26HEFFECTIVE EARTH TEMP-TE(K)F8.2/ X,18HLOOK AN
5GLE-AL(DEC)F8.2,23X,22HENGINE ALTITUDE-HO(FT)F9.0,17X,23HROCKET VE
6LOCITY-VR(FPS)F9.2/ X,27HVERT. POLAR ANGLE-PHIU(DEC)F8.2,14X,30HV
7ERT. AZIMUTHAL ANGLE-XIU(DEC)F8.2,10X,19HCHAMBER TEMP.-TC(K)F8.2/
8 X,26HVEL. POLAR ANGLE-PHIR(DEC)F8.2,15X,29HVEL. AZIMUTHAL ANGLE-X
9IR(DEC)F8.2,16H BAND LIMITS AREF6.2,3H TOF6.2,19X,30HRADIATING SPE
FCIIS MASS FRAC-XM F8.5,// 23H CONTROL FLAGS F
AOLLO--//8H IND =15,25H (1 FOR DAY, 0 FOR NIGHT)/8H ISL =15,79
RH (1,-1 FOR PARTICLES; 0,-2 FOR GAS ONLY; NEGATIVE TO SKIP INITIAL
G CALCULATIONS) /8H INT =15,56H (APPROXIMATE
CNUMBER OF INTEGRATION STEPS THROUGH LAYER)/8H LINES =15,56H (APPRO

```

```

        DIMATE NUMBER OF LINES OF SIGHT IN ONE DIMENSION)/8H      NR =15,44H
        E (NUMBER OF ENTRIES IN FIRING HISTORY TABLE)/8H IPLOT =15,42H (0 F
        FOR NO PLOT, 1 FOR ISOINTENSITY PLOTS))
C ***** INITIALIZATION OF SD 4020 CAMERA
        IF(IPLOT.NE.0) GO TO 30
31 IR = (NR+2)/3
        K = 1
        DO 2 I = 1,IR
        DO 2 J =1,3
        L = I + (J-1)*IR
        TIDP(K) = TID(L)
        PDEP(K) = PDE(L)
        2 K = K +1
        WRITE (6,101)(TIDP(I),PDEP(I),I =1,NR)
101 FORMAT(72H0 FIRING HISTORY TABLE FOLLOWS, EXIT-PLANE DENSITY (G/CM
        13) VS TIME (SEC) //3(4X,4HTIME,5X,7HDENSITY,5X)/(3(F10.4,E11.3,4X)
        2)
C
C ***** REGION 1, INITIAL CALCULATIONS - PARTICLES *****
C
        CC =(WLMAX - WLMIN)/2.
        DD =(WLMAX + WLMIN)/2.
        IF(ISL.NE.1) GO TO 3
        DO 4 I = 1,20
        T = TET(I)
C FIND WAVELENGTH WHERE PLANCK FUNCTION IS GREATEST. TRANSFORM INTEGRAL
        WW = 2898./T
        C = 3.9*WW
        D = 4.1*WW
C PERFORM GAUSSIAN INTEGRATION (EQUATIONS 4-1,-2,-4,-5)
        RPASS(I) = 0.
        R1X = 0.
        DO 7 K = 1,NP
        EPT(I,K) = 0.
        7 ETL(I,K) = 0.
        DO 5 J = 1,8
        WL = C*XP(J) + D
        WWL = CC*XP(J) + DD
        R1 = 3.742E4/WL**5/(EXP(14390./WL/T)-1.)
        R2 = 3.742E4/WWL**5/(EXP(14390./WWL/T)-1.)
        R1X = R1X + W(J)*R1
        RPASS(I) = RPASS(I) + R2*W(J)
C SUM OVER PARTICLES
        DO 5 K = 1,NP
        EM1 = EMIT(RP(K),WL)
        EM2 = EMIT(RP(K),WWL)
        EPT(I,K) = EPT(I,K) + W(J)*R1*EM1
        5 ETL(I,K) = ETL(I,K) + W(J)*R2*EM2
        DO 6 K = 1,NP
        EPT(I,K) = EPT(I,K)/R1X
        6 ETL(I,K) = ETL(I,K)/RPASS(I)
        4 RPASS(I) = RPASS(I)*CC/3.14159
        WRITE(6,103) (RP(I),I=1,NP)
103 FORMAT(46H0 INITIAL CALCULATIONS FOR PARTICLE PROPERTIES//35X,16HT
        10TAL EMISSIVITY//16X,27HPARTICLE RADIUS (CM X 10E4)/13H      TEMP(K)
        2 ,4P10F10.3//)
        DO 20 I = 1,20
        20 WRITE(6,104) TET(I),(EPT(I,J),J=1,NP)
104 FORMAT(1H ,F9.1,3X,10F10.5)

```

```

      WRITE(6,107) WLMIN,WLMAX,(RP(I),I=1,NP)
107 FORMAT(////13H0 RAND DATA,(,F5.1,4H TO ,F5.1,1H)//32X,18HEMISSIVIT
      1Y IN RAND//12X,42HPLANCK FUNC. PARTICLE RADIUS (CM X 10E4)/25H
      2 TEMP(K) (W/CM2/SR) ,4P10F9.3//)
      DO 22 I = 1,20
      22 WRITE(6,108) TET(I),RPASS(I),(FTL(I,J),J = 1,NP)
108 FORMAT(1H ,F9.1,F13.4,2X,10F9.5)
C COOLING CURVES FOR SOLID PARTICLES (SECTION 4.1.2,EQS. 4-3 AND 4-8)
      DT = (TF-100.)/19.
      TCOO(1) = TF
      DO 8 K = 1,NP
      TCIX(1,K) = 0.
      T1(K) = 1./TF**4*RP(K) /TABLE(TF,TET,EPT(1,K),20,0,1)
      8 DELT(K) = T1(K)*DF/1.7016E-11
      T = TF
      DO 9 I = 2,20
      T = T - DT
      TCOO(I) = T
      DO 9 K = 1,NP
      T2(K) = 1./T**4*RP(K) /TABLE(T,TET,EPT(1,K),20,0,1)
      TCIX(I,K) = TCIX(I-1,K)+(T2(K) + T1(K))/2.*DT
      9 T1(K) = T2(K)
      DO 10 K = 1,NP
      10 TCIX(20,K) = TCIX(20,1)
C REARRANGE TO CHANGE INDEPENDENT VARIABLE TO TIME.
      DT = TCIX(20,1)/722700.
      TIME = -DT
      DO 11 I = 1,20
      TIME = TIME + DT*FLOAT(I**4)
      TCI(I) = TIME/1.7016E-11*CP
      DO 11 K = 1,NP
      11 TCOOL(I,K) = TABLE(TIME,TCIX(1,K),TCOO,20,0,1)
      TCI(20) = 1.E5
      WRITE(6,105) (RP(I),I=1,NP)
105 FORMAT(///25X,34HPARTICLE TEMPERATURE HISTORIES (K)//16X,27HPARTIC
      1LE RADIUS (CM X 10E4)/13H TIME(SEC) ,4P10F10.3//)
      DO 21 I = 1,20
      21 WRITE(6,106) TCI(I),(TCOOL(I,J),J = 1,NP)
106 FORMAT(1H ,F10.3,2X,10F10.1)
      3 IF(1SL.LE.0) GO TO 15
      TO = 400.
      DO 12 K = 1,NP
      TCOOL(20,K) = 70.
C NEWTON-RAPHSON SOLUTION TO EQ. 4-7 (AND EQ. 4-6)
      C = 5.964E10*TABLE(6000.,TET,EPT(1,K),20,0,1)*FLOAT(IND) +
      1TABLE(TET,TET,EPT(1,K),20,0,1)*TE**4/2.
      DO 13 I = 1,100
      TN = (2.*TO + C /TO**3/TABLE(TO,TET,EPT(1,K),20,0,1))/3.
      IF (ABS(TN-TO).LT..1) GO TO 14
      13 TO = TN
      WRITE (6,102) TN,TO,K
102 FORMAT(5RHOERROR, SOLUTION OF EQ.(4-7) DOES NOT CONVERGE. NEW VALU
      1E= G8.2, 12H, OLD VALUE= G8.2,15,18H-TH PARTICLE GROUP)
      14 TMIN(K) = TN
      12 CONTINUE
      WRITE(6,109) (RP(I),I = 1,NP)
109 FORMAT(////27H PARTICLE RADIUS (CMX10E4) ,4P10F10.3)
      WRITE(6,110) (DELT(I),I = 1,NP)
110 FORMAT(27H SOLIDIFICATION TIME (SEC) ,10F10.3)

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      WRITE(6,111) (TMIN(I),I = 1,NP)
111 FORMAT(27H EQUILIBRIUM TEMPERATURE(K) ,10F10.2)
C
C ***** REGION 2, INITIAL CALCULATIONS - GAS *****
C
C IF BAND LIMITS UNCHANGED FROM PREVIOUS RUN, SKIP BAND-ENERGY CALC.
  15 IF(ISL.L1.0) GO TO 16
C DETERMINE ET (EQ. 4-11)
    Q1 = 0.5E4/R/WLMAX
    Q2 = 0.5E4/R/WLMIN
    DO 18 I = 1,20
      T = TR(I)
      SIG = 1.4387*R/T
      XMIN = Q1*Q1*SIG
      XMAX = Q2*Q2*SIG
      ET(I) = 9.928E-4*R/SIG*(Q1*EXP(-XMIN) - Q2*EXP(-XMAX)
    1) + .88623/SORT(SIG)*(ERF(Q2*SORT(SIG)) - ERF(Q1*SORT(SIG)))
C DETERMINE BAND RADIANCE (EQ. 4-15)
    JM(I) = 0.
    DO 17 J = 1,8
      WWL = CC*XP(J) + DD
    17 JM(I) = JM(I) + W(J)*ABCO(WWL,T)/WWL**5/(EXP(14390./WWL/T)-1.)
    18 JM(I) = JM(I)*CC*3.742E4/PI
      WRITE(6,112) (TR(I),ET(I),JM(I),I = 1,20)
112 FORMAT(////43HOROTATIONAL RADIATION FROM WATER MOLECULES./97H
  1AVAILABLE BAND ENERGY = ET (10E-20 W SEC/MOLEC), BAND RADIATIVE RA
  2TE = J (10E-20 W/SR MOLEC) //24X,4HT(K),8X,2HET,12X,1HJ//(22X,F7.1
  3,2X,2E13.3))
C CALCULATION OF RADIATIVE COOLDOWN (EQ. 4-16) (J VS TIME)
    ETX = ET(I)
    TO = 0.
    DET = ET(20)/190.
    JT(I) = JM(I)
    TJ(I) = 0.
    DO 19 I = 2,20
      ETX = ETX + DET
      JT(I) = TABLE(ETX-ET,JM,20,0,1)
      TN = .0/9578/JT(I)
      TJ(I) = TJ(I-1) + (TO + TN)/2.*DET
    19 TO = TN
C FIT DATA WITH FUNCTION L/NEAR IN TIME (EQ. 4-17)
    CALL LSOPOL(TJ,JT,WP,TIDP,20,ETX,2,CU)
    A = CU(2)
    DO 23 I = 1,20
      TCOO(I) = CU(1) + CU(2)*IJ(I)
      TIDP(I) = TABLE(JT(I),JM,ET,20,0,1)
    23 PDEP(I) = TABLE(JT(I),JM,TR,20,0,-1)
      WRITE(6,113) (TJ(I),JT(I),TCOO(I),TIDP(I),PDEP(I),I = 1,20)
113 FORMAT(///28H RADIATIVE DECAY IN BAND //9X, 9HTIME(SEC),6X,1H
  1J,8X,6HJ(FIT),7X,2HET,8X,4HT(K)//(6X,F11.5,2X,4E11.3))
C SOLUTION FOR LIMITING VELOCITY (EQS. 4-9 AND 4-10)
  16 VP = 2.*SORT(TC/WMOL/(2.-2./GAMMA)*8.3144E7)
C
C PREPARATION FOR PLUME INTEGRATION
C
    X1 = 0.
    Y1 = 0.
C CALCULATION OF MAXIMUM PLUME CONE HALF-ANGLE AND MAXIMUM PLUME RADIUS
    GG = SORT((GAMMA+1.)/(GAMMA-1.))

```

```

GM = SORT(ME**2 - 1.)
CONF = LIP + RA*(GG-1) - RD*(GG*ATAN(GM/GG) - ATAN(GM))
RM = VP*TIG
DELS = (RM + RM*COS(AL/RD))*0.05/(1.1**((INT/2) - 1.))
IA = 0.
P1 = PHIR/RD
P2 = PHIU/RD
XI1 = XIR/RD
XI2 = XIU/RD
IF(AL.GT.45.) IA = 1.
DELX = DELS*COS((RA*IA - AL)/RD)
DELY = (RM + RM*SIN(AL/RD))*0.1/(1.1**LINES - 1.)
CD = COS((RA*IA-AL)/RD)
DELXX = DELY/CD
DEL = DELY
JO = 10000
C SELECTION OF ISOINTENSITY LEVELS TO BE PLOTTED
FC(1) = 1.E-11
X10 = SORT(10.)
DO 29 I = 1,19
29 FC(I+1) = FC(I)*X10
C
C ***** REGION 3. SELECTION OF INITIAL VALUE OF X1 *****
C
VV = VR*30.4801
V1 = VV*SIN(P1)*COS(XI1)
V2 = VV*SIN(P1)*SIN(XI1)
V3 = VV*COS(P1)
INOUT = 2
IF(IA.EQ.1.) GO TO 36
IF(XM.EQ.0.) GO TO 27
IF(CONE.GT.90.) GO TO 25
IF(90. - CONE.GE.AL) GO TO 27
34 X1 = RM*COS((CONE + AL)/RD)/COS(AL/RD)
GO TO 27
25 IF(90. - CONE.GE.AL - 90.) GO TO 34
X1 = -RM/COS(AL/RD)
GO TO 27
36 IF(XM.NE.0.) GO TO 33
X1 = -RM*COS((AL - 33.52)/RD)/SIN(AL/RD)
GO TO 27
33 IF(CONE.LT.90..AND.CONE.LT.AL) GO TO 35
X1 = -RM/SIN(AL/RD)
GO TO 27
35 X1 = -RM*COS((AL - CONE)/RD)/SIN(AL/RD)
27 XPLOT = 0.
PX(1) = XPLOT
DELX1 = (DELXX + ABS(X1)/10.)/1.1
R = 1./1.1
C PLOT LIMITS FOR GRID
XL = 0.
XR = RM - X1*COS((RA*IA - AL)/RD)
IF(IA.EQ.1..AND.XM.EQ.0.) XR = 2.*RM*SIN(33.52/RD)
1 *SIN(AL/RD)
YB = 0.
YT = XR
C
C ***** REGION 4. CALCULATE INTEGRATION BOUNDARIES *****
C

```

```

24 CALL STAR(PMIN,PMAX,INOUT)
   DELY = DELY*1.1
   IF(INOUT.EQ.0) GO TO 26
   Y1 = 0.
   DELY = DEL
   JO = MIN0(JO,JP)
   KP = KP + 1
   IF(KP.EQ.4.AND.IPLOT.NE.0) CALL CONTUR(KP)
   JP = 0
   IF (X1.GT.0.) R = 1.1
   DELX1 = DELX1*R
   X1 = X1 + DELX1
   XPLOT1 = XPLOT + DELX1*CD
C  STURE VALUES (PX,PL,PY) FOR PLOTTING
   PX(KP)=XPLOT
   CALL STAR(PMIN,PMAX,INOUT)
   IF(INOUT.EQ.0) GO TO 26
   GO TO 1

C
C ***** REGION 5, INITIAL EVALUATION OF INTEGRAND *****
C
C   26 CALL FUN(SOLD,PMIN)
C
C ***** REGION 6, INTEGRATE THROUGH PLUME *****
C
   CALL SIMP(SOLD,SINT,PMAX,PMIN,DELX)
   SINT = SINT*.6666667
   WRITE(6,114) XPLOT,Y1,SINT
114 FORMAT(4X,2F10.2,E11.3)
   JP = JP + 1
   PL(KP,JP)= SINT
   PY(JP) = Y1
   Y1 = Y1 + DELY
   GO TO 24
30 IF(IPT.EQ.0) CALL CAMRAV(IPLOT)
   IPT = 1
   GO TO 31
END

```



```

SUBROUTINE SIMP(SOLD,SINT,PMAX,PMIN,DEL)
SINT = SOLD
DELX = ( ABS(PMIN) /10. + DEL)/1.1
A = 1./1.1
X = PMIN
1 SUM = 0.
C = 4.
DELX = DELX*A
4 X = X + DELX
IF(X.GT.PMAX) GO TO 2
CALL FUN(SS,X)
SS = SS*DELX
SUM = SUM + C*SS
IF(C.EQ.4.) GO TO 3
IF(X.GT.0.) GO TO 6
5 SOLD = SS
SINT = SINT + SUM
GO TO 1
3 C = 2.
GO TO 4
2 SINT = SINT - SOLD
RETURN
6 A = 1.1
GO TO 5
END

```

```

SUBROUTINE STAB(PMIN,PMAX,INOUT)
  DIMENSION X(2)
  COMMON/STC/X,      XX
  COMMON/MSTCF/IA,AL,RD,CONE,RM,X1,Y1,RA
  REAL IA
  IF(INOUT.NE.2) GO TO 1
  A0 = ABS (TAN((RA*IA - AL)/RD))
  A1 = A0*A0
  A2 = (TAN((RA*(IA-1.) - CONE)/RD))**2
1 INOUT = 0
  B1 = X1*A0/(1. + A1)
  B = Y1**2
  C1 = B + X1**2      --RM*RM
  IF(IA.EQ.0.)B = -B*A2
  C2 = X1*X1 + B
  RAD = B1*B1 - C1/(1. + A1)
  JI = 1
C  CHECK FOR LINE-OF-SIGHT BEYOND R=RP.
  IF(RAD.LT.0.) GO TO 5
  RAD = SQRT(RAD)
  XX = -B1 + RAD
  CALL CHECK(JI)
  IF(JI.GT.2) GO TO 4
  XX = -B1 - RAD
  CALL CHECK(JI)
  IF(JI.GT.2) GO TO 4
3 RD = A1 - A2
  B2 = X1*A0/RD
C  CHECK FOR LINE-OF-SIGHT BEING PARALLEL TO PLUME BOUNDARY.
  IF(ABS(BD).LT.1.E-8) GO TO 2
  RAD = B2*B2 - C2/BD
  IF (RAD.LT.0.) GO TO 5
  RAD = SQRT(RAD)
  XX = -B2 + RAD
  CALL CHECK(JI)
  IF(JI.GT.2) GO TO 4
  XX = -B2 - RAD
7 CALL CHECK(JI)
  IF(JI.GT.2) GO TO 4
5 INOUT = 1
  RETURN
2 IF(AL.EQ.0..OR.AL.EQ.RA) GO TO 5
  XX = C2/B2
  GO TO 7
4 IF (X(1).GT. X(2)) GO TO 6
  PMIN = X(1)
  PMAX = X(2)
  RETURN
6 PMIN = X(2)
  PMAX = X(1)
  RETURN
END

```

```

SUBROUTINE CHECK(I)
  DIMENSION X(2)
  COMMON/STC/X,      XX
  COMMON/MSTCF/IA,AL,RD,CONE,RM,X1,Y1,RA
  REAL IA
C  CHECK TO SEE IF XX REPRESENTS X OR Z.
  IF(IA.EQ.0.) GO TO 1
  ZZ = XX
  RR =  SORT(Y1*Y1 + (ZZ*TAN((RA-AL)/RD) + X1)**2 + ZZ*ZZ)
  2 IF (RR.EQ.0.) RETURN
  THETA = RD*ACOS(ZZ/RR)
C  CHECK TO SEE IF INTERSECTION IS ON CONICAL PLUME BOUNDARY.
  IF(ABS((THETA-CONE)/CONE).LT..001) GO TO 3
C  INTERSECTION IS ON SPHERICAL CAP OF PLUME.
  IF(THETA.GT.CONE) RETURN
C  DO THE FOLLOWING IF THIS POINT IS INDEED ONE OF THE INTEGRATION LIMITS
  4 X(I) = XX
  I = I+1
  RETURN
  1 ZZ =  XX*TAN(AL/RD) + X1
  RR = SORT(Y1*Y1 + XX*XX + ZZ*ZZ)
  GO TO 2
C  INTERSECTION IS ON CONICAL BOUNDARY
  3 IF(RR/RM.LT.1.001) GO TO 4
  RETURN
END

```

```

SUBROUTINE FUN(O,P)
  DIMENSION PDF(50),TID(50),TET(20),TCOOL(20,10),TCI(20),BPASS(20),
  1FTL(20,10),TMIN(10),DELT(10),ET(20),TR(20),RP(10),N(10),AB(10),B(1
  20),C(10),TH(10),D(10),E(10),F(10),H(20),TMT(20),CAT(20),XP(8)
  3,W(N)
  REAL MF,IA,J,JG,N
  COMMON/MF/ISL,NP,R,PE,GAMMA,MF,AN,TIG,TID,PDF,NR,XML,DELT,TMIN,TCI
  1,TCOOL,P1,RP,TF1,FTL,BPASS,TF,VP,PHIU,XIU,TR,ET,A,V1,V2,V3,H0,WMO
  2,XM,XP,W,GL
  COMMON/MSTCF/IA,AL,RD,CONE,RM,X1,Y1,RA
  DATA WMO,XML,DP/17. , .36 , 4.0045/,S,CF/.7E-15,3.531445E-5/
  C CURVE FIT COEFFICIENT V AND RHO CORRELATIONS FOR TWO-PHASE FLOW
  DATA AB/2.6007E5,2.6764E5,2.7520E5,2.8248E5,2.8932E5,2.9701E5,3.07
  176F5/,B/-2.9745E4,-2.9337E4,-2.9355E4,-3.0367E4,-3.2891E4,-3.8094E
  24,-4.6692E4/,C/.02092,.02064,.02036,.02012,.01976,.01930,.01874/,N
  3/-.6425,-1.0308,-1.3175,-1.4886,-1.5419,-1.3494,-.9048/,D/.5396,.8
  487,1.3339,1.0308,.5395,.4053,.2201/,E/2.04,2.06,2.08,2.12,2.17,2.2
  54,2.33/,F/-2.05,6.30,6.95,7.30,6.75,5.65,2.84/,TH/13.56,14.75,16.8
  61,19.76,23.53,28.22,33.52/
  C H(FT),T, AND LOG(N/FT3) FROM U.S.STANDARD ATMOSPHERE (1962)
  DATA HT,TMT,CAT/3.E5,3.5F5,4.E5,4.5E5,5.F5,5.5F5,6.F5,6.5E5,7.E5,8
  1.E5,9.F5,1.F6,1.1F6,1.3F6,1.5E6,1.7E6,1.9F6,2.1F6,2.3E6,3.0E6,184.
  295,241.74,385.34,663.20,924.16,1084.09,1170.47,1230.3,1275.21,1346
  3.67,1396.14,1435.68,1455.53,1485.69,1489.35,1496.29,1500.57,1506.3
  43,2*1507.58,18.15,16.98,16.05,15.41,15.01,14.75,14.55,14.36,14.20,
  513.90,13.64,13.40,13.17,12.76,12.41,12.08,11.76,11.49,11.20,10.3/
  X = P
  JG = 0.
  S1 = 0.
  G2 = GAMMA - 1.
  Z = X1 + ABS(TAN((RA*IA-AL)/RD))*X
  IF(IA.EQ.1.) GO TO 6
  7 R = SQRT(X*X + Y1*Y1 + Z*Z)
  AN = ACOS(Z/R)*RD
  IF(IABS(ISL).NE.1) GO TO 1
  G1 = (GAMMA + 1.)/2.
  C CALCULATE THROAT RADIUS FROM EXIT RADIUS AND EXPANSION RATIO
  RT = RE/SQRT(G1**(-G1/G2)*(1.+G2*ME*ME/2.))*G1/G2/ME)
  CT= COS(AN/RD)
  DO 2 I = 1,NP
  C CALCULATE PARTICLE TIME OF FLIGHT, EQ. 4-25
  IF(TH(I).LT.AN) GO TO 2
  C ARE WE BEYOND THE LIMITING STREAMLINE FOR THIS SIZE PARTICLE
  TI = R/CT**N(I)/(AB(I)+ R(I)*EXP(-C(I)*R/RT))
  IF(TI.GT.TIG) GO TO 2
  PDO = TARLE(TIG-TI,TID,PDE,NR,0,1)*XML
  C PARTICLE CONCENTRATION
  PD = PDO*D(I)*(RT/R)**E(I)*CT**F(I)/4.18879/DP/RP(I)**3
  C CHECK FOR INCOMPLETE SOLIDIFICATION OF PARTICLE
  IF(TI.LT.DELT(I)) GO TO 4
  T = AMAX1(TMIN(I),TARLE(TI-DELT(I),TCI,TCOOL(1,I),20,0,1))
  3 RO = TARLE(T,TF1,BPASS ,20,0,1)
  C PARTICLE RADIANCE ACCUMULATION, EQ. 4-27
  S1 = S1 + PD*PI*RP(I)**2 *TARLE(T,TET,ETL(1,I),20,0,-1)*RO
  2 CONTINUE
  GO TO 1
  4 T = TF
  GO TO 3
  C GAS-PHASE CALCULATIONS.

```

```

1  TI= R/VP
   IF(TI.GT.TIG.OR.AN.GE.CONE) GO TO 5
   GDO = TABLE(TIG-TI,TID,PDE,NR,0,1)*6.0254E23/WMOL
   XK = GAMMA*G2*MF*ME
C  GAS DENSITY, EQ 4-29 (MODIFIED FOR TWO-PHASE FLOW)
   GD = GDO*XK/2.*(RE/R,**2*XM*(COS(RA/CONE*AN/RD))**(XK-2.))
   IF(GD .LT.1.)GO TO 5
   VX = X*VP/R
   VZ = Z*VP/R
C  RELATIVE VELOCITY, EQ 3-16
   VREL = SORT((VX+V1)**2 + (VZ+V3)**2 + (Y1*VP/R +V2)**2)
C  FIND ALTITUDE IN PLUME EQ. 4-40
   H = HO +(SIN(PHIU))*(X*COS(XIU) + Y1*SIN(XIU)) +Z*COS(PHIU))/30.48
   H = AMAX1(H,3.E5)
   H = AMIN1(H,3.E6)
C  FIND EXCITATION TEMPERATURE, EQ. 4-35
   TA = TABLE(H,HT,TMT,20,0,1)
   TM = AMAX1(TA,VREL*(9.61194E-9*VREL + 1.07636E-3))
C  ATMOSPHERIC CONCENTRATION
   CA =(EXP(TABLE(H,HT,CAT,20,0,-1)*GL))*CF
C  CALCULATION OF ATTENUATION FACTOR, EQ. 2-7
   HD = (H - HO)/2.
   AT = 0.
   DO 8 I = 1,8
   HS = HO + HD*(XP(I) + 1.)
   8 AT = AT + W(I)*EXP(TABLE(HS,HT,CAT,20,0,1)*GL)
   AT = AT*R/2.*CF*S
C  RADIATIVE RELAXATION TIME, EQ. 4-18 AND 4-36
   EM = TABLE(TM,TR,ET,20,0,1)
   TAU = SORT(EM/2./PI/A)
C  COMPUTATION OF EMISSION, EQ. 4-38 AND 4-39
   TX = 1. - TI/TAU
   J = A*TAU/2.
   IF(TX.GT.0.) J = J + J*TX
C  RADIANCE CALCULATION, EQ. 4-37
   JG = VREL*CA*S*TAU*(1.-EXP(-TI/TAU))*GD*J*1.E-20*EXP(-AT)
5  Q = S1 + JG
   RETURN
6  TX = Z
   Z = X
   X = TX
   GO TO 7
END

```

```

      FUNCTION ARCO(WL,TE)
C     COMPUTES ABSORPTION COEFFICIENT PER MOLECULE OF WATER IN THE
C     ROTATIONAL SPECTRUM, SCALED BY 10E20
      IF(T1.EQ.TE) GO TO 1
      T1 = TE
      T = SORT(T1)
      W0 = 10.6*T
      Y0 = 279.96/(T + 50.031) - 3.208
      A = 1.0303E-4 - 105.582E-4/(T-2.896)
1     W = 10000./WL
      ARCO = EXP(2.302585*(Y0 + A*(W - W0)**2/SORT(W)))*1.3803E-2*TE
      RETURN
      END

```

```

      FUNCTION EMIT(R,W)
C COMPUTES THE EMISSIVITY OF ALUMINA PARTICLES OF RADIUS R (CM) AT A
C WAVELENGTH W (MICRONS) BASED ON CORRELATION OF MIE CALCULATIONS.
      IF (A.EQ.0.) GO TO 1
      IF (R.EQ.R0) GO TO 2
      1 A = (11.32 - 2220.*R) *R + .00689
      R0 = R
      2 EMIT = A*(W + 2.5)
      RETURN
      END

```

```

SUBROUTINE CONTOUR(KP)
  DIMENSION TT(9),XP(21),YP(21),PX(3),PY(100),PL(3,100),FC(20),CA(6)
  1,IP2(20)
  DATA TT/9*10H /
  COMMON/MCG/ JO,IPL,PX,PL,PY,FC,XL,XR,YR,YT,IP,IP2
  JO = ((JO + 1)/2)*2 - 1
  IF (IPL.EQ.0) CALL GRID
  J1 = JO - 2
  FD = 0.
  DX = (PX(3) - PX(1))/20.
  DO 1 L = 1, J1, 2
C CURVE FIT TO 3 X 3 MATRIX OF PLUME POINTS
  CALL CF(PX,PY(L),PL(1,L),CA)
C HERE TO STATEMENT 7 -- FINDS WHAT ISOINTENSITY LINES ARE IN MATRIX
  BIG = 0.
  SMALL = 1.E20
  L2 = L+2
  DO 2 I = 1,3
  DO 2 J = L,L2
  A = PL(I,J)
  IF(A.EQ.0.) GO TO 2
  BIG = AMAX1(BIG,A)
  SMALL = AMIN1(SMALL,A)
2 CONTINUE
  YS = PY(L)
  YL = PY(L2)
  DO 9 I = 1,20
  IF (FC(I).LT.SMALL.OR.FC(I).GT.BIG) GO TO 9
  IF (FD.GT.FC(I)) GO TO 7
  FD = FC(I)
  ID = I
  N = L
C HERE TO STATEMENT 3 -- SOLVES FOR COORDINATES OF ISOINTENSITY LINES
  7 K = 0
  X = PX(1)
  DO 3 J = 1,20
  B = CA(4) + CA(5)*X
  A = CA(3)
  C = CA(1)*X**2 + CA(2)*X - FC(I) + CA(6)
  D = B**2 - 4.*A*C
  IF (D.LT.0.) GO TO 4
  D = SQRT(D)
  Y = (-B + D)/2./A
  IF(Y.GT.YS.AND.Y.LT.YL) GO TO 5
  Y = (-B - D)/2./A
  IF(Y.GT.YS.AND.Y.LT.YL) GO TO 5
4 X = X + DX
3 CONTINUE
  IF (K.LE.1) GO TO 9
  CALL AICRT3(0,0,XP,YP,K,1,1,1,58,TT,TT,TT,2,1,16.,16.,2,XL,XR,2,
  1YR,YT)
9 CONTINUE
1 CONTINUE
  DO 6 M = 1,IP
6 IF(IP2(M).EQ.ID) GO TO 10
  IP = IP + 1
  IP2(IP) = ID
  WRITE (6,100) FD ,PX(1),PY(N)
100 FORMAT(22H ISOINTENSITY LINE FOR E9.2,18HW/CM2/SR IS NEAR (E9.2,1H

```



```

      1,E9.2,1H))
C  INITIALIZATION FOR NEXT GROUP OF 3X3 MATRICES
10  PX(1) = PX(3)
    DO 8 I = 1,J0
      8  PL(I,1) = PL(3,I)
        KP = 2
        J0 = 10000
        RETURN
5   K = K + 1
    XP(K) = X
    YP(K) = Y
    GO TO 4
END)

```

```

      SUBROUTINE GRID
C THIS ROUTINE PREPARES THE GRAPH LIMITS FOR PLOTTING ISOINTENSITY LINES
      DIMENSION TITLE(8),XT(9),YT(9),PX(3),PY(100),PL(3,100),FC(20)
      1,IP2(20)
      COMMON/MCG/ JO,IPL,PX,PL,PY,FC,XL,XR,YR,Y2,IP,IP2
      DATA TITLE/3*1H ,6HPLUME ,6HCONTOU,2HRS/,XT/4*1H ,6HX (CM)/,
      1 YT/4*1H ,6HY (CM)/
      CALL ALCRT3(0,0,X,Y,2,3,1,1,42,TITLE,XT,Y1,1,1,16.,16.,2,XL,XR,2,
      1YR,Y2)
      IPL = 1
      RETURN
      END

```

```

      SUBROUTINE CF(X,Y,F,C)
C THIS ROUTINE PERFORMS A LEAST-SQUARE CURVE FIT TO 9 (X,Y) POINTS OF
C FUNCTION F, WHERE
C   F(FIT) = C(1)X**2 + C(2)X + C(3)Y**2 + C(4)Y + C(5)XY + C(6)
      DIMENSION X(3),Y(3),F(3,3),C(6),A(6,6),RT(18),IT(18)
      DO 1 I = 1,6
        C(I) = 0.
      DO 1 J = 1,6
1     A(I,J) = 0.
      DO 2 I = 1,3
        X2 = X(I)**2
        A(1,1) = A(1,1) + X2*X2
        A(1,2) = A(1,2) + X2*X(I)
        A(2,6) = A(2,6) + X(I)
2     A(2,2) = A(2,2) + X2
      DO 3 I = 1,3
        Y2 = Y(I)**2
        A(3,3) = A(3,3) + Y2*Y2
        A(3,4) = A(3,4) + Y2*Y(I)
        A(4,6) = A(4,6) + Y(I)
3     A(4,4) = A(4,4) + Y2
        A(1,1) = 3.*A(1,1)
        A(1,2) = 3.*A(1,2)
        A(2,2) = 3.*A(2,2)
        A(3,3) = 3.*A(3,3)
        A(3,4) = 3.*A(3,4)
        A(4,4) = 3.*A(4,4)
        A(4,6) = 3.*A(4,6)
        A(2,6) = 3.*A(2,6)
        A(2,1) = A(1,2)
        A(4,3) = A(3,4)
      DO 4 I = 1,3
        DO 4 J = 1,3
          X2 = X(I)**2
          Y2 = Y(J)**2
          XY = X(I)*Y(J)
          A(1,3) = A(1,3) + X2*Y2
          A(1,4) = A(1,4) + X2*Y(J)
          A(1,5) = A(1,5) + X2*XY
          A(2,3) = A(2,3) + XY*Y(J)
          A(2,4) = A(2,4) + XY
          A(3,5) = A(3,5) + XY*Y2
          C(1) = C(1) + F(I,J)*X2
          C(2) = C(2) + F(I,J)*X(I)
          C(3) = C(3) + F(I,J)*Y2
          C(4) = C(4) + F(I,J)*Y(J)
          C(6) = C(6) + F(I,J)
4     C(5) = C(5) + F(I,J)*XY
        A(3,1) = A(1,3)
        A(3,2) = A(2,3)
        A(4,1) = A(1,4)
        A(4,2) = A(2,4)
        A(4,5) = A(2,3)
        A(5,4) = A(2,3)
        A(5,1) = A(1,5)
        A(5,2) = A(1,4)
        A(2,5) = A(1,4)
        A(5,3) = A(3,5)
        A(6,6) = 9.

```

```

A(1,6) = A(2,2)
A(6,1) = A(2,2)
A(6,2) = A(2,6)
A(3,6) = A(4,4)
A(6,3) = A(4,4)
A(6,4) = A(4,6)
A(5,5) = A(1,3)
A(5,6) = A(2,4)
A(6,5) = A(2,4)
CALL SID(A,6,6,6,C,1,1,S,IR,RT,IT,SC)
IF(S.LT.3.) WRITE(6,100) X(1),Y(1)
100 FORMAT (38H CURVEFIT COEFFICIENTS INACCURATE AT (E9.2,2H, E9.2,1H)
1)
RETURN
END

```

```

      FUNCTION ERF(ETA)
C     SEE HANDBOOK OF MATHEMATICAL FUNCTIONS WITH FORMULAS, GRAPHS AND
C     MATHEMATICAL TABLES. NBS, APPLIED MATHEMATICAL SERIES 55,
C     PAGE 299, SECTION 7.1.26 (1964)
      DATA P/ 0.3275911/
      DATA A1,A2,A3,A4,A5 / 0.254829592,-0.284496736,1.421413741,
1-1.453152027,1.06140543 /
      T = 1.0/(1. + P*ABS(ETA))
      S = T*(A1 + T*(A2 + T*(A3 + T*(A4 + A5*T))))
      ERF = 1. - S*EXP(-ETA**2)
      IF (ETA.LT. 0.0) ERF = -ERF
      RETURN
      END

```

FUNCTION TABLE(XX,X,FX,M0,L,IFX)	0010
DIMENSION X(M0),FX(M0)	0020
M=M0	0030
IF (XX-X(M)) 1,1,3	0040
1 IF (XX-X(1)) 2,6,6	0050
2 TABLE=FX(1)	0060
GO TO 4	0070
3 TABLE=FX(M)	0080
4 WRITE(6,5)	0090
5 FORMAT(1H0,10Y,3H*** ERROR, RANGE OF TABLE EXCEEDED ***)	0100
WRITE (6,101) XX,X(1),X(M),M0	
101 FORMAT (3F13.6,110)	
IF (1EX.F0.0) GO TO 100	0110
STOP	0120
6 IF (1FX.11.0) GO TO 14	
M1 = 1	
M3=M	0140
M2=M/2	0150
7 IF (X(M2)-XX) 9,8,8	0160
8 M3=M2	0170
GO TO 10	0180
9 M1=M2	0190
10 J=(M1+M3)/2	0200
IF (J.E0.M2) GO TO 11	0210
M2=J	0220
GO TO 7	0230
11 IF (X(J)-XX) 12,12,13	0240
12 M3=J+1	0250
M1=J	0260
GO TO 14	0270
13 M3=J	0280
M1=J-1	0290
14 IF (L.NE.0) GO TO 15	0300
TABLE=FX(M1)+(XX-X(M1))*(FX(M3)-FX(M1))/(X(M3)-X(M1))	0310
GO TO 100	0320
15 IF (M1.GT.1) GO TO 16	0330
I1=1	0340
I2=2	0350
I3=3	0360
GO TO 17	0370
16 IF (M.GT.M3) GO TO 18	0380
I3=M	0390
I2=M-1	0400
I1=M-2	0410
17 DF32=FX(I3)-FX(I2)	0420
DF21=FX(I2)-FX(I1)	0430
DENOM=(X(I3)-X(I2))*(X(I3)-X(I1))*(X(I2)-X(I1))	0440
A=((X(I2)-X(I1))*DF32-(X(I3)-X(I2))*DF21)/DENOM	0450
R=((X(I2)**2-X(I1)**2)*DF32-(X(I3)**2-X(I2)**2)*DF21)/DENOM	0460
TABLE=FX(I3)-A*(X(I3)**2-XX**2)+R*(X(I3)-XX)	0470
GO TO 100	0480
18 I4=M3+1	0490
I3=M3	0500
I2=M1	0510
I1=M1-1	0520
DF43=FX(I4)-FX(I3)	0530
DF32=FX(I3)-FX(I2)	0540
DF21=FX(I2)-FX(I1)	0550
A1=(X(I3)-X(I2))*(X(I3)-X(I1))*(X(I2)-X(I1))*DF43	0560

A2=(X(I4)-X(I3))*(X(I2)-X(I1))*(X(I4)+X(I3)-X(I2)-X(I1))*DF32	0570
A3=(X(I4)-X(I3))*(X(I4)-X(I2))*(X(I3)-X(I2))*DF21	0580
DENOM=(X(I4)-X(I3))*(X(I4)-X(I2))*(X(I4)-X(I1))*	0590
1 (X(I3)-X(I2))*(X(I3)-X(I1))*(X(I2)-X(I1))	0600
A=(A1-A2+A3)/DENOM	0610
R1=(X(I3)-X(I2))*DF43-(X(I4)-X(I3))*DF32	0620
R2=(X(I4)-X(I3))*(X(I3)-X(I2))*(X(I4)-X(I2))	0630
R=R1/142-A*(X(I4)+X(I3)+X(I2))	0640
C=DF43/(X(I4)-X(I3))-A*(X(I4)**2+X(I3)**2+X(I4)*X(I3))	0650
1 -R*(X(I4)+X(I3))	0660
D=FX(I4)-(A*FX(I4)**2+R*FX(I4)+C)*X(I4)	0670
TARLE=A*YY**3+R*XX**2+C*XX+D	0680
100 RETURN	0690
END	0700

"
 \$IN TIG= 1.,RF = 5.7,GAMMA = 1.33, ME = 4.,LIP =15.,TE = 250.,AL = 0.,
 H0= 5.E5, VR = 2.E4,PHIU = 90.,XIU = 90.,PHIR = 90.,XIR = 90.,TC=1255.,
 WLMIN=8.,WLMAX=14., IND =0, ISL =1, INT = 20, LINES = 30, NR = 5,
 TID= -.1,0.,2.,2.1,10.,POE = 0.,2*1.1E-3,2*0., ILOT = 1,XM = 0.0 \$
 BROADSIDE LOOK, PARTICLE ONLY, 600 LB THRUST, AT NIGHT
 \$IN END = 1\$

Appendix B
OUTPUT FOR SAMPLE PROBLEM

GOE MAP	17.52.37	NORMAL	CONTROL	000100	055005	000000	000000
TIME	LOAD MODE	L1	L2	TYPE	CALL	CALL	CALL
FMA LOADER	073771	FMA TABLES	066023	USER	CALL	CALL	CALL
PROGRAM	ADDRESS	LA	VE	LE	COMMON	COMMON	COMMON
P2170	002347						
GR10	016337						
CF	016441						
ST4P	017110						
ALCO	012316						
EMIT	017324						
STAB	017366						
CHECK	017622						
SUN	012343						
ERF	021046						
YASLE	021160						
CONTR	022025						
ACSPER	022304						
EX2	022416						
SINCOS	022375						
TAN	022676						
ALN-OG	023003						
ASINCOS	023074						
ATAN	023232						
RADEX	023336						
RADEX	023367						
RUCFEO	023446						
GET9A	023552						
CAMRAV	023521						
CHSIZV	023737						
FLAGSV	024537						
HU-LV	024604						
INT8CD	024665						
LIVEV	024711						
PLDTOD	025303						
PHINTV	025433						
ELISEV	025445						
RTSTV	026211						
SETCIV	026243						
SETMIV	026328						
SMFTV	026412						
TAB-IV	026435						
XAKISV	026731						
CHEKIT	027141						
LSJPOL	030335						
WGMTR	030526						
CRPOLC	030563						
CRPOLC	030762						
SID	031635						
VLV	033013						
SOAT	033021						
INBJTC	033066						

IOCMK	033172
OUTPTC	033272
KOWALL	033374
KODER	033443
SIOB	034651
JARUIN	036231
SYSTEM	037513
ASCI	040640
MDID	040667
DXDYV	042462
AICRT3	043032
KRAKER	044736
LOGF	045777
DAES	046002
DALEX	046017
OUTPTS	046074
ENDFIL	046160
APLOTV	046234
CNTTBM	045405
CTLAV	046515
ERBLNV	046702
ERRVLV	047012
FORHV	047117
FRAMEV	047272
GRIDIV	047477
HOLDIV	050516
LINRV	050545
LXVV	051162
NASORT	051212
NOVLNV	051276
NXV	051635
OPMSGV	052036
PLJIND	052071
PLJTV	052344
POIVTV	052466
SCERRV	052634
VC4ARV	052656
VECTRV	053223
XBDVV	053252
XC4ARV	053307
XSCALV	053331
LA3LV	053534
FN3CDV	054173
CATCDC	054335
FN3KXV	054457
NOJTV	054506
PAGE4C	054607
SCOUTV	054745

REFERENCES

-----UNSATISFIED EXTERNALS-----

TOO FEW CONSTANTS FOR UNSUBSCRIPTED ARRAY
 ERROR NUMBER 0149 DETECTED BY INPJN AT ADDRESS 036654
 CALLED FROM P217 AT 002260

TOO FEW CONSTANTS FOR UNSUBSCRIPTED ARRAY
 ERROR NUMBER 0349 DETECTED BY INPJN AT ADDRESS 036654
 CALLED FROM P217 AT 002260

BROADSIDE LOCK, PARTICLE ONLY, 600 LB THRUST, AT NIGHT

```
TIME AFTER IGNITION-TIG(SEC)      1.0000
EXIT PACH NUMBER-ME                0.0000
LOOK ANGLE=AL(DEG)                 0.00
VERT. POLAR ANGLE=BPHI(DEG)        0.00
VELOCITY VECTOR-VR(VPS)            25000.00
ROCKET VELOCITY-VR(FPS)            20000.00
EFFECTIVE EARTH TEMP-TE(K)         1.3300
GAUSSIAN TEMPERATURE-TC(K)         1258.00
CHAMBER TEMP-TC(K)                 1258.00
VELOCITY VECTOR-VX(VPS)            90.00
VELOCITY VECTOR-VY(VPS)            90.00
AZIMUTHAL ANGLE=XI(USEC)           90.00
ALTITUDE-HO(FT)                    500000
LIP ANGLE=LPI(DEG)                  15.0000
NOZZLE EXIT RADIUS=RE(CM)           5.7000
NOZZLE EXIT AREA=AE(SQCM)          1.0000
```

```

IUD = 0 (1 FOR DAY, 5 FOR NIGHT)
ISL = 1 (1-1 FOR PARTICLES) 0-2 FOR GAS ONLY; NEGATIVE TO SKIP INITIAL CALCULATIONS)
IYT = 20 (APPROXIMATE NUMBER OF INTEGRATION STEPS THROUGH LAYER)
LINES = 30 (APPROXIMATE NUMBER OF LINES OF SIGHT IN ONE DIMENSION)
ICR = 5 (NUMBER OF ENTRIES IN FIRING HISTORY TABLE)
IPLT = 1 (0 FOR NO PLOT, 1 FOR ISOINTENSITY PLOTS)

```

TIME	DENSITY	TIME	DENSITY
0.1000	0.	2.0000	1.1000-03
			10.0000 0.

TOTAL EMISSIONS

TEMPERATURE	PARTICLE RADIUS (CM $\times 10^4$)				
	4.460	4.600	4.740	4.880	5.020
99.0	.6195	.5824	.5472	.5138	.4820
150.0	.41620	.39400	.37101	.34762	.32391
200.0	.31945	.30245	.28476	.26642	.24786
250.0	.26140	.24732	.23302	.21801	.20282
300.0	.22370	.21097	.19952	.18574	.17280
350.0	.19564	.18425	.17388	.16268	.15135
400.0	.17333	.16307	.15340	.14339	.13266
450.0	.15200	.14282	.13403	.12475	.11535
500.0	.13230	.12354	.11519	.10633	.09692
600.0	.10594	.11924	.11228	.10505	.09773
700.0	.11213	.10413	.09696	.09352	.08701
800.0	.11177	.08614	.08172	.08488	.07219
1000.0	.08326	.06261	.07778	.07277	.06770
1200.0	.07756	.07344	.06716	.06470	.05949
1400.0	.07067	.06832	.06300	.05824	.05383
1700.0	.06335	.05990	.05648	.05284	.04916
2000.0	.05723	.05511	.05191	.04857	.04518
2500.0	.05243	.04944	.04621	.04373	.04163
3000.0	.04735	.04448	.04121	.03949	.03674
4000.0	.03368	.03082	.02781	.02643	.02387
4900.0	.02781	.02466	.02147	.02017	.01761

BAND DATA, (0.0 TO 14.0)

MISSIVITY IN BAND

TEMP (K)	PLANK FUNC. (W/CM2/SR)	PARTICLE RADIUS (CM X 10E4)	2.660	2.000	1.330	.690
99.0	1.2463E-06	4.660	3.332	2.660	1.330	.690
150.0	7.3650E-05	1.7608	.16672	.15696	.14685	.13594
200.0	6.0061E-04	1.6254	.15419	.14516	.13581	.12635
250.0	2.2775E-03	1.5971	.15056	.14174	.13261	.12338
300.0	5.4808E-03	1.5613	.14302	.13816	.13318	.12810
350.0	1.0038E-02	1.5411	.14621	.13765	.12873	.11981
400.0	1.7685E-02	1.5300	.14487	.13639	.12760	.11871
450.0	2.6500E-02	1.5193	.14307	.13469	.12601	.11723
500.0	3.6897E-02	1.4919	.14116	.13290	.12501	.11630
600.0	8.3531E-02	1.4650	.13850	.13028	.12183	.11340
800.0	1.2370E-01	1.4374	.13499	.12671	.11822	.10967
1000.0	1.9236E-01	1.4061	.13098	.12254	.11401	.10509
1200.0	2.6768E-01	1.3726	.12661	.11817	.10961	.10109
1400.0	3.4625E-01	1.3373	.12193	.11349	.10493	.09643
1700.0	4.6794E-01	1.2919	.11661	.10817	.09961	.09109
2000.0	5.9247E-01	1.2468	.11061	.10217	.09361	.08509
2500.0	8.0365E-01	1.1615	.10338	.09483	.08628	.07775
3200.0	1.1034E+00	1.0594	.09594	.08739	.07884	.07031
6001.0	2.3194E+00	1.4543	.13789	.12981	.12145	.11299

PARTICLE TEMPERATURE HISTORIES (K)

TIME (SEC)	PARTICLE RADIUS (CM X 10E4)	2.660	2.000	1.330	.690
0.000	4.660	3.332	2.660	1.330	.690
0.002	2323.0	2323.0	2323.0	2323.0	2323.0
	2264.1	2264.1	2264.1	2264.1	2264.1
.614	21.7.5	2064.6	2030.4	1974.9	1893.6
.092	1759.9	1724.3	1681.2	1617.8	1534.7
.145	1364.6	1326.4	1273.2	1214.4	1150.1
.336	1032.7	1001.1	957.5	901.7	833.3
.691	791.2	765.5	727.9	680.4	631.7
1.296	622.2	595.7	562.6	534.6	489.1
2.266	497.0	473.3	448.6	426.4	393.4
3.744	4.9.3	392.2	369.8	336.3	319.1
5.977	329.4	322.1	311.9	296.7	272.1
8.972	292.6	281.4	268.0	242.9	216.4
13.192	241.9	225.5	213.7	213.7	211.3
18.540	212.4	210.7	205.8	204.6	202.2
26.391	222.8	201.2	195.5	197.7	195.7
36.036	190.4	189.1	187.4	185.9	184.2
48.378	174.7	173.5	172.3	170.9	169.5
63.692	154.8	154.0	153.1	152.1	151.0
83.191	130.2	129.3	128.7	128.1	127.5
100.000	111.0	110.0	109.0	108.0	107.0

PARTICLE RADIUS (CMX10E4) 4.663 4.000 3.332 2.660 2.000 1.330 .690
 SOLIDIFICATION TIME (SEC) .071 .050 .040 .028 .017
 EQUILIBRIUM TEMPERATURE(K) 199.92 199.92 199.92 199.92 199.92 199.92

ROTATIONAL RADIATION FROM WATER MOLECULES.
 AVAILABLE BAND ENERGY = ET (10E-20 = SEC/MOLEC). BAND RADIATIVE RATE = J (10E-20 W/SR MOLEC)

T(K)	ET	J
100.0	2.724E-04	3.333E-25
400.0	9.408E-15	4.554E-07
800.0	6.744E-10	2.177E-03
1200.0	3.223E-06	5.185E-02
1600.0	2.437E-04	2.807E-01
2000.0	3.447E-03	5.136E-01
2400.0	2.091E-02	1.208E+00
2700.0	5.796E-02	2.586E+00
3000.0	1.327E-01	3.634E+00
3300.0	2.638E-01	4.808E+00
3400.0	4.719E-01	6.070E+00
3900.0	7.776E-01	7.384E+00
4200.0	1.200E+00	8.715E+00
4500.0	1.758E+00	1.004E+01
5000.0	3.031E+00	1.216E+01
5500.0	4.781E+00	1.414E+01
6000.0	7.049E+00	1.593E+01
7000.0	1.322E+01	1.896E+01
9000.0	3.204E+01	2.343E+01
11000.0	5.836E+01	2.715E+01

RADIATIVE DECAY IN BAND

TIME(SEC)	J	J(FIT)	ET	T(K)
0.0000	3.333E-25	5.541E+00	2.724E-04	1.000E+02
.00241	5.071E+00	4.109E+00	3.072E-01	3.362E+03
.00665	6.632E+00	3.036E+00	6.144E-01	3.740E+03
.01004	7.837E+00	7.885E+00	9.215E-01	4.002E+03
.01299	8.783E+00	8.572E+00	1.229E+00	4.215E+03
.01567	9.511E+00	9.193E+00	1.536E+00	4.380E+03
.01815	1.018E+01	9.775E+00	1.843E+00	4.535E+03
.02050	1.069E+01	1.032E+01	2.150E+00	4.654E+03
.02273	1.121E+01	1.081E+01	2.457E+00	4.725E+03
.02486	1.172E+01	1.134E+01	2.765E+00	4.895E+03
.02691	1.221E+01	1.181E+01	3.072E+00	5.012E+03
.02888	1.256E+01	1.227E+01	3.379E+00	5.109E+03
.03080	1.290E+01	1.272E+01	3.686E+00	5.187E+03
.03267	1.325E+01	1.316E+01	3.993E+00	5.275E+03

03449	1.340E+01	1.354E+01	4.301E+00	5.163E+01
03627	1.394E+01	1.399E+01	4.608E+00	5.151E+03
03630	1.424E+01	1.437E+01	4.915E+00	5.530E+03
03971	1.449E+01	1.473E+01	5.222E+00	5.537E+03
04138	1.473E+01	1.514E+01	5.529E+00	5.665E+03
04303	1.497E+01	1.556E+01	5.836E+00	5.733E+03
0.00	0.00	0.00	0.00	0.00
0.00	1487.44	0.	0.	0.
0.00	3123.62	0.	0.	0.
0.00	4923.42	0.	0.	0.
0.00	6973.21	0.	0.	0.
0.00	9080.97	0.	0.	0.
0.00	11476.50	0.	0.	0.
0.00	14111.59	0.	0.	0.
0.00	17010.19	0.	0.	0.
0.00	25198.65	0.	0.	0.
0.00	23795.95	0.	0.	0.
0.00	27563.99	0.	0.	0.
0.00	3187.83	0.	0.	0.
0.00	36476.05	0.	0.	0.
0.00	41611.09	0.	0.	0.
0.00	47259.64	0.	0.	0.
0.00	53473.04	0.	0.	0.
0.00	60307.79	0.	0.	0.
0.00	67826.01	0.	0.	0.
0.00	76096.05	0.	0.	0.
0.00	85193.09	0.	0.	0.
0.00	95199.84	0.	0.	0.
0.00	10627.26	0.	0.	0.
0.00	118315.43	0.	0.	0.
0.00	131634.41	0.	0.	0.
0.00	146285.29	0.	0.	0.
0.00	162401.26	0.	0.	0.
0.00	18128.82	0.	0.	0.
0.00	19929.14	0.	0.	0.
0.00	221079.50	0.	0.	0.
0.00	0.00	0.	0.	0.
0.00	1352.22	1352.22	0.	0.
0.00	1352.22	2839.66	0.	0.
0.00	1352.22	4475.84	0.	0.
0.00	1352.22	6275.64	0.	0.
0.00	1352.22	8255.42	0.	0.
0.00	1352.22	10433.18	0.	0.
0.00	1352.22	12828.72	0.	0.
0.00	1352.22	15463.81	0.	0.
0.00	1352.22	18362.41	0.	0.
0.00	1352.22	21559.87	0.	0.
0.00	1352.22	25058.17	0.	0.
0.00	1352.22	28316.21	0.	0.
0.00	1352.22	3160.04	0.	0.
0.00	1352.22	37828.27	0.	0.
0.00	1352.22	42963.31	0.	0.
0.00	1352.22	48611.86	0.	0.
0.00	1352.22	54825.25	0.	0.
0.00	1352.22	61440.00	0.	0.
0.00	1352.22	69174.22	0.	0.

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1352.22	77448.26	0.
1352.22	86542.31	0.
1352.22	96532.05	0.
1352.22	107559.48	0.
1352.22	119447.64	0.
1352.22	132984.63	0.
1352.22	147637.51	0.
1352.22	163753.47	0.
1352.22	181461.04	0.
1352.22	200981.36	0.
2839.66	0.00	0.
2839.66	1352.22	0.
2839.66	2839.66	0.
2839.66	4475.84	0.
2839.66	6275.64	0.
2839.66	8255.42	0.
2839.66	107559.48	0.
2839.66	12828.72	0.
2839.66	15463.81	0.
2839.66	18362.41	0.
2839.66	21550.87	0.
2839.66	25059.17	0.
2839.66	2839.66	0.
2839.66	33160.04	0.
2839.66	37828.27	0.
2839.66	42963.31	0.
2839.66	48611.86	0.
2839.66	54825.26	0.
2839.66	61178.22	0.
2839.66	69178.22	0.
2839.66	77448.26	0.
2839.66	86542.31	0.
2839.66	96532.05	0.
2839.66	107559.48	0.
2839.66	119447.64	0.
2839.66	132984.63	0.
2839.66	147637.51	0.
2839.66	163753.47	0.
2839.66	181461.04	0.
2839.66	200981.36	0.
4475.84	0.00	0.
4475.84	1352.22	0.
4475.84	2839.66	0.
4475.84	4475.84	0.
4475.84	6275.64	0.
4475.84	8255.42	0.
4475.84	107559.48	0.
4475.84	12828.72	0.
4475.84	15463.81	0.
4475.84	18362.41	0.
4475.84	21550.87	0.
4475.84	25059.17	0.
4475.84	2839.66	0.
4475.84	33160.04	0.
4475.84	37828.27	0.
4475.84	42963.31	0.

4475.84	48611.86	0.
4475.84	54825.26	0.
4475.84	61460.00	0.
4475.84	69178.22	0.
4475.84	77448.26	0.
4475.84	86545.31	0.
4475.84	96552.05	0.
4475.84	107559.48	0.
4475.84	119467.64	0.
4475.84	132986.63	0.
4475.84	147637.51	0.
4475.84	163753.47	0.
4475.84	181481.04	0.
4475.84	200981.36	0.
6275.54	0.00	0.
6275.64	1352.22	0.
6275.64	2839.66	0.
6275.64	4475.84	0.
6275.64	6275.64	0.
6275.64	8255.42	0.
6275.54	10433.18	0.
6275.54	12828.72	0.
6275.54	15463.81	0.
6275.54	18362.41	0.
6275.64	21550.87	0.
6275.54	25058.17	0.
6275.54	28316.21	0.
6275.54	33160.04	0.
6275.64	37828.27	0.
6275.54	42931.31	0.
6275.54	48611.86	0.
6275.54	54825.26	0.
6275.54	61460.00	0.
6275.54	69178.22	0.
6275.54	77443.26	0.
6275.54	86545.31	0.
6275.64	96552.05	0.
6275.64	107559.48	0.
6275.54	119467.64	0.
6275.54	132986.63	0.
6275.54	147637.51	0.
6275.54	163753.47	0.
6275.64	181481.04	0.
6275.54	200951.36	0.
8255.42	1.00	0.
8255.42	1352.22	0.
8255.42	2837.66	0.
8255.42	4475.84	0.
8255.42	6275.64	0.
8255.42	8255.42	0.
8255.42	10433.18	0.
8255.42	12828.72	0.
8255.42	15463.81	0.
8255.42	18362.41	0.
8255.42	21550.87	0.
8255.42	25058.17	0.

8295.42	28914.21	0.
8295.42	33160.04	0.
8295.42	37828.27	0.
8295.42	42943.31	0.
8295.42	48611.86	0.
8295.42	54825.26	0.
8295.42	61440.00	0.
8295.42	69178.22	0.
8295.42	77448.26	0.
8295.42	86545.31	0.
8295.42	96552.05	0.
8295.42	107559.48	0.
8295.42	119442.44	0.
8295.42	132986.63	0.
8295.42	147637.51	0.
8295.42	163753.47	0.
8295.42	181481.04	0.
8295.42	200981.36	0.
10433.18	0.00	2.514E-04
10433.18	1352.22	2.437E-06
10433.18	2839.66	0.
10433.18	4479.84	0.
10433.18	6273.64	0.
10433.18	8255.42	0.
10433.18	11433.18	0.
10433.18	12828.72	0.
10433.18	15463.81	0.
10433.18	18362.41	0.
10433.18	21590.87	0.
10433.18	25058.17	0.
10433.18	28914.21	0.
10433.18	33160.04	0.
10433.18	37828.27	0.
10433.18	42943.31	0.
10433.18	48611.86	0.
10433.18	54825.26	0.
10433.18	61440.00	0.
10433.18	69178.22	0.
10433.18	77448.26	0.
10433.18	86545.31	0.
10433.18	96552.05	0.
10433.18	107559.48	0.
10433.18	119442.44	0.
10433.18	132986.63	0.
10433.18	147637.51	0.
10433.18	163753.47	0.
10433.18	181481.04	0.
10433.18	200981.36	0.
12828.72	0.00	2.514E-04
12828.72	1352.22	1.787E-06
12828.72	2839.66	1.456E-06
12828.72	4479.84	1.438E-06
12828.72	6273.64	0.
12828.72	8255.42	0.
12828.72	10433.18	0.
12828.72	12828.72	0.

18362.41	4475.84	2.539E-06
18362.41	6275.64	2.154E-06
18362.41	8255.42	5.224E-07
18362.41	13433.18	0.
18362.41	12828.72	0.
18362.41	15443.81	0.
18362.41	18362.41	0.
18362.41	21550.87	0.
18362.41	25058.17	0.
18362.41	28916.21	0.
18362.41	33160.04	0.
18362.41	37828.27	0.
18362.41	42963.31	0.
18362.41	48611.86	0.
18362.41	54825.24	0.
18362.41	61660.00	0.
18362.41	69178.22	0.
18362.41	77448.26	0.
18362.41	86545.31	0.
18362.41	96552.05	0.
18362.41	107559.48	0.
18362.41	119667.64	0.
18362.41	132986.63	0.
18362.41	147637.51	0.
18362.41	163753.47	0.
18362.41	181491.04	0.
18362.41	200881.16	0.
21550.87	0.00	1.375E-05
21550.87	1352.22	1.354E-05
21550.87	2832.66	1.285E-05
21550.87	4475.84	4.657E-06
21550.87	6275.64	1.587E-06
21550.87	8255.42	1.328E-06
21550.87	13433.18	3.063E-07
21550.87	12828.72	0.
21550.87	15463.81	0.
21550.87	18362.41	0.
21550.87	21550.87	0.
21550.87	25051.17	0.
21550.87	28916.21	0.
21550.87	33160.04	0.
21550.87	37828.27	0.
21550.87	42963.31	0.
21550.87	48611.86	0.
21550.87	54825.24	0.
21550.87	61660.00	0.
21550.87	69178.22	0.
21550.87	77448.26	0.
21550.87	86545.31	0.
21550.87	96552.05	0.
21550.87	107559.48	0.
21550.87	119667.64	0.
21550.87	132986.63	0.
21550.87	147637.51	0.
21550.87	163753.47	0.
21550.87	181491.04	0.

2150.87	200981.36	0	ISOINTENSITY LINE FOR 1.00E-05W/CH2/SR IS NEAR (1.55E-04, 2.84E+03)
25098.17	0.00	2.059E-05	
25098.17	1352.22	2.034E-05	
25098.17	2839.66	9.325E-06	
25098.17	4475.64	6.653E-06	
25098.17	6275.64	3.022E-06	
25098.17	8255.42	1.094E-06	
25098.17	10433.18	4.230E-07	
25098.17	12628.72	1.996E-07	
25098.17	15463.61	0.	
25098.17	18362.41	0.	
25098.17	21553.67	0.	
25098.17	25058.17	0.	
25098.17	28916.21	0.	
25098.17	33160.34	0.	
25098.17	37823.27	0.	
25098.17	42963.31	0.	
25098.17	48611.86	0.	
25098.17	54825.26	0.	
25098.17	61640.00	0.	
25098.17	69178.22	0.	
25098.17	77448.26	0.	
25098.17	86545.31	0.	
25098.17	96552.09	0.	
25098.17	107559.48	0.	
25098.17	119667.64	0.	
25098.17	132986.63	0.	
25098.17	147637.51	0.	
25098.17	163753.47	0.	
25098.17	181481.04	0.	
25098.17	200981.36	0.	
28916.21	0.00	2.029E-05	
28916.21	1352.22	2.011E-05	
28916.21	2839.66	1.404E-05	
28916.21	4475.64	1.325E-05	
28916.21	6275.64	5.342E-06	
28916.21	8255.42	2.195E-06	
28916.21	10433.18	7.351E-07	
28916.21	12628.72	6.103E-07	
28916.21	15463.61	1.298E-07	
28916.21	18362.41	0.	
28916.21	21553.67	0.	
28916.21	25058.17	0.	
28916.21	28916.21	0.	
28916.21	33160.34	0.	
28916.21	37823.27	0.	
28916.21	42963.31	0.	
28916.21	48611.86	0.	
28916.21	54825.26	0.	
28916.21	61640.00	0.	
28916.21	69178.22	0.	
28916.21	77448.26	0.	
28916.21	86545.31	0.	
28916.21	96552.09	0.	
28916.21	107559.48	0.	

28916.21	119607.64	0.	
28916.21	132986.63	0.	
28916.21	147637.51	0.	
28916.21	163753.47	0.	
28916.21	181481.04	0.	
28916.21	201981.36	0.	
33160.04	0.00	1.953E-05	
33160.04	1352.22	1.843E-05	
33160.04	2839.66	1.465E-05	
33160.04	4475.64	1.724E-05	
33160.04	6225.64	6.528E-06	
33160.04	8255.42	4.082E-06	
33160.04	10433.18	1.505E-06	
33160.04	12828.72	4.195E-07	
33160.04	15463.81	4.056E-07	
33160.04	18362.41	6.397E-08	
33160.04	21350.87	0.	
33160.04	25059.17	0.	
33160.04	28916.21	0.	
33160.04	33160.04	0.	
33160.04	37828.27	0.	
33160.04	42963.31	0.	
33160.04	48611.86	0.	
33160.04	54825.26	0.	
33160.04	61667.00	0.	
33160.04	69178.22	0.	
33160.04	77449.26	0.	
33160.04	86545.31	0.	
33160.04	96552.05	0.	
33160.04	107552.48	0.	
33160.04	119607.64	0.	
33160.04	132986.63	0.	
33160.04	147637.51	0.	
33160.04	163753.47	0.	
33160.04	181481.04	0.	
33160.04	201981.36	0.	
37826.27	0.00	1.448E-05	
37826.27	1352.22	1.442E-05	
37826.27	2839.66	1.420E-05	
37826.27	4475.64	1.380E-05	
37826.27	6225.64	7.219E-06	
37826.27	8255.42	6.651E-06	
37826.27	10433.18	3.061E-06	
37826.27	12828.72	1.316E-06	
37826.27	15463.81	4.204E-07	
37826.27	18362.41	2.675E-07	
37826.27	21350.87	8.043E-08	
37826.27	25059.17	0.	
37826.27	28916.21	0.	
37826.27	33160.04	0.	
37826.27	37828.27	0.	
37826.27	42963.31	0.	
37826.27	48611.86	0.	
37826.27	54825.26	0.	
37826.27	61667.00	0.	
37826.27	69178.22	0.	

37828.27	77449.26	0.
37828.27	86545.31	0.
37828.27	94552.05	0.
37828.27	107559.48	0.
37828.27	119667.64	0.
37828.27	132986.63	0.
37828.27	147637.51	0.
37828.27	163757.47	0.
37828.27	181451.04	0.
37828.27	200931.36	0.
42943.31	0.00	1.184E-05
42943.31	1352.22	1.180E-05
42943.31	2839.66	1.164E-05
42943.31	4475.84	1.138E-05
42943.31	6275.64	9.759E-06
42943.31	8255.42	5.650E-06
42943.31	10433.18	2.347E-06
42943.31	12828.72	2.388E-06
42943.31	15443.81	4.795E-07
42943.31	18362.41	2.794E-07
42943.31	21551.67	6.185E-08
42943.31	25093.17	5.190E-08
42943.31	28916.21	0.
42943.31	33160.04	0.
42943.31	37828.27	0.
42943.31	42963.31	0.
42943.31	48611.86	0.
42943.31	54825.26	0.
42943.31	61669.00	0.
42943.31	69178.22	0.
42943.31	77449.26	0.
42943.31	86545.31	0.
42943.31	94552.05	0.
42943.31	107559.48	0.
42943.31	119667.64	0.
42943.31	132986.63	0.
42943.31	147637.51	0.
42943.31	163757.47	0.
42943.31	181451.04	0.
42943.31	200931.36	0.
48611.86	0.00	9.941E-06
48611.86	1352.22	9.935E-06
48611.86	2839.66	9.173E-06
48611.86	4475.84	9.307E-06
48611.86	6275.64	8.320E-06
48611.86	8255.42	6.462E-06
48611.86	10433.18	3.893E-06
48611.86	12828.72	2.022E-06
48611.86	15443.81	6.529E-07
48611.86	18362.41	4.464E-07
48611.86	21551.67	3.414E-07
48611.86	25093.17	3.932E-08
48611.86	28916.21	3.214E-09
48611.86	33160.04	0.
48611.86	37828.27	0.
48611.86	42963.31	0.

48611.86	48611.86	0.
48611.86	54825.26	0.
48611.86	61660.00	0.
48611.86	69178.22	0.
48611.86	77448.26	0.
48611.86	86545.31	0.
48611.86	96552.05	0.
48611.86	107559.48	0.
48611.86	119667.64	0.
48611.86	132966.63	0.
48611.86	147637.51	0.
48611.86	163751.47	0.
48611.86	181461.04	0.
48611.86	200921.36	0.
54825.26	9.00	7.08E-06
54825.26	1352.22	7.073E-06
54825.26	2839.66	7.024E-06
54825.26	4475.64	6.904E-06
54825.26	6225.64	6.756E-06
54825.26	8255.42	6.477E-06
54825.26	10433.18	6.445E-06
54825.26	12828.72	6.668E-06
54825.26	15463.91	6.479E-06
54825.26	18362.41	6.422E-07
54825.26	21550.87	2.944E-07
54825.26	25059.17	1.205E-07
54825.26	28916.71	2.647E-06
54825.26	33160.04	2.230E-08
54825.26	37828.27	0.
54825.26	42963.31	0.
54825.26	48611.86	0.
54825.26	54825.26	0.
54825.26	61660.00	0.
54825.26	69178.22	0.
54825.26	77448.26	0.
54825.26	86545.31	0.
54825.26	96552.05	0.
54825.26	107559.48	0.
54825.26	119667.64	0.
54825.26	132966.63	0.
54825.26	147637.51	0.
54825.26	163751.47	0.
54825.26	181461.04	0.
54825.26	200921.36	0.
61660.00	0.00	5.301E-06
61660.00	1352.22	5.262E-06
61660.00	2839.66	5.262E-06
61660.00	4475.64	5.205E-06
61660.00	6225.64	5.115E-06
61660.00	8255.42	4.764E-06
61660.00	10433.18	4.769E-06
61660.00	12828.72	2.308E-06
61660.00	15463.91	2.323E-06
61660.00	18362.41	9.919E-07
61660.00	21550.87	1.252E-07
61660.00	25059.17	1.746E-07

61660.00	28916.21	P.491E-08
61660.00	33160.04	1.836E-08
61660.00	37828.27	1.547E-08
61660.00	42963.31	0.
61660.00	48611.86	0.
61660.00	54825.26	0.
61660.00	61660.00	0.
61660.00	69178.22	0.
61660.00	77446.26	0.
61660.00	86545.31	0.
61660.00	96552.05	0.
61660.00	107559.48	0.
61660.00	119667.64	0.
61660.00	132986.63	0.
61660.00	147637.51	0.
61660.00	163753.47	0.
61660.00	181481.04	0.
61660.00	200981.36	0.
6178.22	0.00	3.954E-06
6178.22	152.22	3.940E-06
6178.22	2839.66	3.931E-06
6178.22	4475.84	3.897E-06
6178.22	6273.64	3.848E-06
6178.22	8259.42	3.763E-06
6178.22	10433.18	3.656E-06
6178.22	12828.72	3.557E-06
6178.22	15463.81	3.468E-06
6178.22	18362.41	3.388E-06
6178.22	21550.87	3.316E-07
6178.22	25058.17	3.247E-07
6178.22	28916.21	3.185E-08
6178.22	33160.04	3.129E-08
6178.22	37828.27	3.076E-08
6178.22	42963.31	3.025E-08
6178.22	48611.86	2.976E-08
6178.22	54825.26	2.929E-08
6178.22	61660.00	2.884E-08
6178.22	69178.22	2.840E-08
6178.22	77446.26	2.797E-08
6178.22	86545.31	2.755E-08
6178.22	96552.05	2.714E-08
6178.22	107559.48	2.674E-08
6178.22	119667.64	2.635E-08
6178.22	132986.63	2.596E-08
6178.22	147637.51	2.558E-08
6178.22	163753.47	2.520E-08
6178.22	181481.04	2.483E-08
6178.22	200981.36	2.447E-08
77446.26	0.00	3.019E-06
77446.26	1352.22	3.015E-06
77446.26	2639.66	3.005E-06
77446.26	4475.84	2.994E-06
77446.26	6273.64	2.984E-06
77446.26	8255.42	2.974E-06
77446.26	10433.18	2.964E-06
77446.26	12828.72	2.954E-06

77448.26	15463.81	2.270E+06
77448.26	16362.41	1.545E+06
77448.26	21550.87	6.270E+07
77448.26	25058.17	3.323E+07
77448.26	28916.21	1.831E+07
77448.26	33160.64	5.700E+08
77448.26	37829.27	4.188E+08
77448.26	42963.31	1.341E+08
77448.26	48411.86	7.142E+09
77448.26	54825.26	0.
77448.26	61660.00	0.
77448.26	69179.22	0.
77448.26	77440.26	0.
77448.26	86545.31	0.
77448.26	96552.05	0.
77448.26	107553.48	0.
77448.26	119667.64	0.
77448.26	132986.63	0.
77448.26	147637.51	0.
77448.26	163753.47	0.
77448.26	181441.04	0.
77448.26	200931.36	0.
86545.31	0.00	2.379E+06
86545.31	1352.22	2.377E+06
86545.31	2839.66	2.370E+06
86545.31	4475.84	2.356E+06
86545.31	6212.64	2.332E+06
86545.31	8259.42	2.308E+06
86545.31	10433.10	2.259E+06
86545.31	12828.72	2.208E+06
86545.31	15403.81	1.973E+06
86545.31	18362.41	1.647E+06
86545.31	21550.87	1.146E+06
86545.31	25058.17	6.321E+07
86545.31	28916.21	2.365E+07
86545.31	33160.64	1.285E+07
86545.31	37829.27	5.356E+08
86545.31	42963.31	3.150E+08
86545.31	48411.86	9.185E+09
86545.31	54825.26	4.749E+09
86545.31	61660.00	0.
86545.31	69179.22	0.
86545.31	77440.26	0.
86545.31	86545.31	0.
86545.31	96552.05	0.
86545.31	107553.48	0.
86545.31	119667.64	0.
86545.31	132986.63	0.
86545.31	147637.51	0.
86545.31	163753.47	0.
86545.31	181441.04	0.
86545.31	200931.36	0.
96552.05	0.00	2.332E+06
96552.05	1352.22	2.038E+06
96552.05	2839.66	2.125E+06
96552.05	4475.84	2.016E+06

96552.05	6225.64	2.301E-06
96552.05	8259.42	1.228E-06
96552.05	17433.18	1.498E-06
96552.05	12828.72	1.358E-06
96552.05	15463.61	1.570E-06
96552.05	18362.41	1.342E-06
96552.05	21352.67	1.332E-06
96552.05	25053.17	5.364E-07
96552.05	28916.21	4.435E-07
96552.05	33150.04	1.751E-07
96552.05	37823.27	9.395E-08
96552.05	42983.31	3.789E-08
96552.05	48631.86	2.232E-08
96552.05	54825.26	6.756E-09
96552.05	61660.00	1.130E-09
96552.05	59178.22	0.
96552.05	77449.26	0.
96552.05	96545.31	0.
96552.05	96552.05	0.
96552.05	107559.46	0.
96552.05	119667.64	0.
96552.05	132956.63	0.
96552.05	147637.51	0.
96552.05	163753.47	0.
96552.05	181451.04	0.
ISOTENSITY LINE FOR 1.00E-06W/CM2/SR IS NEAR (7.74E+04, 2.16E+04)		
107559.48	0.00	1.757E-06
107559.48	1352.22	1.750E-06
107559.48	2839.86	1.753E-06
107559.48	4475.84	1.747E-06
107559.48	6225.64	1.732E-06
107559.48	8259.42	1.722E-06
107559.48	12433.18	1.702E-06
107559.48	12828.72	1.535E-06
107559.48	15463.61	1.465E-06
107559.48	18362.41	1.281E-06
107559.48	21350.87	1.141E-06
107559.48	25053.17	7.574E-07
107559.48	28916.21	4.270E-07
107559.48	33150.04	2.244E-07
107559.48	37828.27	1.382E-07
107559.48	42943.31	9.274E-08
107559.48	48444.86	2.460E-08
107559.48	54825.26	9.074E-09
107559.48	61660.00	4.106E-09
107559.48	69178.22	2.163E-09
107559.48	77449.26	0.
107559.48	86545.31	0.
107559.48	96552.05	0.
107559.48	107559.48	0.
107559.48	119667.64	0.
107559.48	132956.63	0.
107559.48	147637.51	0.
107559.48	163753.47	0.
107559.48	181451.04	0.
107559.48	0.00	1.102E-06

119A37.64	1352.22	1.401E-06
119A37.64	2839.66	1.399E-06
119A37.64	4475.64	1.395E-06
119A37.64	6275.64	1.388E-06
119A37.64	8255.42	1.378E-06
119A37.64	10433.18	1.365E-06
119A37.64	12626.72	1.346E-06
119A37.64	15463.81	1.322E-06
119A37.64	18362.41	1.153E-06
119A37.64	21550.67	9.965E-07
119A37.64	25038.17	8.503E-07
119A37.64	28916.21	5.257E-07
119A37.64	33160.04	3.263E-07
119A37.64	37828.27	1.684E-07
119A37.64	42963.31	7.387E-08
119A37.64	48611.86	4.804E-08
119A37.64	54825.26	1.972E-08
119A37.64	61660.00	1.140E-08
119A37.64	69178.22	3.434E-09
119A37.64	77448.26	3.454E-09
119A37.64	86545.31	0.
119A37.64	96552.05	0.
119A37.64	107559.48	0.
119A37.64	119667.64	0.
119A37.64	132986.63	0.
119A37.64	147637.51	0.
119A37.64	163753.47	0.
119A37.64	181461.04	0.
132986.63	0.00	1.155E-06
132986.63	1352.22	1.154E-06
132986.63	2839.66	1.153E-06
132986.63	4475.64	1.150E-06
132986.63	6275.64	1.103E-06
132986.63	8255.42	1.097E-06
132986.63	10433.18	1.087E-06
132986.63	12626.72	1.075E-06
132986.63	15463.81	1.047E-06
132986.63	18362.41	9.475E-07
132986.63	21550.67	9.220E-07
132986.63	25058.17	7.614E-07
132986.63	28916.21	7.284E-07
132986.63	33160.04	4.111E-07
132986.63	37828.27	3.237E-07
132986.63	42963.31	1.336E-07
132986.63	48611.86	6.254E-08
132986.63	54825.26	3.449E-08
132986.63	61660.00	1.575E-08
132986.63	69178.22	8.755E-09
132986.63	77448.26	2.362E-09
132986.63	86545.31	9.889E-10
132986.63	96552.05	0.
132986.63	107559.48	0.
132986.63	119667.64	0.
132986.63	132986.63	0.
132986.63	147637.51	0.
132986.63	163753.47	0.

[A] have only $s < 8$ significant digits, the user may compute sigrev, the corrected number of significant digits in the inverse by the formula

$$\text{sigrev} = \text{sigdig} - 8.3 + s$$

To get some idea how accurate sigrev itself was, the difference between sigrev and sigact, the actual average number of significant digits, was computed for many matrices of random numbers. When also all the elements of the matrix were of the same order of magnitude, the results (for a batch of 120 matrices) were:

$$-1.17 < (\text{sigrev} - \text{sigact}) < 0.75$$

However, when the elements of the matrix ranged over three orders of magnitude, the results for a batch of 160 matrices were:

$$-1.77 < (\text{sigrev} - \text{sigact}) < 0.78$$

Both batches contained 100 5 by 5 matrices, the other matrices all being 10 by 10.

A DOUBLE PRECISION substitute for this subprogram is available.

- I. Error Indications - ierror is an INTEGER variable which is used to flag errors.

ierror = -1 indicates both of the following:

1. Neither $[A]^{-1}$ nor $[X]$ could be computed.
2. [A] was either singular* or very nearly so (usually the former.

ierror = 1 indicates that no errors were detected.

Due to round-off errors, a singular matrix will rarely be flagged as singular (by ierror = -1) unless it has a row or column of zeros. However, the value of sigdig will seldom be greater than 1.5 for singular matrices.

Division by zero cannot occur in this subprogram.

No test for floating-point overflow is made in this subprogram.

*When [A] is singular, a system of simultaneous equations $[A][X] = [B]$ will have a solution if and only if [B] can be expressed as a linear combination of the columns of [A].

1P14A1.04 0.00 5.972E+07
1A14A1.04 1352.22 5.231E+07
1A14A1.04 2839.66 5.966E+07
1A14A1.04 4475.84 5.957E+07
1A14A1.04 6275.64 5.943E+07
1A14A1.04 8255.42 5.923E+07
1A14A1.04 10433.18 5.394E+07
1A14A1.04 12328.72 5.654E+07
1A14A1.04 15463.81 5.643E+07
1A14A1.04 18362.41 5.222E+07
1A14A1.04 21550.67 4.987E+07
1A14A1.04 25050.17 4.891E+07
1A14A1.04 28916.21 4.772E+07
1A14A1.04 33145.04 4.312E+07
1A14A1.04 37826.27 3.595E+07
1A14A1.04 42963.31 2.781E+07
1A14A1.04 48611.86 1.790E+07
1A14A1.04 54825.26 8.160E+06
1A14A1.04 61660.00 6.544E+06
1A14A1.04 69178.22 2.414E+06
1A14A1.04 77448.26 1.375E+06
1A14A1.04 86549.31 6.644E+05
1A14A1.04 96542.05 2.355E+05
1A14A1.04 107559.48 1.124E+05
1A14A1.04 119667.64 1.646E+04
ISOMINELITY LINE FOR 3.14E+07W/CM2/SR IS NEAR (1.48E+05, 1.78E+04)

2009A1.36 0.00 4.611E+07
2009A1.36 1352.22 4.611E+07
2009A1.36 2839.66 4.608E+07
2009A1.36 4475.84 4.602E+07
2009A1.36 6275.64 4.593E+07
2009A1.36 8255.42 4.580E+07
2009A1.36 10433.18 4.561E+07
2009A1.36 12828.72 4.535E+07
2009A1.36 15463.81 4.404E+07
2009A1.36 18362.41 4.364E+07
2009A1.36 21550.67 4.300E+07
2009A1.36 25050.17 4.132E+07
2009A1.36 28916.21 3.939E+07
2009A1.36 33160.04 3.715E+07
2009A1.36 37826.27 3.125E+07
2009A1.36 42963.31 2.528E+07
2009A1.36 48611.86 2.063E+07
2009A1.36 54825.26 1.361E+07
2009A1.36 61660.00 5.377E+06
2009A1.36 69178.22 4.337E+06
2009A1.36 77448.26 1.384E+06
2009A1.36 86549.31 5.872E+05
222431.71 0.00 0.

EMCR SUMMARY
ERROR 0049
TIMES 6.02

Appendix C

AUXILIARY LIBRARY ROUTINES

Operating descriptions and listings are presented here for LSQPØL and SID-- two subroutines from the MDAC Library. These were not programmed under this contract, but are used by the FLAME code.

C.1 LSQPØL

Least Squares Polynomial SUBROUTINE

- A. Description - This subroutine subprogram computes the coefficients of the polynomial which best fits (in the least squares sense) a given function of one independent variable.
- B. Use - CALL LSQPØL (x, y, w, r, n, s, m, c) where:
 - x is the REAL array, dimensioned at least n, of the values of the independent variable of the given function.
 - y is the REAL array, dimensioned at least n, of the values of the dependent variable of the given function.
 - w is the REAL array, dimensioned at least n, of the least squares weighting for each point of the given function. (If an unweighted least squares fit is desired, each element of this array must be a 1.)
 - r is the REAL array, dimensioned at least n, which this subprogram sets equal to the residuals at each point of the given function.
 - n is an INTEGER variable (or constant) denoting the number of points defining the given function. This number must be ≤ 50 .
 - s is the REAL variable which this subprogram sets equal to the sum of the squares of the residuals.
 - m is the INTEGER variable (or constant) denoting how many coefficients the least square polynomial will have. That is, it is the degree plus one of the least squares polynomial desired. This number must be input ≤ 10 .

c is the REAL array, dimensioned at least m whose elements are set equal to the coefficients of the least squares polynomial beginning with the constant term.

- C. Support Level - Supported by Programming Systems and Support Branch of Information Processing Systems.
- D. Language Used - FORTRAN.
- E. Availability - On FORTRAN library.
- F. Extent - 4052₈ words.
- G. Timing - Not available.
- H. Restrictions - None.
- I. Error Indications - The least squares problem is mathematically inconsistent unless $m \leq n$, therefore, if $m > n$ the following note is printed and execution of the program is terminated.

LSQPOL ERROR - MORE COEFFICIENTS THAN POINTS

- J. Method - The method used is completely described in "Polynomial Curve Fitting With Constraint" by J. E. L. Peck in SIAM Review, volume 4, number 2, April, 1962.

This subprogram finds the c_i 's such that:

$$s = \sum_{k=1}^n w_k r_k^2 \text{ is a minimum.}$$

$$r_k = y_k - \sum_{i=1}^m c_i x_k^{i-1}$$

- K. Examples - None.

C.2 SID

Single Precision Simultaneous Equation Solution and Matrix Inversion SUBROUTINE

- A. Description - This subroutine subprogram will invert a REAL nonsingular square matrix, [A], and if desired, also solve linear system(s) of simultaneous equations, $[A][X] = [B]$, where [B] may have any number of columns. Every column of [B] must have at least one non-zero element. The dimensions of [A] and [B] are limited only by the available core storage.

B. Use - CALL SID (amat, nrow, maxrow, mxcola, bmat, ncolb, mxcolb, sigdig, ierror, rtemp, itemp, scaleb) where:

1. Input

amat is a REAL two dimensional array which contains [A]. [A] will be replaced by $[A]^{-1}$ during the execution of this subprogram.

nrow is an INTEGER variable or constant which denotes the number of rows in [A]. nrow may be as small as one (1) or as large as available core storage will permit.

maxrow is an INTEGER variable or constant which denotes the maximum number of rows which may be stored in the amat array. Note that a matrix may have fewer rows and/or columns than the array which contains it.

mxcola is an INTEGER variable or constant which denotes the maximum number of columns which may be stored in the amat array.

bmat is a Real array which contains [B] if [B] is present. If [B] is present, the bmat array must have exactly maxrow rows. [B] will be replaced by [X] during the execution of this subprogram. If no [B] is present, bmat may be any variable or constant of any type.

ncolb is an INTEGER variable or constant which denotes the number of columns in [B]. If there is no [B], use ncolb = 0. ncolb may be as small as zero or as large as available core storage will permit.

mxcolb is an INTEGER variable or constant which denotes the maximum number of columns which may be stored in the bmat array. If no [B] is present, mxcolb may have any value.

2. Output

amat will contain $[A]^{-1}$.

bmat will contain [X] if [B] was present.

sigdig is a REAL variable which will be set equal to an estimate of the number of significant digits in the elements of the inverse. However, this estimate is based on the assumption that the elements of [A] are all accurate to eight significant digits. The Restrictions section describes a simple adjustment of sigdig which the user must make when the elements of [A] have fewer than eight significant digits.

ierror is an INTEGER variable which is used to flag errors. Its meaning is explained in Error Indications.

3. Intermediate - The following three arrays are used during calculation for the temporary storage of intermediate results. It makes no difference whether the arrays are one, two, or three dimensional. The three arrays need not all be of the same dimension.

rtemp is a REAL array which contains at least $3 \cdot nrow$ elements.

itemp is an INTEGER array which has at least $3 \cdot nrow$ elements.

scaleb is a REAL array which contains at least $ncolb$ elements.

- C. Support Level - Supported by Programming Systems and Support Branch of Information Processing Systems.
- D. Language Used - MSSD Standard FORTRAN.
- E. Availability - On the FORTRAN library.
- F. Extent - 766 words, 19 words for ALOG10, and 46 words for ALOG.
- G. Timing - Not available.
- H. Restrictions - This subprogram cannot find a solution to the system of simultaneous equations $[A] [X] = [0]$.

An estimate of the average number of significant digits in the elements of the inverse is given by the argument sigdig. (See the Use section.) However, this estimate is only valid when the elements of A have eight significant digits. When the elements of

[A] have only $s < 8$ significant digits, the user may compute sigrev, the corrected number of significant digits in the inverse by the formula

$$\text{sigrev} = \text{sigdig} - 8.3 + s$$

To get some idea how accurate sigrev itself was, the difference between sigrev and sigact, the actual average number of significant digits, was computed for many matrices of random numbers. When also all the elements of the matrix were of the same order of magnitude, the results (for a batch of 120 matrices) were:

$$-1.17 < (\text{sigrev} - \text{sigact}) < 0.75$$

However, when the elements of the matrix ranged over three orders of magnitude, the results for a batch of 160 matrices were:

$$-1.77 < (\text{sigrev} - \text{sigact}) < 0.78$$

Both batches contained 100 5 by 5 matrices, the other matrices all being 10 by 10.

A DOUBLE PRECISION substitute for this subprogram is available.

- I. Error Indications - ierror is an INTEGER variable which is used to flag errors.

ierror = -1 indicates both of the following:

1. Neither $[A]^{-1}$ nor $[X]$ could be computed.
2. $[A]$ was either singular* or very nearly so (usually the former.

ierror = 1 indicates that no errors were detected.

Due to round-off errors, a singular matrix will rarely be flagged as singular (by ierror = -1) unless it has a row or column of zeros. However, the value of sigdig will seldom be greater than 1.5 for singular matrices.

Division by zero cannot occur in this subprogram.

No test for floating-point overflow is made in this subprogram.

*When $[A]$ is singular, a system of simultaneous equations $[A][X] = [B]$ will have a solution if and only if $[B]$ can be expressed as a linear combination of the columns of $[A]$.

- J. Method - The numerical method used is basically the Gauss-Jordan elimination with selection of maximum pivotal elements (full pivoting). A description of the Gauss-Jordan method is contained in Numerical Analysis by Kaiser S. Kunz; McGraw-Hill, 1957.

Special procedures are present to improve accuracy and also to minimize the frequency of overflow and underflow when the magnitudes of any of the elements in the A or B matrices are either large or small.

- K. Example - The coding below will invert a matrix, solve a system of simultaneous equations, and print a warning note if the inverse has less than about one significant digit. The elements of the matrix originally in AMAT are assumed to have an average accuracy of eight significant digits.

```
.  
.   
.   
DIMENSION AMAT (10,10), BMAT(10,1 ), RTEMP(10,3),  
ITEMP(10,3), 1 SCALEB(1)
```

```
.   
.   
.   
NRØW = 6
```

```
.   
.   
.   
CALL SID (AMAT,NRØW, 10, 10, BMAT, 1, 1, SIGDIG, IERRØR,  
1 RTEMP, ITEMP, SCALEB)  
IF (SIGDIG .LT. 2.) WRITE (6,180) SIGDIG  
180 FØRMAT (40H1WARNING - MATRIX INVERSE HAS ØNLY ABOUT,  
1 F5.1, 20H SIGNIFICANT DIGITS.)
```

```

SUBROUTINE SID (A, N, NDRW, NDCOLA, B, M, NDCOLB, SIGDIG, IERROR,
. PIVOT, INDEX, SCALEB )
C
C SID - A SINGLE PRECISION SIMULTANEOUS EQUATION SOLVER, INVERSE
C FINDER, AND DETERMINANT SUBROUTINE
C
  DIMENSION A(NDRW,NDCOLA), B(NDRW,NDCOLB), PIVOT(N,3),
. SCALEB(M), INDEX(N,3)
  DOUBLE PRECISION DRIGP2
  DATA DRIGP2
. / 7378697629483829.D4
C
  EPS = 1.F-3
712 EPS = EPS/2.
  EPSP15= EPS + 1.5
  IF (EPSP15 .NE. 1.5) GO TO 712
  SIGMCH = ALOG10(1.522/EPS)
  RIGPW2 = DRIGP2
  PIVOT(1,1) = 0.
C
C SCALE ROWS
C
  DO 38 I=1,N
    ROWMX = 0.
    DO 28 J=1,N
      IF ((ABS(A(I,J))) .GT. ROWMX) ROWMX = ABS(A(I,J))
28 CONTINUE
      IF ( ROWMX) 29, 750, 29
29 CONTINUE
      ROWMXI = 1. / ROWMX
      DO 32 J=1,N
        AIJ = A(I,J)
        A(I,J) = (A(I,J) * ROWMXI) * BIGPW2
        IF (A(I,J) .EQ. 0.) A(I,J) = (AIJ * BIGPW2) * ROWMXI
32 CONTINUE
        IF (M) 34, 38, 34
34 DO 36 J=1,M
        RIJ = B(I,J)
        B(I,J) = (B(I,J) * ROWMXI) * BIGPW2
        IF (B(I,J) .EQ. 0.) B(I,J) = (RIJ * BIGPW2) * ROWMXI
36 CONTINUE
38 PIVOT(I,2) = ROWMXI
C
C SCALE COLUMNS
C
  DO 10 J=1,N
    COLMX = 0.
    DO 4 I=1,N
      IF (ABS(A(I,J)) .GT. COLMX) COLMX = ABS(A(I,J))
4 CONTINUE
      IF ( COLMX ) 5, 750, 5
5 CONTINUE
      COLMXI = 1./COLMX
      DO 8 I=1,N
        AIJ = A(I,J)
        A(I,J) = (A(I,J) * COLMXI) * BIGPW2
        IF (A(I,J) .EQ. 0.) A(I,J) = (AIJ * BIGPW2) * COLMXI
8 CONTINUE
10 PIVOT(J,3) = BIGPW2 * COLMXI

```

```

      IF (M) 14,24,14
14 DO 22 J=1,M
      COLMX = 0.
      DO 16 I=1,N
      IF (ABS(R(I,J)).GT. COLMX ) COLMX = ABS(R(I,J))
16 CONTINUE
      IF (COLMX ) 17, 22, 17
17 CONTINUE
      SCALER(J) = COLMX / BIGPW2
      COLMXI = 1./COLMX
      DO 20 I=1,M
      RIJ = R(I,J)
      R(I,J) = (R(I,J) * COLMXI) * BIGPW2
      IF (R(I,J) .EQ. 0.) R(I,J) = (RIJ * BIGPW2) * COLMXI
20 CONTINUE
22 CONTINUE
24 CONTINUE
C
C   INITIALIZATION
C
      PMONE=1.
      DO 42 J=1,N
      PIVOT(J,1) = 0.
42 INDEX(J,3) = 0
      DO 550 I=1,N
C
C   SEARCH FOR PIVOT ELEMENT
C
      ABPIVI=0.
45 DO 105 J=1,N
50 IF (INDEX(J,3)-1) 60,105,60
60 DO 100 K=1,N
70 IF (INDEX(K,3)-1) 80,100,80
80 IF (ABS(A(J,K)) - ABPIVI) 100,100,85
85 IROW=J
90 ICOLUM=K
      ABPIVI=ABS(A(J,K))
100 CONTINUE
105 CONTINUE
      IF (I-1) 115,120,115
115 IF ( ABPIVI .GE. PIVMIN ) GO TO 123
120 PIVMIN=ABPIVI
      IF (ABPIVI) 123,750,123
123 CONTINUE
      INDEX(ICOLUM,3)=1
      PIVOTI=A(IROW,ICOLUM)
      PIVOT(I,1) = PIVOTI
C
C   INTERCHANGE ROWS TO PUT PIVOT ELEMENT ON DIAGONAL
C
130 IF (IROW-ICOLUM) 140, 260, 140
140 PMONE=-PMONE
150 DO 200 L=1,N
160 SWAP=A(IROW,L)
170 A(IROW,L)=A(ICOLUM,L)
200 A(ICOLUM,L)=SWAP
205 IF (M) 260, 260, 210
210 DO 250 L=1, M
220 SWAP = R(IROW,L)

```

```

230 R(IROW,L) = B(ICOLUM,L)
250 B(ICOLUM,L) = SWAP
260 INDEX(I,1)=IROW
270 INDEX(I,2)=ICOLUM
C
C   DIVIDE PIVOT ROW BY PIVOT ELEMENT
C
   PIVINV=1.0/PIVOT1
330 A(ICOLUM,ICOLUM) = BIGPW2
340 DO 350 L=1,N
350 A(ICOLUM,L)= A(ICOLUM,L)*PIVINV
355 IF (M) 380, 380, 360
360 DO 370 L=1,M
370 B(ICOLUM,L) = B(ICOLUM,L)*PIVINV
C
C   REDUCE NON-PIVOT ROWS
C
380 DO 550 L1=1,N
390 IF (L1-ICOLUM) 400, 550, 400
400 T=A(L1,ICOLUM)
   IF (T) 420,550,420
420 A/L1,ICOLUM)=0.0
430 DO 450 L=1,N
450 A(L1,L)=A(L1,L)-A(ICOLUM,L)*T
455 IF (M) 550, 550, 460
460 DO 500 L=1,M
500 B(L1,L) = B(L1,L) - B(ICOLUM,L)*T
550 CONTINUE
C
C   INTERCHANGE COLUMNS
C
600 DO 710 I=1,N
610 L=N+1-I
620 IF (INDEX(L,1)-INDEX(L,2)) 630, 710, 630
630 JROW=INDEX(L,1)
640 JCOLUM=INDEX(L,2)
650 DO 705 K=1,N
660 SWAP=A(K,JROW)
670 A(K,JROW)=A(K,JCOLUM)
700 A(K,JCOLUM)=SWAP
705 CONTINUE
710 CONTINUE
C
C   PIVOT(1,1) = PIVOT(1,1) * PMONE
C
SIGDIG = SIGMCH - ALOG10(BIGPW2/PIVMIN)
IF (SIGDIG .LT. .85) SIGDIG = 0.
C
C   UNSCALE INVERSE AND SOLUTION(S)
C
DO 720 J=1,N
ROWMXI = PIVOT(J,2)
DO 720 I=1,N
IF (ROWMXI .LT. 1.) GO TO 715
A(I,J) = (A(I,J) * PIVOT(I,3)) * ROWMXI
GO TO 720
715 A(I,J) = A(I,J) * (PIVOT(I,3) * ROWMXI)
720 CONTINUE
   IF (M) 725, 735, 725

```

```

725 DO 730 J=1,M
    ROWMXI = SCALER(J)
    DO 730 I=1,N
        IF (ROWMXI .LT. 1.) GO TO 728
        B(I,J) = (B(I,J) * PIVOT(I,3)) * ROWMXI
        GO TO 730
728 B(I,J) = B(I,J) * (PIVOT(I,3) * ROWMXI)
730 CONTINUE
735 CONTINUE
C
    IERROR = 1
    RETURN
750 IERROR = -1
    SIGDIG = 0.
    RETURN
    END

```



```

SUBROUTINE LSOPUL (XI,YI,WI,RO,N,SUM,M,COEF)
DIMENSION XI(1),YI(1),WI(1),RO(1),COEF(1)
COMMON / Z/LSQP/
      X(50),Y(50),W(50),B(11),A(11),S(11),DSQ(11)
      ,P(50),PO(50),C(11)
DOUBLE PRECISION X,Y,W,B,A,S,PO,P,DSQ,C,G
DO 100 I = 1,N
X(I) = XI(I)
Y(I) = YI(I)
100 W(I) = WI(I)
CALL ORPOLG(M-1,X,Y,W,N,B,A,S,DSQ,PO,P,G,E)
IF(F) 1100,150,1100
150 CALL ORPOLC(M-1,C,S,A,B,G,P,PO)
DO 200 I = 1,M
200 COEF(I) = C(I)
SUM = 0.0
DO 500 I = 1,N
K = M-1
Z = COEF(M)
IF (K) 300,350,300
300 Z = Z*XI(I) + COEF(K)
K = K-1
IF (K) 350,350,300
350 RO(I) = 7 - YI(I)
500 SUM = SUM + RO(I)**2
RETURN
1100 WRITE (6,1200)
1200 FORMAT(43H1LSOPUL ERROR-MORE COEFFICIENTS THAN POINTS)
STOP
END

```

```

      SUBROUTINE ORPOLG(N,X,Y,W,B,A,S,DSQ,PO,P,GAMMA,ERROR)

      DOUBLE PRECISION X,Y,W,B,A,PO,P,DSQ,WPP,WXPP,WYP,TEMP
      DOUBLE PRECISION WPP0,SUM,GAMMA,1,MEAN,S
      EQUIVALENCE (T,TEMP,SUM)
      DIMENSION X(1),Y(1),W(1),B(1),A(1),S(1),PO(1),P(1),
1      DSQ(1)
      SUM=0.
      DO 6 I=1,M
        SUM=SUM+X(I)
6      MEAN=SUM/FLOAT(M)
      T=0.
      DO 7 I=1,M
        T=AMAX1(T,ABS(X(I)-MEAN))
7      TT = DABS(X(I) - MEAN)
      IF (TT .GT. T) T = TT
      GAMMA=T/2.
      DO 8 I=1,M
        X(I)=(X(I)-MEAN)/GAMMA
8      ERROR=0.
      NO=M-N-1
101      IF(NO)102,103,103
102      ERROR=1.
      GO TO 105
103      B=0.
      DSQ=0.
      WPP=0.
      DO 106 J=1,M
        P(J)=1.
        PO(J)=0.
        WPP=WPP+W(J)
        IF(NO)106,106,107
107      DSQ=DSQ+W(J)*Y(J)**2
106      CONTINUE
      NA=N+1
      DO 109 I=1,NA
        WXPP=0.
        WYP=0.
        DO 110 J=1,M
          TEMP=W(J)*P(J)
          IF(NA-I)111,111,112
112      WXPP=WXPP+TEMP*X(J)*P(J)
111      IF(M - I)110,113,113
113      WYP=WYP+TEMP*Y(J)
110      CONTINUE
          IF(M - I)114,115,115
115      S(I)=WYP/WPP
114      IF(NA-M)666,117,117
666      IF(I-1)102,667,116
667      DSQ=DSQ-S**2*WPP
      GO TO 117
116      DSQ(I)=DSQ(I-1)-S(I)**2*WPP
117      IF(NA-I)109,109,119
119      A(I)=WXPP/WPP
      WPP0=WPP
      WPP=0.
      DO 120 J=1,M
        TEMP=(X(J)-A(I))*P(J)-B(I)*PO(J)
        WPP=WPP+W(J)*TEMP**2

```

```

      PO(J)=P(J)
120  P(J)=TEMP
      R(I+1)=WPP/WPP0
109  CONTINUE
      IF (N) 140,155,140
140  DO 150 I=1,N
      A(I)=A(I)*GAMMA + MEAN
150  B(I)=B(I)*GAMMA
155  DO 160 I=1,M
160  X(I)=GAMMA*X(I)+MEAN
105  RETURN
      END

```

```

SUBROUTINE ORPOLC (N,C,S,A,R,G,P,PM)
DOUBLE PRECISION C,S,A,R,G,P,PM,T1,T2
DIMENSION C(1),S(1),A(1),R(1),P(1),PM(1)
M=N+1
DO 300 L=1,M
C(L)=C.
PM(L)=0.
300 P(L)=0.
P=1.
C=S
IF (N) 301,303,301
301 DO 302 L=1,N
T2=0.
J=L+1
DO 302 M=1,J
T1=(T2-A(L)*P(M)-B(L)*PM(M))/G
T2=P(M)
PM(M)=P(M)
P(M)=T1
302 C(M)=C(M)+T1*S(L+1)
303 RETURN
END

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13. ABSTRACT As part of the Optical Signatures Program, McDonnell Douglas Astronautics Company-West has developed the initial working model of a code to describe the gross features of rocket-plume radiation for altitudes above 75 n mi. The main effort is the construction of a scheme for integration of an arbitrary function through an arbitrary axisymmetric rocket plume, with any specified look angle, plume direction, and vehicle velocity direction. Radiances are presented as integrated values in a specified spectral band. The equations used and a printout of the code and of a sample application are included.		

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